

Last time roots of Bessel functions:

$$J_\nu(x_{\nu n}) = 0, \quad n=1, 2, 3, \dots$$

$x_{\nu n}$  is the  $n$ th root of Bessel function  $J_\nu(x)$

⇒ Bessel functions form a complete & orthonormal set

$$\int_0^a dp \cdot p \cdot J_\nu(x_{\nu n} \frac{p}{a}) J_\nu(x_{\nu m} \frac{p}{a}) = \frac{a^\nu}{2} \delta_{nm} [J_{\nu+1}(x_{\nu n})]^2$$

on the set of functions  $f(p)$ ,  $p \in [0, a]$ ,  $f(0) = 0 = f(a)$ .

(Defined) modified Bessel functions:

$$\begin{aligned} I_\nu(x) &\equiv i^{-\nu} J_\nu(ix) \\ K_\nu(x) &\equiv \frac{\pi}{2} i^{\nu+1} [J_\nu(ix) + i N_\nu(ix)] \end{aligned}$$

Note that:  $I_\nu(x) \Big|_{x \ll 1} = \frac{1}{\Gamma(\nu+1)} \left(\frac{x}{2}\right)^\nu$

$$K_\nu(x) \Big|_{x \ll 1} = \begin{cases} -\ln \frac{x}{2} - 0.5772\dots, & \nu = 0 \\ \frac{\Gamma(\nu)}{2} \left(\frac{2}{x}\right)^\nu, & \nu \neq 0 \end{cases}$$

$$I_\nu(x) \Big|_{x \gg 1} \simeq \frac{1}{\sqrt{2\pi x}} e^x$$

$$K_\nu(x) \Big|_{x \gg 1} \simeq \sqrt{\frac{\pi}{2x}} e^{-x}$$



Another useful special functions are  
modified Bessel functions  $I_\nu(z)$  &  $K_\nu(z)$

(Definition)

$$I_\nu(x) = e^{-\nu} J_\nu(ix)$$

$$K_\nu(x) = \frac{\pi}{2} i^{\nu+1} [J_\nu(ix) + i N_\nu(ix)]$$

obey  $\frac{d^2 R}{dx^2} + \frac{1}{x} \frac{dR}{dx} - \left(1 + \frac{\nu^2}{x^2}\right) R = 0$  diff. equation.

Very useful formula:

$$\frac{1}{k} \delta(k - k') = \int_0^\infty dx \cdot x \cdot J_\nu(kx) J_\nu(k'x)$$

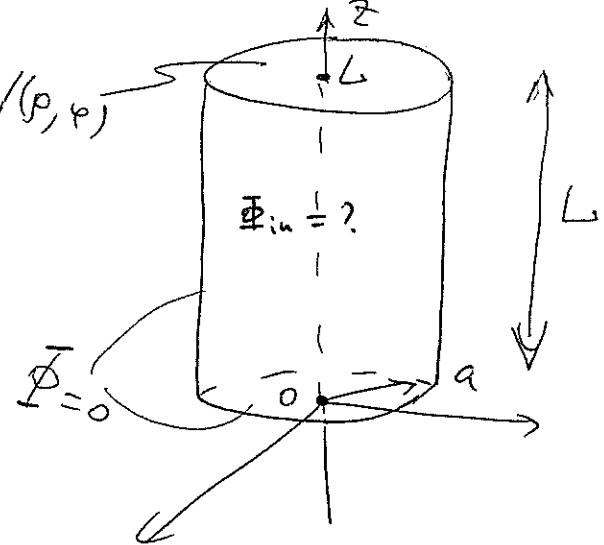
(Example of a Boundary-Value Problem:

$$Q(\varphi) = A \sin(m\varphi) + B \cos(m\varphi)$$

$$Z(z) = C \sinh(Kz) + D \cosh(Kz)$$

$$\Rightarrow D = 0 \text{ as } Z(0) = 0.$$

$$R(r) = E J_m(kr) + F N_m(kr)$$



$$\Rightarrow \text{finite at } r=0 \Rightarrow F=0$$

$$R(a)=0 \Rightarrow K \rightarrow k_{mn} = \frac{x_{mn}}{a} \text{ - roots}$$

$$\Rightarrow \Phi(r, \varphi, z) = \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} J_m(k_{mn}r) \sinh(k_{mn}z) \cdot [A_{mn} \sin(m\varphi) + B_{mn} \cos(m\varphi)].$$

Finally,  $\Phi(s, \varphi, z=L) = V(p, \varphi)$

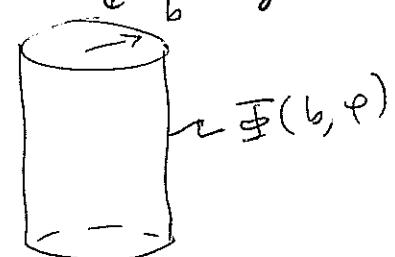
$$\Rightarrow V(p, \varphi) = \sum_{m,n} J_m(K_{mn} p) \sinh(K_{mn} L) [A_{mn} \sin(m\varphi) + B_{mn} \cos(m\varphi)] \Rightarrow \text{invert the Fourier and Fourier-Bessel } (\overset{\circ}{p}) \text{ series to get}$$

$$\begin{pmatrix} A_{mn} \\ B_{mn} \end{pmatrix} = \frac{2}{\pi a^2 \sinh(K_{mn} L) J_{m+1}^2(K_{mn} a)} \int_0^{2\pi} d\varphi \cdot \int_0^\infty dp \cdot p \cdot V(p, \varphi) J_m(K_{mn} p) \begin{pmatrix} \sin(m\varphi) \\ \cos(m\varphi) \end{pmatrix}.$$

for  $m=0$  use  $2B_{0n}$  on the right-hand side  
(that is, divide  $B_{0n}$  from the f-ls by 2 to obtain true  $B_{0n}$ )

Jackson problem 2.12: infinitely long cylinder

$$\begin{aligned} \Phi(p, \varphi) &= A_0 + b_0 \ln p + \sum_{n=1}^{\infty} a_n p^n \sin(n\varphi + \alpha_n) + \\ &+ \sum_{n=1}^{\infty} b_n p^{-n} \sin(n\varphi + \beta_n) \end{aligned}$$



$\Rightarrow b_n = 0$  as  $\Phi$  is finite at  $p=0$  ( $b_0 = 0 + \infty$ )

$$\Rightarrow \Phi(p, \varphi) = \frac{A_0}{2} + \sum_{n=1}^{\infty} (A_n \cos(n\varphi) + B_n \sin(n\varphi)) p^n$$

For  $p=b$  we have

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$$\Phi(b, \varphi) = \frac{A_0}{2} + \sum_{m=1}^{\infty} (A_m \cos(m\varphi) + B_m \sin(m\varphi)) b^m$$

$$\Rightarrow \begin{pmatrix} A_m \\ B_m \end{pmatrix} b^m = \frac{1}{\pi} \int_0^{2\pi} d\varphi' \begin{pmatrix} \cos(m\varphi') \\ \sin(m\varphi') \end{pmatrix} \Phi(b, \varphi'), \quad A_0 = \frac{1}{\pi} \int_0^{2\pi} d\varphi' \Phi(b, \varphi')$$

$$\Rightarrow \Phi(p, \varphi) = \frac{1}{\pi} \int_0^{2\pi} d\varphi' \Phi(b, \varphi') \sum_{m=1}^{\infty} [\cos m\varphi \cos m\varphi' +$$

$$+ \sin m\varphi \sin m\varphi'] \cdot b^{-m} p^m + \frac{1}{2\pi} \int_0^{2\pi} d\varphi' \Phi(b, \varphi').$$

$$\text{Now, } [\quad] = \cos m(\varphi - \varphi') = \frac{1}{2} [e^{im(\varphi-\varphi')} + e^{-im(\varphi-\varphi)}]$$

$$\Rightarrow \sum_{m=1}^{\infty} e^{im(\varphi-\varphi')} \left(\frac{p}{b}\right)^m = \frac{p}{b} \cdot e^{i(\varphi-\varphi')} \frac{1}{1 - \frac{p}{b} e^{i(\varphi-\varphi')}} =$$

$$= \frac{1}{\frac{b}{p} e^{-i(\varphi-\varphi')} - 1} \Rightarrow \sum_{m=1}^{\infty} \left(\frac{p}{b}\right)^m \cos m(\varphi - \varphi') =$$

$$= \frac{1}{2} \left[ \frac{1}{\frac{b}{p} e^{-i(\varphi-\varphi')} - 1} + \frac{1}{\frac{b}{p} e^{i(\varphi-\varphi')} - 1} \right] = \frac{\frac{b}{p} \cos(\varphi - \varphi') - 1}{1 + \frac{b^2}{p^2} - 2 \frac{b}{p} \cos(\varphi - \varphi')}$$

$$= \frac{b p \cos(\varphi - \varphi') - p^2}{p^2 + b^2 - 2 p b \cos(\varphi - \varphi')}$$

$$\text{So, } \Phi(p, \varphi) = \frac{1}{2\pi} \int_0^{2\pi} d\varphi' \Phi(b, \varphi') \left[ 1 + 2 \frac{b p \cos(\varphi - \varphi') - p^2}{b^2 + p^2 - 2 p b \cos(\varphi - \varphi')} \right]$$

$$= \frac{1}{2\pi} \int_0^{2\pi} d\varphi' \Phi(b, \varphi') \frac{b - p^c}{b^2 + p^2 - 2bp \cos(\varphi - \varphi')} \quad \text{as desired!} \quad (163)$$

Green function in cylindrical coordinates:

need to solve  $\nabla^2 G(\vec{x}, \vec{x}') = -4\pi \delta^3(\vec{x} - \vec{x}') =$

$$= -\frac{4\pi}{\rho} \delta(\rho - \rho') \delta(\varphi - \varphi') \delta(z - z').$$

write  $\delta(z - z') = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ik(z-z')} = \int_0^{\infty} \frac{dk}{\pi} \cos[k(z-z')]$

$$\delta(\varphi - \varphi') = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{im(\varphi - \varphi')}$$

$$\Rightarrow G(\vec{x}, \vec{x}') = \frac{1}{2\pi^2} \sum_{m=-\infty}^{\infty} \int_0^{\infty} dk e^{im(\varphi - \varphi')} \cos[k(z - z')].$$

$$g_m(k, \rho, \rho')$$

Plug this back into eqns for  $G$  to get

$$\frac{1}{\rho} \frac{d}{dp} \left( \rho \frac{d}{dp} g_m(k, \rho, \rho') \right) - \left( k^2 + \frac{m^2}{\rho^2} \right) g_m = -\frac{4\pi}{\rho} \delta(\rho - \rho')$$

Same story as in rectangular coord's:

$$\text{if } \rho < \rho' \Rightarrow \text{get } g_m \sim A I_m(k\rho) + B K_m(k\rho)$$

$\Rightarrow$  we want regular behavior as  $\rho \rightarrow 0 \Rightarrow$

$$\Rightarrow g_m \sim I_m(k\rho) \text{ for } \rho < \rho'$$