

Start out with microscopic equations for magnetostatics ($\vec{b} \sim$ magnetic field, $\vec{j} \sim$ current)

$$\boxed{\vec{\nabla} \cdot \vec{b} = 0} \quad \text{and} \quad \boxed{\vec{\nabla} \times \vec{b} = \mu_0 \vec{j}}$$

Introduce averaging in space:

$$\langle \vec{j}(\vec{x}) \rangle = \int \frac{d^3x'}{V} \vec{j}(\vec{x} - \vec{x}') f(\vec{x}')$$

where $f(\vec{x}')$ is some profile function: 

e.g. $f(\vec{x}) = \begin{cases} 1, & |\vec{x}| < R \\ 0, & \text{otherwise} \end{cases} ; V = \frac{4}{3}\pi R^3.$

$\Rightarrow$ Averaging represents "smearing" over some region of radius R !

$$\text{Now, } \vec{j}(\vec{x}) = \sum_n \sum_j q_i (\vec{V}_n + \vec{V}_{nj}) \delta(\vec{x} - \vec{x}_n - \vec{x}_{nj})$$

$\uparrow$ all molecules $\uparrow$ all electrons/protons in molecules

$\vec{x}_n \sim$ position of n th molecule

$\vec{V}_n \sim$ velocity ———

$\vec{x}_{nj} \sim$ position of charge i in the n th molecule with respect to its center

$\vec{V}_{nj} \sim$ velocity ———

$$\langle \vec{j}(\vec{x}) \rangle = \left\langle \sum_{n,j} q_j (\vec{v}_n + \vec{v}_{n,j}) \delta(\vec{x} - \vec{x}_n - \vec{x}_{n,j}) \right\rangle =$$

$$= \left\langle \sum_{n,j} q_j (\vec{v}_n + \vec{v}_{n,j}) \left[\delta(\vec{x} - \vec{x}_n) - \vec{x}_{n,j} \cdot \vec{\nabla} \delta(\vec{x} - \vec{x}_n) + \dots \right] \right\rangle$$

Def. Macroscopic current

$$\vec{j}(\vec{x}) \equiv \left\langle \sum_{n,j} q_j (\vec{v}_n + \vec{v}_{n,j}) \delta(\vec{x} - \vec{x}_n) \right\rangle$$

Def. Macroscopic magnetization

$$\vec{M}(\vec{x}) \equiv \left\langle \sum_n \vec{m}_n \delta(\vec{x} - \vec{x}_n) \right\rangle =$$

$$= \left\langle \sum_{n,j} \frac{1}{2} q_j (\vec{x}_{n,j} \times \vec{v}_{n,j}) \delta(\vec{x} - \vec{x}_n) \right\rangle$$

Remember: $\vec{\mu}(\vec{x}) = \frac{1}{2} \vec{x} \times \vec{j}(\vec{x})$ microscopically

$$\Rightarrow \text{macroscopically } \vec{M}(\vec{x}) = \langle \vec{\mu}(\vec{x}) \rangle$$

$$= \frac{1}{2} \langle \vec{x} \times \vec{j}(\vec{x}) \rangle = \frac{1}{2} \int \frac{d^3x'}{V} (\vec{x} - \vec{x}') \times \vec{j}(\vec{x} - \vec{x}') f(\vec{x}')$$

$$= \int \frac{d^3x'}{V} \frac{1}{2} \sum_{n,j} q_j (\underbrace{\vec{x} - \vec{x}'}_{-\vec{x}_n \sim \text{shift the origin}}) \times (\vec{v}_n + \vec{v}_{n,j}) \delta(\vec{x} - \vec{x}' - \vec{x}_n - \vec{x}_{n,j}) f(\vec{x}')$$

$$|\vec{v}_n| \ll |\vec{v}_{n,j}| \Rightarrow \text{neglect}$$

$$\Rightarrow \vec{M}(\vec{x}) = \frac{1}{2V} \sum_{n,j} g_j (\vec{x}_{nj} \times \vec{v}_{nj}) f(\vec{x} - \vec{x}_n - \vec{x}_{nj})$$

$$\approx \frac{1}{2} \frac{1}{V} \sum_{n,j} g_j (\vec{x}_{nj} \times \vec{v}_{nj}) f(\vec{x} - \vec{x}_n) + o(x_{nj})$$

$$\Rightarrow \vec{M}(\vec{x}) = \left\langle \frac{1}{2} \sum_{n,j} g_j (\vec{x}_{nj} \times \vec{v}_{nj}) \delta(\vec{x} - \vec{x}_n) \right\rangle$$

as defined $\Rightarrow$ it's consistent!

With the above definitions:

$$\boxed{\langle \vec{J}(\vec{x}) \rangle = \vec{J}(\vec{x}) + \vec{\nabla} \times \vec{M}(\vec{x})}$$

Check:

$$\begin{aligned} \vec{\nabla} \times \vec{M} &= \left\langle \frac{1}{2} \sum_{n,j} g_j \vec{\nabla} \times [(\vec{x}_{nj} \times \vec{v}_{nj}) \delta(\vec{x} - \vec{x}_n)] \right\rangle \\ &= \left\langle \sum_{n,j} \frac{1}{2} g_j \left[\vec{x}_{nj} (\vec{v}_{nj} \cdot \vec{\nabla}) \delta(\vec{x} - \vec{x}_n) - \vec{v}_{nj} (\vec{x}_{nj} \cdot \vec{\nabla}) \delta(\vec{x} - \vec{x}_n) \right] \right\rangle \\ &= \left(\text{similar to } \int d^3x [x_i J_i + x_j J_j] = 0 \right) \\ &= - \left\langle \sum_{n,j} g_j \vec{v}_{nj} (\vec{x}_{nj} \cdot \vec{\nabla}) \delta(\vec{x} - \vec{x}_n) \right\rangle \approx \left(\text{as } v_n \ll v_{nj} \right) \\ &\approx - \left\langle \sum_{n,j} g_j (\vec{v}_n + \vec{v}_{nj}) (\vec{x}_{nj} \cdot \vec{\nabla}) \delta(\vec{x} - \vec{x}_n) \right\rangle \end{aligned}$$

Finally, defining macroscopic magnetic field

$$\vec{B}(\vec{r}) \equiv \langle \vec{b}(\vec{r}) \rangle$$

$$\text{get: } \vec{\nabla} \cdot \vec{b} = 0 \Rightarrow \boxed{\vec{\nabla} \cdot \vec{B} = 0}$$

$$\vec{\nabla} \times \vec{b} = \mu_0 \vec{j} \Rightarrow \vec{\nabla} \times \vec{B} = \mu_0 \langle \vec{j} \rangle = \mu_0 [\vec{J} + \vec{\nabla} \times \vec{M}]$$

$$\Rightarrow \boxed{\vec{\nabla} \times \vec{B} = \mu_0 [\vec{J} + \vec{\nabla} \times \vec{M}]}$$

In principle one can define $\vec{M}$ by $\vec{J}_M = \vec{\nabla} \times \vec{M}$ with $\vec{J}_M$ effective current due to bound currents.

Then $\vec{M} \rightarrow \vec{M} + \vec{\nabla} \psi$ leaves $\vec{J}_M$ unchanged $\Rightarrow$ ambiguity in definition. $\vec{M} = \frac{1}{2} \vec{x} \times \vec{J}_M$ and $\vec{J} = \vec{\nabla} \times \vec{M}$ are related by "gauge" transformations:

$$\begin{aligned} m_i &= \frac{1}{2} \int d^3x (\vec{x} \times \vec{J}_M)_i = \frac{1}{2} \int d^3x [\vec{x} \times (\vec{\nabla} \times \vec{M})]_i = \\ &= \frac{1}{2} \int d^3x [x_j \nabla_i M_j - x_j \nabla_j M_i] \stackrel{(\text{parts})}{=} \\ &= \frac{1}{2} \int d^3x [-M_i + 3M_i] = \int d^3x M_i \end{aligned}$$

$$\Rightarrow \boxed{\int d^3x \frac{1}{2} \vec{x} \times \vec{J}_M = \int d^3x \vec{M}}$$