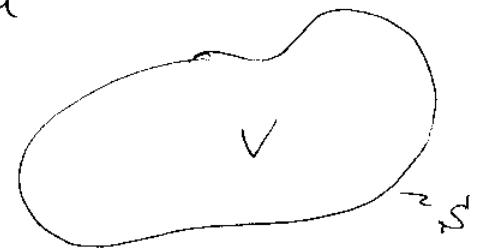


Green's Theorem.

Start from divergence theorem:

$$\int_V \vec{\nabla} \cdot \vec{A} d^3x = \oint_S \vec{A} \cdot \hat{n} da$$



Put $\vec{A} = \phi \vec{\nabla} \psi$, with ϕ, ψ

two arbitrary scalar fields:

$$\int_V \vec{\nabla} \cdot (\phi \vec{\nabla} \psi) d^3x = \oint_S \phi \hat{n} \cdot \vec{\nabla} \psi da$$

as $\vec{\nabla} \cdot (\phi \vec{\nabla} \psi) = \phi \nabla^2 \psi + (\vec{\nabla} \phi)(\vec{\nabla} \psi)$

and denoting $\hat{n} \cdot \vec{\nabla} \psi = \frac{\partial \psi}{\partial n}$ we get

$$\int_V d^3x [\phi \nabla^2 \psi + (\vec{\nabla} \phi)(\vec{\nabla} \psi)] = \oint_S \phi \frac{\partial \psi}{\partial n} da$$

Green's first identity.

Swap $\phi \leftrightarrow \psi$:

$$\int_V d^3x [\psi \nabla^2 \phi + (\vec{\nabla} \psi)(\vec{\nabla} \phi)] = \oint_S \psi \frac{\partial \phi}{\partial n} da$$

& subtract from the 1st identity:

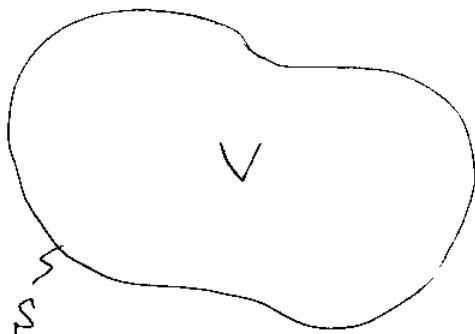
$$\int_V d^3x [\phi \nabla^2 \psi - \psi \nabla^2 \phi] = \oint_S \left[\phi \frac{\partial \psi}{\partial n} - \psi \frac{\partial \phi}{\partial n} \right] da$$

Green's second identity or Green's theorem.

Solution of Poisson Equation:

Dirichlet & Neumann Boundary Conditions

$$\nabla^2 \phi = -\frac{\rho}{\epsilon_0} \quad \text{in volume } V.$$



① ϕ is specified on S

~ Dirichlet boundary condition

② $\frac{\partial \phi}{\partial n}$ is specified on S

~ Neumann boundary condition

Uniqueness of the solution:

Suppose there are 2 solutions $\sim \phi_1$ & ϕ_2

$$\nabla^2 \phi_1 = -\frac{\rho}{\epsilon_0} \quad \& \quad \nabla^2 \phi_2 = -\frac{\rho}{\epsilon_0} \Rightarrow \text{define}$$

$u = \phi_1 - \phi_2 \Rightarrow \nabla^2 u = 0 \Rightarrow$ put $\phi = \psi = u$ in
first Green's identity

$$\int_V d^3x \left[u \underbrace{\nabla^2 u}_{=0} + |\vec{\nabla} u|^2 \right] = \oint_S u \frac{\partial u}{\partial n} da$$

$$\Rightarrow \int_V d^3x |\vec{\nabla} u|^2 = 0 \quad \left. \begin{array}{l} \Rightarrow u = 0 \\ \text{in } V \Rightarrow \\ \text{for Neumann } \frac{\partial u}{\partial n} = 0 \text{ on } S' \end{array} \right\} \Rightarrow \text{solution is unique, } \phi_1 = \phi_2.$$

(in case of Neumann one may have $u = \text{const}$
 ~ not important, ϕ is defined up to a constant
 anyway)

Green Functions (Green had no formal math education when he published it all in 1828 at the age of *35)

Suppose you have a linear differential equation

$$\hat{L}_x \psi(x) = J(x)$$

where $J(x)$ is known, \hat{L}_x is some differential operator and $\psi(x)$ is to be found.

If we know the Green function of operator \hat{L}_x defined by $\hat{L}_x G(\vec{x}, \vec{x}') = \delta^3(\vec{x} - \vec{x}')$, then

$\psi(x) = \int d^3x' J(\vec{x}') \cdot G(\vec{x}, \vec{x}')$ would be the solution of our equation. Works for any linear operator \hat{L}_x , and any "good" function $J(x)$.

In our case, define Green function by

$$\left\{ \nabla^2 G(\vec{x}, \vec{x}') = -4\pi \delta^3(\vec{x} - \vec{x}') \right\}$$

We know that $G(\vec{x}, \vec{x}') = \frac{1}{|\vec{x} - \vec{x}'|} + F(\vec{x}, \vec{x}')$

for any F such that $\nabla^2 F(\vec{x}, \vec{x}') = 0$ in V .

Substitute $\phi = \phi$ the potential and

$\psi = G(\vec{x}, \vec{x}')$ into the second Green's identity:

$$\int_V d^3x' \left[\underbrace{\phi \nabla'^2 G(\vec{x}, \vec{x}')}_{-4\pi \delta^3(\vec{x} - \vec{x}')} - G(\vec{x}, \vec{x}') \underbrace{\nabla'^2 \phi}_{-\rho/\epsilon_0} \right] =$$

$$= \oint_S \left[\phi(\vec{x}') \frac{\partial G(\vec{x}, \vec{x}')}{\partial n'} - G(\vec{x}, \vec{x}') \frac{\partial \phi}{\partial n'} \right] da'$$

$$\boxed{\phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' G(\vec{x}, \vec{x}') \rho(\vec{x}') +}$$

"Master formula"

$$\boxed{+ \frac{1}{4\pi} \oint_S \left[G(\vec{x}, \vec{x}') \frac{\partial \phi}{\partial n'} - \phi(\vec{x}') \frac{\partial G(\vec{x}, \vec{x}')}{\partial n'} \right] da'}$$

Use the freedom of redefining $G \rightarrow G + F$, where (21)

$\nabla^2 F = 0$, to fix boundary conditions for $G(\vec{x}, \vec{x}')$.

Example: conductors are equipotential

(if not \Rightarrow get $\vec{E} \neq 0 \Rightarrow$ will become equipotential)

\Rightarrow natural candidate for Dirichlet boundary

conditions \sim conducting surfaces as boundaries

\vec{x}'' \vec{x}' \Rightarrow interested in potential outside
 \ominus - \oplus conductor

$$G = \frac{1}{|\vec{x} - \vec{x}'|} \quad \text{outside}$$

with $(-x'_1, y'_1, z'_1) = \vec{x}''$

Can add $F = -\frac{1}{|\vec{x} - \vec{x}''|}$ \checkmark as

$$\nabla^2 F = 0 \quad \text{in the } \underline{\text{volume}} \\ \underline{\text{of interest}}$$

One gets $G' = \frac{1}{|\vec{x} - \vec{x}'|} - \frac{1}{|\vec{x} - \vec{x}''|} \Rightarrow G' = 0$ on the surface

① To solve Dirichlet b.c. problem choose

$G_D(\vec{x}, \vec{x}') = 0$ for \vec{x}' on $S \Rightarrow$ using master formula

$$\phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' G_D(\vec{x}, \vec{x}') \rho(\vec{x}') -$$

$$- \frac{1}{4\pi} \oint_S \phi(\vec{x}') \frac{\partial G_D(\vec{x}, \vec{x}')}{\partial n} d\alpha'$$