

Last time

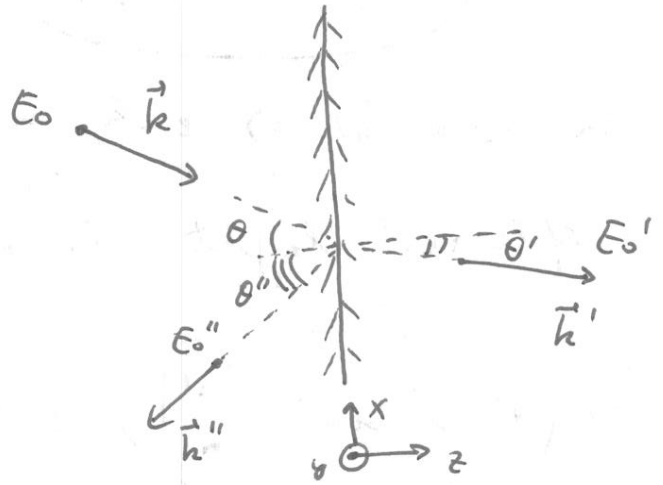
Reflection and Refraction

$$\omega = \omega' = \omega''$$

$$k = k'' = \sqrt{\mu\epsilon} \omega, \quad k' = \sqrt{\mu'\epsilon'} \omega$$

$$\theta = \theta''$$

$$n \sin \theta = n' \sin \theta'$$



① $\vec{E}_0 \parallel \hat{y}$ ($\vec{E}_0 \perp$ plane of incidence)

$$\frac{E_0'}{E_0} = \frac{2n \cos \theta}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

$$\frac{E_0''}{E_0} = \frac{n \cos \theta - \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

② \vec{E}_0 in xz -plane (plane of incidence)

$$\frac{E_0'}{E_0} = \frac{2n n' \cos \theta}{\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

$$\frac{E_0''}{E_0} = \frac{-\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}{\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

\Rightarrow for $\theta_B = \tan^{-1}\left(\frac{n'}{n}\right)$ (for $\mu = \mu'$) the reflected wave vanishes

in case (ii) \Rightarrow only have polarization \perp xz-plane,
reflected light is polarized

$\Rightarrow \theta > \sin^{-1}\left(\frac{n'}{n}\right) \sim$ total internal reflection

$$\Rightarrow \vec{E} \sim e^{-|k_z|z}$$

evanescent wave



if $\theta = 0$ get

$$\frac{E_0'}{E_0} = \frac{2n}{n + \frac{\mu}{\mu'} n'}$$

$$\frac{E_0''}{E_0} = \frac{n - \frac{\mu}{\mu'} n'}{n + \frac{\mu}{\mu'} n'}$$

transmission coefficient

$$T = \frac{|\vec{S}'|}{|\vec{S}|} = \sqrt{\frac{\mu \epsilon'}{\mu' \epsilon}} \frac{4n^2}{\left(n + \frac{\mu}{\mu'} n'\right)^2}$$

$$R = \frac{|\vec{S}''|}{|\vec{S}|} = \left(\frac{n - \frac{\mu}{\mu'} n'}{n + \frac{\mu}{\mu'} n'} \right)^2 \quad (\theta = 0)$$

in case (I) $\frac{E_o''}{E_o}$ never vanishes (always < 0)
if $n' > n$

in case (II) $\frac{E_o''}{E_o} = 0$ for $\theta_B = \tan^{-1}\left(\frac{n'}{n}\right)$ Brewster's angle

$$\left[\begin{aligned} n^2 n'^2 - n^4 \sin^2 \theta &= n'^4 \cos^2 \theta \Rightarrow \\ \Rightarrow n^2 n'^2 (1 + \tan^2 \theta) - n^4 \tan^2 \theta &= n'^4 \Rightarrow n^2 n'^2 - n^4 = n^2 (n'^2 - n^2) \tan^2 \theta \Rightarrow \tan^2 \theta = \frac{n'^2}{n^2} \end{aligned} \right]$$

\Rightarrow reflected light is polarized.

if $\theta = \theta_B \Rightarrow$ polarization is linear, \perp to the plane of incidence.

(fish in the ocean reflect light \sim squids with polarized vision can see them)

Total internal reflection: Snell's law:

$$n \sin \theta = n' \sin \theta' \leq n' \Rightarrow \theta \leq \sin^{-1}\left(\frac{n'}{n}\right) \Rightarrow$$

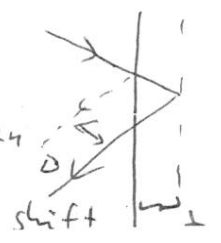
\Rightarrow for $\theta > \sin^{-1}\left(\frac{n'}{n}\right)$ get total reflection "evanescent wave"

$$\begin{aligned} k' &= \sqrt{\mu' \epsilon'} \omega \Rightarrow k'_x = -\sqrt{\mu' \epsilon'} \omega \sin \theta' = -\sqrt{\mu_0 \epsilon_0} \omega n' \sin \theta' = \\ & \quad \frac{n' \omega}{c} \qquad \qquad \qquad k'_y = 0 \qquad \qquad \qquad = -\frac{\omega}{c} n \sin \theta \\ k'_z &= \sqrt{k'^2 - k_x'^2} = \sqrt{n'^2 - n^2 \sin^2 \theta} \frac{\omega}{c} \end{aligned}$$

\Rightarrow for $\theta > \sin^{-1}\left(\frac{n'}{n}\right)$: k'_z becomes imaginary $k'_z = +i|k'_z|$

$\Rightarrow e^{i k'_z z} \sim e^{-|k'_z| z} \sim$ exponential falloff

effectively ^{the wave} gets reflected from a different surface \sim violation of geom. optics, Goos-Hänchen effect



transmission coefficient $\vec{S} = \vec{E} \times \vec{H} = \text{Re}[\vec{E} \times \vec{H}^*]$

$$T = \frac{|\vec{S}'|}{|\vec{S}|} = \frac{E_0' H_0' \cdot \frac{1}{2}}{E_0 H_0 \cdot \frac{1}{2}} = \frac{\mu}{\mu'} \frac{E_0' B_0'}{E_0 B_0} = \frac{\mu}{\mu'} \frac{\sqrt{\mu' \epsilon'} (E_0')^2}{\sqrt{\mu \epsilon} (E_0)^2}$$

$\uparrow \langle \cos^2 \rangle$ phase $T = \frac{4n n'}{(n+n')^2}$; $R = \left(\frac{n-n'}{n+n'}\right)^2$

$= |_{\theta=0} = \sqrt{\frac{\mu \epsilon'}{\mu' \epsilon}} \cdot \frac{4n^2}{(n + \frac{\mu}{\mu'} n')^2}$ \nearrow fraction of incident power that got through

reflection coefficient ~ fraction of inc. power reflected. (T+R=1)

$$R = \frac{|\vec{S}''|}{|\vec{S}|} = \frac{E_0'' H_0''}{E_0 H_0} = \frac{E_0'' B_0''}{E_0 B_0} = \frac{|E_0''|^2}{|E_0|^2} = \left(\frac{n - \frac{\mu}{\mu'} n'}{n + \frac{\mu}{\mu'} n'}\right)^2$$

Electromagnetic Waves in Conductors

Maxwell equations: $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ $\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$

$\vec{\nabla} \cdot \vec{D} = \rho = 0$ $\vec{\nabla} \cdot \vec{B} = 0$

\uparrow assume no free charges

Ohm's law: $\vec{J} = \sigma \vec{E}$; $\vec{B} = \mu \vec{H}$, $\vec{D} = \epsilon \vec{E}$

Look for plane-wave solutions:

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{x} - \omega t)} ; \vec{B} = \vec{B}_0 e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

$$\Rightarrow \vec{\nabla} \times \vec{E} = i\omega \vec{B} \quad \vec{\nabla} \times \vec{H} = \underbrace{\sigma \vec{E} - i\omega \epsilon \vec{E}}_{-i\omega(\epsilon + \frac{i}{\omega} \sigma)}$$

$$i\omega \mu \vec{H}$$

$$\Rightarrow \vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = -\nabla^2 \vec{E} = i\omega \mu \vec{\nabla} \times \vec{H} = i\omega \mu (\sigma - i\omega \epsilon) \vec{E}$$

=> we get $(\nabla^2 + k^2) \vec{E} = 0$ with $k^2 = \mu \epsilon \omega^2 + i \omega \mu \sigma$

=> $k = \pm \sqrt{\mu \epsilon} \omega \sqrt{1 + \frac{i \sigma}{\epsilon \omega}} \equiv k_1 + i k_2$

As $\vec{\nabla} \cdot \vec{E} = 0 \Rightarrow \vec{k} \cdot \vec{E}_0 = 0$. ~ still transverse.
 $\vec{\nabla} \cdot \vec{B} = 0 \Rightarrow \vec{k} \cdot \vec{B}_0 = 0$.

Good conductor: $\frac{\sigma}{\epsilon \omega} \gg 1$

Bad conductor: $\frac{\sigma}{\epsilon \omega} \ll 1$

Assume that $\vec{k} \parallel \hat{z}$ and $\vec{E} \parallel \hat{x}$ (linear polarization)

$\vec{E} = \hat{x} E_0 e^{i(kz - \omega t)} = \hat{x} E_0 e^{i k_1 z - k_2 z - i \omega t} = \vec{E}$

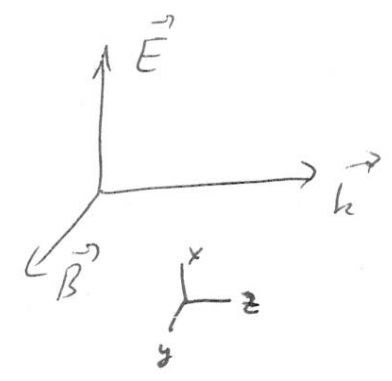
=> wavelength $\lambda = \frac{2\pi}{k_1}$

$E \sim e^{-k_2 z}$ ~ exponentially decaying

\vec{B} -field: $\vec{B} = \frac{-i}{\omega} \vec{\nabla} \times \vec{E} = \frac{1}{\omega} \vec{k} \times \vec{E}_0 e^{i(kz - \omega t)} \Rightarrow$

$\vec{B} = \frac{k}{\omega} \hat{y} E_0 e^{i(kz - \omega t)}$

=> $\vec{B} \perp \vec{k}$, $\vec{B} \perp \vec{E}$



but: as k is complex, \vec{B} & \vec{E}

are out of phase.

Time-averaged Poynting vector:

$\langle S_z \rangle = \frac{1}{2} \text{Re} (\vec{E} \times \vec{H}^*)_z = \frac{1}{2\mu} \text{Re} \left(\frac{k^*}{\omega} |E_0|^2 e^{-2k_2 z} \right) =$

$$= \frac{1}{2\mu} \frac{k_1 |E_0|^2}{\omega} e^{-2k_2 z} \propto e^{-z/\delta} \Rightarrow \delta = \frac{1}{2k_2} \text{ "skin depth"}$$

$$k_2 = \text{Im} \left[\sqrt{\mu \epsilon} \omega \sqrt{1 + \frac{i\sigma}{\epsilon \omega}} \right]$$

Bad conductor: $k_2 \approx \sqrt{\mu \epsilon} \omega \frac{\sigma}{2\epsilon \omega} = \sqrt{\frac{\mu}{\epsilon}} \frac{\sigma}{2}$

$$\Rightarrow \delta = \frac{1}{\sigma} \sqrt{\frac{\epsilon}{\mu}} \text{ (if } \sigma \text{ is small } \Rightarrow \delta \text{ is large)}$$

Good conductor: $k_2 \approx \sqrt{\mu \epsilon} \omega \frac{1}{\sqrt{2}} \sqrt{\frac{\sigma}{\epsilon \omega}} = \sqrt{\frac{\sigma \mu \omega}{2}}$

$$\Rightarrow \delta = \sqrt{\frac{2}{\sigma \mu \omega}} \text{ (if } \sigma \text{ is large } \Rightarrow \delta \text{ is small)}$$

Frequency - dependent ϵ, μ, σ .

we just showed that $k = \sqrt{\mu \epsilon} \omega \sqrt{1 + \frac{i\sigma}{\epsilon \omega}} =$

$$= \omega \sqrt{\mu \left(\epsilon + \frac{i\sigma}{\omega} \right)} \Rightarrow \text{if we want } k = \sqrt{\mu \epsilon} \omega$$

as in non-conductors, we have to, in general, assume that $\epsilon = \epsilon(\omega)$, $\mu = \mu(\omega)$, $\sigma = \sigma(\omega)$

and, here redefine $\epsilon \rightarrow \epsilon(\omega) = \epsilon_0 + \frac{i\sigma}{\omega}$ due to bound charges

"Complex dielectric function" (not a constant!)

$$n(\omega) = \sqrt{\frac{\mu(\omega) \epsilon(\omega)}{\mu_0 \epsilon_0}} \sim \text{"complex index of refraction"}$$