

Last time

## The Lagrangian and the Action (cont'd)

Constructed the Lagrangian for a relativistic particle in electromagnetic field:

$$L = -mc^2 \sqrt{1 - \frac{v^2}{c^2}} - e\Phi + \frac{e}{c} \vec{v} \cdot \vec{A}$$

Wrote down equations of motion for this Lagrangian:

$$\frac{dp^i}{dt} = \frac{e}{c} \left[ c \left( \frac{\partial A^0}{\partial x_i} - \frac{\partial A^i}{\partial x^0} \right) - v^j \left( \frac{\partial A^j}{\partial x_i} - \frac{\partial A^i}{\partial x_j} \right) \right]$$

(Def.) Electric field  $\vec{E} = -\vec{\nabla}\Phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \Rightarrow E^i = \partial^i A^0 - \partial^0 A^i$

(Def.) Magnetic field  $\vec{B} = \vec{\nabla} \times \vec{A} \Rightarrow \partial^j A^i - \partial^i A^j = \epsilon^{ijk} B^k$

$\Rightarrow$  using these definitions we wrote the EOM as

$$\frac{d\vec{p}}{dt} = e\vec{E} + \frac{e}{c} \vec{v} \times \vec{B} \quad \text{Lorentz force!}$$

$\Rightarrow$  We obtained the Lorentz force from the Lorentz invariance (④ homogeneity & isotropy of space-time ⑤ existence of a 4-vector field  $A^\mu(x)$ ).



Def.

Field-strength tensor:

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$F_{\mu\nu} = -F_{\nu\mu} \text{ anti-symmetric}$$

Components are:  $F^{0i} = \partial^0 A^i - \partial^i A^0 =$

$$= \frac{\partial A^i}{\partial x^0} - \frac{\partial A^0}{\partial x^i} = -E^i \Rightarrow F^{0i} = -E^i$$

$$F^{ij} = \partial^i A^j - \partial^j A^i = \frac{\partial A^j}{\partial x_i} - \frac{\partial A^i}{\partial x_j} = -\epsilon^{ijk} B^k$$

$$\Rightarrow F^{ij} = -\epsilon^{ijk} B^k$$

$\Rightarrow$  components of the field-strength tensor are the electric and magnetic fields!

$$F^{0i} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix}, \quad \text{where } \begin{aligned} E_x &= E^1, E_y = E^2, \\ E_z &= E^3, \\ B_x &= B^1, B_y = B^2, B_z = B^3. \end{aligned}$$

$$F_{\mu\nu} = \begin{pmatrix} 0 & E_x & E_y & E_z \\ -E_x & 0 & -B_z & B_y \\ -E_y & B_z & 0 & -B_x \\ -E_z & -B_y & B_x & 0 \end{pmatrix}$$

will revisit  
later

$\Rightarrow \vec{E}$  &  $\vec{B}$  fields transform as components of a rank-2 tensor under Lorentz transformations!

Using  $F^{\mu\nu}$  we can write the Lorentz force (30)  
as

$$\frac{dp^i}{dt} = \frac{e}{c} \left[ -c F^{0i} - v^j F^{ij} \right] = \frac{e}{c} \left[ c F^{i0} - v^j F^{ij} \right]$$

$$= \frac{e}{c} \frac{1}{\gamma} u_\mu F^{i\mu} \Rightarrow \text{as } dt = \gamma d\tau$$

$$\Rightarrow \boxed{\frac{dp^i}{d\tau} = \frac{e}{c} u_\mu F^{i\mu}}$$

What about  $\frac{dp^0}{d\tau}$ ? as  $p^\mu = m u^\mu$  and  $u_\mu u^\mu = c^2$

$$\Rightarrow u_\mu \frac{dp^0}{d\tau} = 0 \quad (\text{why?}) \Rightarrow c \frac{dp^0}{d\tau} - v^i \frac{dp^i}{d\tau} = 0$$

$$\Rightarrow \frac{dp^0}{d\tau} = \frac{1}{c} v^i \frac{dp^i}{d\tau} = \frac{1}{c} \frac{e}{c} \gamma v^i \left[ c F^{i0} - v^j F^{ij} \right]$$

as  $v^i v^j F^{ij} = 0$

$$= \frac{e}{c} \gamma v^i F^{i0} = \frac{e}{c} u_\mu F^{0\mu} \leftarrow (\text{as } F^{00} = 0)$$

$$\Rightarrow \boxed{\frac{dp^0}{d\tau} = \frac{e}{c} u_\mu F^{0\mu}}$$

$\Rightarrow$  Can combine the two into

$$\boxed{\frac{dp^\mu}{d\tau} = \frac{e}{c} u_\nu F^{\mu\nu}}$$

Lorentz-covariant form  
of the Lorentz force

In the non-relativistic notation we write for  $\rho^o$ : (31)

$$\frac{dp^o}{dt} = \frac{1}{c} v^i \frac{dp^i}{dt} \Rightarrow \frac{dp^o}{dt} = \frac{1}{c} v^i \frac{dp^i}{dt} \Rightarrow \text{as } \rho^o = \frac{E \text{ energy}}{c \text{ of the particle}}$$

$$\Rightarrow \frac{dE}{dt} = \vec{v} \cdot \frac{d\vec{p}}{dt} = \vec{v} \cdot \left[ e\vec{E} + \frac{e}{c} \vec{v} \times \vec{B} \right] = e \vec{v} \cdot \vec{E}$$

$$\Rightarrow \boxed{\frac{dE}{dt} = e \vec{v} \cdot \vec{E}} \Rightarrow \text{only electric field does the work on the particle.}$$

(Note that  $\frac{dE}{dt} = \vec{F} \cdot \vec{v}$  as we saw before  $\Rightarrow$

$\Rightarrow$  as  $\vec{F} = \frac{d\vec{p}}{dt}$  (the force)  $\Rightarrow$  above is consistent with the formulas for energy change studied earlier.)

We can also construct the Hamiltonian for the particle in the E&M fields. Since the canonical momentum is  $\vec{P} = \vec{p} + \frac{e}{c} \vec{A}$

$$\Rightarrow H = \vec{P} \cdot \vec{v} - L \quad \begin{matrix} \\ \text{with } \vec{v} \text{ replaced} \\ \text{by } \vec{P} \end{matrix} \Rightarrow$$

$$\text{as } \vec{P} = \frac{m\vec{v}}{\sqrt{1 - \frac{v^2}{c^2}}} + \frac{e}{c} \vec{A} \Rightarrow \left( \vec{P} - \frac{e}{c} \vec{A} \right)^2 = \frac{m^2 v^2}{1 - \frac{v^2}{c^2}}$$

$$\Rightarrow v^2 \left[ m^2 + \frac{1}{c^2} (\vec{P} - \frac{e}{c} \vec{A})^2 \right] = (\vec{P} - \frac{e}{c} \vec{A})^2$$

$$\Rightarrow \vec{v} = \frac{c(\vec{P} - \frac{e}{c} \vec{A})}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} \quad (\text{also need that } \vec{v} \parallel \vec{P} - \frac{e}{c} \vec{A})$$

$$\Rightarrow H = \vec{P} \cdot \frac{c(\vec{P} - \frac{e}{c} \vec{A})}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} + mc^2 \sqrt{1 - \frac{v^2}{c^2}} + e \Phi - \frac{e}{c} \vec{v} \cdot \vec{A}$$

$$= \vec{P} \cdot \frac{c(\vec{P} - \frac{e}{c} \vec{A})}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} + \frac{m^2 c^3}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} + e \Phi$$

$$- \frac{e}{c} \vec{A} \cdot \frac{c(\vec{P} - \frac{e}{c} \vec{A})}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} = c \frac{(\vec{P} - \frac{e}{c} \vec{A})^2 + m^2 c^2}{\sqrt{m^2 c^2 + (\vec{P} - \frac{e}{c} \vec{A})^2}} + e \Phi$$

$\Rightarrow$  the Hamiltonian is

$$H = \underbrace{\sqrt{m^2 c^4 + c^2 (\vec{P} - \frac{e}{c} \vec{A})^2}}_{\sqrt{m^2 c^4 + p^2 c^2}} + e \Phi \Rightarrow \text{total energy}$$

of the particle!

Non-relativistic limit can be obtained by expansion:

$$H \approx mc^2 + \frac{(\vec{P} - \frac{e}{c} \vec{A})^2}{2m} + e \Phi$$

$\sim$  may be familiar from QM

In QM quantize the system by replacing  $\vec{P} \rightarrow -i\hbar \vec{v}$ .