

General Solution of Maxwell Equations (in Lorenz gauge)

We want to learn to solve Maxwell equations in a general (dynamic) case:

$$\partial_\mu F^{\mu\nu} = \frac{4\pi}{c} J^\nu \quad (\text{Gauss units})$$

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu \Rightarrow$$

$$\partial_\mu \partial^\mu A^\nu - \partial^\nu \partial_\mu A^\mu = \frac{4\pi}{c} J^\nu$$

Choose $\partial_\mu A^\mu = 0$ Lorenz gauge \Rightarrow

Maxwell equations become

$$\square A^\nu = \frac{4\pi}{c} J^\nu$$

where $\square \equiv \partial_\mu \partial^\mu = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2$. We get

$$\left[\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2 \right] \Phi = 4\pi \rho$$

$$\left[\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2 \right] \vec{A} = \frac{4\pi}{c} \vec{J}$$

Maxwell equations
in Lorenz gauge
(Gauss units).

In SI units get

$$\left[\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2 \right] \Phi = \frac{\rho}{\epsilon_0}$$

$$\left[\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \vec{\nabla}^2 \right] \vec{A} = \mu_0 \vec{J}$$

\sim inhomogeneous wave equations

To solve Maxwell equations for arbitrary ρ & \vec{J} we need the Green function for the \square operator.

$i = \sqrt{-1}$ or $i^2 = -1$ imaginary unit number

$z = x + iy$; $\bar{z} = x - iy$ ~ complex conjugate
(x, y ~ real)

Def. $f(z)$ is analytic at a point z_0 if it is and single-valued differentiable in a neighborhood of z_0 . ($\frac{\partial f}{\partial z}$ exists).

$$\frac{\Delta f}{\Delta z} = \frac{f(z_0 + \Delta z) - f(z_0)}{\Delta z} \Rightarrow \text{assume } f(z) = u(x, y) + i v(x, y)$$

u, v ~ real

$$\frac{\Delta f}{\Delta z} = \frac{\left(\frac{\partial u}{\partial x} + i \frac{\partial v}{\partial x} \right) \Delta x + \left(\frac{\partial u}{\partial y} + i \frac{\partial v}{\partial y} \right) \Delta y}{\Delta x + i \Delta y} \Rightarrow$$

want this independent of direction of the

derivative \Rightarrow

$$\frac{\frac{\partial u}{\partial x} + i \frac{\partial v}{\partial x}}{\frac{\partial u}{\partial y} + i \frac{\partial v}{\partial y}} = \frac{1}{i} \Rightarrow$$

$$\Rightarrow i \frac{\partial u}{\partial x} - \frac{\partial v}{\partial x} = \frac{\partial u}{\partial y} + i \frac{\partial v}{\partial y} \Rightarrow$$

$$\Rightarrow \boxed{\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}}$$

$$\text{and } \boxed{\frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}}$$

Cauchy-Riemann conditions

note that if C-R conditions are satisfied \Rightarrow

$$\Rightarrow \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 v}{\partial x \partial y} - \frac{\partial^2 v}{\partial x \partial y} = 0 \Rightarrow \Delta^2 u = 0 \text{ also } \Delta^2 v = 0$$

$\Rightarrow \frac{\partial^2}{\partial x^2 + \partial y^2} f = 0$

harmonic functions

Cauchy Theorem: if $f(z)$ is analytic $\Rightarrow \oint_C f(z) dz = 0$ (220)



Cauchy Formula:

$$f(z_0) = \frac{1}{2\pi i} \oint_C \frac{f(z)}{z - z_0} dz$$

$f(z)$ analytic function



Define a residue: residue of $f(z)$ at z_0 is

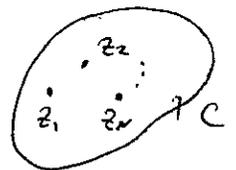
$$\frac{1}{2\pi i} \oint_C f(z) dz = \text{Res } f(z_0)$$



Residue Theorem | $f(z)$ is analytic except for a finite number of isolated singularities z_1, \dots, z_N .

Then

$$\oint_C f(z) dz = 2\pi i \sum_{n=1}^N \text{Res } f(z_n)$$



How to find residues: simple pole $\sim \frac{1}{z-a}$

$$\Rightarrow \text{Res } f(a) = \lim_{z \rightarrow a} (z-a) f(z)$$

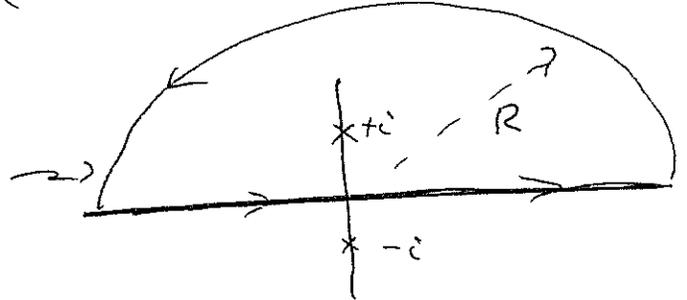
n th order pole: $\frac{1}{(z-a)^n}$

$$\text{Res } f(a) = \lim_{z \rightarrow a} \frac{1}{(n-1)!} \frac{d^{n-1}}{dz^{n-1}} [(z-a)^n f(z)]$$

Example $\int_{-\infty}^{\infty} \frac{dx}{x^2+1} = \arctan x \Big|_{-\infty}^{\infty} = \frac{\pi}{2} - \left(-\frac{\pi}{2}\right) = \pi$ (221)

Residues: $\int_{-\infty}^{\infty} \frac{dx}{x^2+1} = \int_{-\infty}^{\infty} \frac{dx}{(x-i)(x+i)}$

close the contour



if $f(z) \rightarrow 0$ as $z \rightarrow \infty$

faster than $\frac{1}{z}$ (i.e. $f(z) \sim \frac{1}{z^{1+\delta}}$, $\delta > 0$ as $z \rightarrow \infty$)

or $z f(z) \rightarrow 0$ as $z \rightarrow \infty$

\Rightarrow integral over the semi-circle is zero when $R \rightarrow \infty \Rightarrow$ turned line integral into a contour integral

\Rightarrow using residue theorem get $\int_{-\infty}^{\infty} dx \frac{1}{(x-i)(x+i)} =$

$= 2\pi i \lim_{z \rightarrow i} \frac{z-i}{(z-i)(z+i)} = 2\pi i \frac{1}{2i} = \pi$, as desired!

Proof of Cauchy theorem:

$$\oint_C f(z) dz = \oint_C [u + iv] [dx + i dy] = \oint_C [u dx - v dy]$$

$$+ i \oint_C [v dx + u dy]$$

Stokes' theorem: $\oint_C \vec{A} \cdot d\vec{\ell} = \int da \hat{n} \cdot (\vec{\nabla} \times \vec{A})$

\Rightarrow if C is in the x - y plane \Rightarrow

$$\oint_C [A_x dx + A_y dy] = \int dx dy [\partial_x A_y - \partial_y A_x]$$

$$\Rightarrow \oint_C [u dx - v dy] = \int dx dy [-\partial_x v - \partial_y u] = 0$$

↑
due to Cauchy-Riemann conditions

$$\oint_C [v dx + u dy] = \int dx dy [\partial_x u - \partial_y v] = 0$$

↓

\Rightarrow for an analytic function

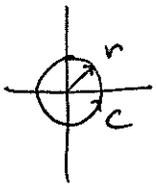
$$\oint_C f(z) dz = 0, \text{ as required.}$$

Proof of Cauchy formula:

(223)

Warmup: $\oint_C dz \cdot z^n = \int_0^{2\pi} r e^{i\varphi} d\varphi \cdot i r^n e^{ni\varphi} = i r^{n+1} \int_0^{2\pi} d\varphi e^{i(n+1)\varphi}$

\uparrow
 $z = r e^{i\varphi}$



$$= i r^{n+1} \frac{1}{i(n+1)} \left[e^{2\pi i(n+1)} - 1 \right] = \begin{cases} 0, & n \neq -1, \\ & n \text{ integer} \\ 2\pi i, & n = -1 \end{cases}$$

$$\Rightarrow \oint_C dz (z - z_0)^n = 2\pi i \delta_{n,-1}$$

$n = -1 \Rightarrow$ simple pole



proof: $f(z)$ is analytic on and inside contour C .

$$\oint_C \frac{f(z)}{z - z_0} dz = \int_0^{2\pi} \frac{f(z_0 + r e^{i\varphi})}{r e^{i\varphi}} r e^{i\varphi} i d\varphi = i \int_0^{2\pi} d\varphi f(z_0 + r e^{i\varphi})$$

\uparrow
 $z = z_0 + r e^{i\varphi}$

$\xrightarrow{r \rightarrow 0}$

$$2\pi i f(z_0)$$

\Rightarrow

$$\oint_C dz \frac{f(z)}{z - z_0} = 2\pi i f(z_0)$$

as desired.

Taylor series: $f(z) = \sum_{n=0}^{\infty} \frac{f^{(n)}(z_0)}{n!} (z - z_0)^n$

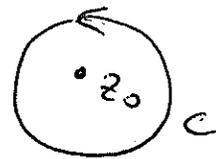
if $f(z)$ is expandable in the Taylor series \Leftrightarrow analytic

Laurent series: $f(z) = \sum_{n=-\infty}^{\infty} a_n (z - z_0)^n$

(if $a_n \neq 0$ for $n \leq -1 \Rightarrow f(z)$ is not analytic at $z = z_0$)

$$\oint_C dz f(z) = \sum_{n=-\infty}^{\infty} a_n \oint_C dz (z-z_0)^n$$

$$= \sum_{n=-\infty}^{\infty} a_n \cdot 2\pi i \delta_{n,-1} = 2\pi i a_{-1}$$

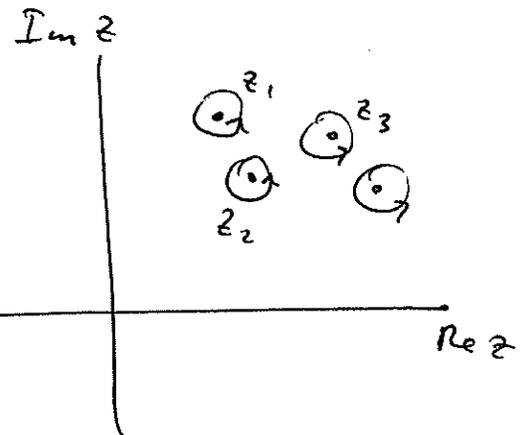
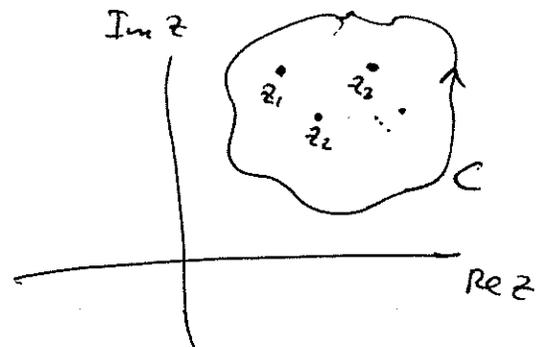


$\Rightarrow a_{-1}$ is called the residue of $f(z)$ at $z=z_0$.

Residue = coefficient of $\frac{1}{z-z_0}$ in Laurent series.

Proof of Residue Theorem:

Can deform the initial contour C to small circles around all the poles inside C :



Hence set

$$\oint_C dz f(z) = 2\pi i [a_{-1}^{(1)} + a_{-1}^{(2)} + \dots]$$

$$\Rightarrow \boxed{\oint_C dz f(z) = 2\pi i \sum_{i=1}^N \text{Res } f(z_i)}$$