

\Rightarrow choose $\vec{n} = (n_x, 0, n_z) \Rightarrow \vec{n} \cdot \vec{x} \Big|_{z=0} = n_x \cdot x \Rightarrow$ no y -dep. (27)

\Rightarrow there should be no y -dependence in $\vec{k}' \cdot \vec{x}$ and in $\vec{k}'' \cdot \vec{x}$ as well $\Rightarrow k'_y = k''_y = 0 \Rightarrow$ all lie in the same plane.

$$k \cdot \sin \theta = k' \cdot \sin \theta' = k'' \cdot \sin \theta''$$

\Rightarrow as $k = k'' \Rightarrow \theta = \theta''$ \approx angle of reflection is equal to angle of incidence!

$$\text{as } k = \sqrt{\mu \epsilon} \omega \text{ and } k' = \sqrt{\mu' \epsilon'} \omega \Rightarrow$$

$$\sqrt{\mu \epsilon} \sin \theta = \sqrt{\mu' \epsilon'} \sin \theta'. \text{ Remember } n = (\sqrt{\mu \epsilon})$$

(index of refraction) $\Rightarrow n \sin \theta = n' \sin \theta'$

Snell's law!

The only thing left is to find \vec{E}_o' & \vec{E}_o'' using b.c.'s:

$$\vec{D}_n \text{ continuous} \Rightarrow \hat{n} \cdot [\epsilon(\vec{E}_o + \vec{E}_o'') - \epsilon' \vec{E}_o'] = 0$$

$$\vec{B}_n \text{ continuous} \Rightarrow \hat{n} \cdot [\vec{k} \times \vec{E}_o + \vec{k}'' \times \vec{E}_o'' - \vec{k}' \times \vec{E}_o'] = 0$$

(and $\omega = \omega' = \omega''$)

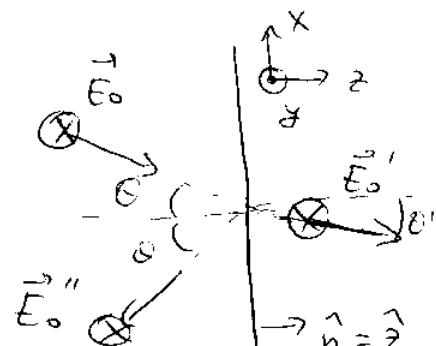
$$\vec{E}_t \text{ continuous} \Rightarrow \hat{n} \times [\vec{E}_o + \vec{E}_o'' - \vec{E}_o'] = 0$$

$$\vec{H}_t \text{ continuous} \Rightarrow \left[\frac{1}{\mu} (\vec{k} \times \vec{E}_o + \vec{k}'' \times \vec{E}_o'') - \frac{1}{\mu'} (\vec{k}' \times \vec{E}_o') \right] \times \hat{n} = 0$$

Consider 2 cases:

① $\vec{E}_o \perp \text{to the plane of incidence}$

$$\vec{E}_o, \vec{E}_o', \vec{E}_o'' \parallel \hat{y}$$



3rd & 4th equations: ($\hat{u} = \hat{z}$)

$$\left\{ \begin{array}{l} E_0 + E_0'' - E_0' = 0 \\ \frac{1}{\mu} (k E_0 \cos \theta - k'' E_0'' \cos \theta') - \frac{1}{\mu'} k' E_0' \cos \theta' = 0 \end{array} \right.$$

as $k = k'' = \sqrt{\mu \epsilon} \omega$, $k' = \sqrt{\mu' \epsilon'} \omega \Rightarrow \text{and } \theta = \theta''$

$$\left\{ \begin{array}{l} E_0 + E_0'' - E_0' = 0 \\ \sqrt{\frac{\epsilon}{\mu}} (E_0 - E_0'') \cos \theta - \sqrt{\frac{\epsilon'}{\mu'}} \cdot E_0' \cos \theta' = 0 \end{array} \right.$$

1st eqn. $0 = 0$; 2nd eqn.: $k E_0 \sin \theta + k'' E_0'' \sin \theta'' - k' E_0' \sin \theta' = 0 \Rightarrow (E_0 + E_0'') \sin \theta - \sqrt{\frac{\mu' \epsilon'}{\mu \epsilon}} E_0' \sin \theta' = 0$

as $\sqrt{\mu \epsilon} \sin \theta = \sqrt{\mu' \epsilon'} \sin \theta'$ (Snell's law) $\Rightarrow E_0 + E_0'' - E_0' = 0$
 \Rightarrow duplicates the 3rd one.

Using Snell's law ($n \sin \theta = n' \sin \theta'$) to get rid of n' we write (work it out yourself):

$$\boxed{\frac{E_0'}{E_0} = \frac{2n \cos \theta}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n^2 - n'^2 \sin^2 \theta}}$$

$$\frac{E_0''}{E_0} = \frac{n \cos \theta - \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

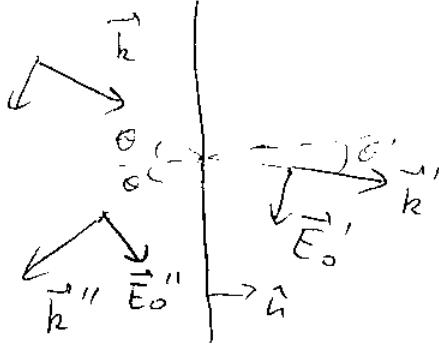
$$n = \sqrt{\frac{\mu \epsilon}{\mu_0 \epsilon_0}}$$

$$n' = \sqrt{\frac{\mu' \epsilon'}{\mu_0 \epsilon_0}}$$

② $\vec{E}_0 \parallel$ plane of incidence ($x-z$ plane)

2 independent equations (3rd & 4th):

$$\begin{cases} (E_0 + E_0'') \cos \theta - E_0' \cos \theta' = 0 \\ \sqrt{\frac{\epsilon}{\mu}} (E_0 + E_0'') - \sqrt{\frac{\epsilon'}{\mu'}} E_0' = 0 \end{cases}$$



(other two can be reduced to those)

Solve:

using
Snell's
Law

$$\frac{E_0'}{E_0} = \frac{2n n' \cos \theta}{\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

$$\frac{E_0''}{E_0} = \frac{-\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}{\frac{\mu}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

Normal incidence: $\theta = 0 \Rightarrow$ both ① and ②

give

$$\frac{E_0'}{E_0} = \frac{2n}{n + \frac{\mu}{\mu'} n'}$$

$$\frac{E_0''}{E_0} = \frac{n - \frac{\mu}{\mu'} n'}{n + \frac{\mu}{\mu'} n'}$$

Polarization by reflection: put $\mu = \mu'$ for simplicity

$$\textcircled{I}: \frac{E_0''}{E_0} = \frac{n \cos \theta - \sqrt{n'^2 - n^2 \sin^2 \theta}}{n \cos \theta + \sqrt{n'^2 - n^2 \sin^2 \theta}} \leftarrow \text{different} \Rightarrow$$

$$\textcircled{II}: \frac{E_0''}{E_0} = \frac{-n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}{n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}} \quad \checkmark \Rightarrow \text{reflected light is polarized!}$$

in case I $\frac{E_0''}{E_0}$ never vanishes (always < 0) (30)

in case II $\frac{E_0''}{E_0} = 0$ for $\boxed{\theta_B = \tan^{-1}\left(\frac{n'}{n}\right)}$ Brewster's angle

~~explosive~~

\Rightarrow reflected light is polarized.

if $\theta = \theta_B \Rightarrow$ polarization is linear, \perp to the plane of incidence.

(fish in the ocean reflect light \sim squids with polarized vision can see them)

Total internal reflection: Snell's law:

$$n \sin \theta = n' \sin \theta' \leq n' \Rightarrow \theta \leq \sin^{-1}\left(\frac{n'}{n}\right) \Rightarrow$$

\Rightarrow for $\boxed{\theta > \sin^{-1}\left(\frac{n'}{n}\right)}$ get total reflection
"evanescent wave"

$$k' = \sqrt{\mu' \epsilon'} \omega \Rightarrow k'_x = -\sqrt{\mu' \epsilon'} \omega \sin \theta' = -\sqrt{\mu_0 \epsilon_0} \omega n' \sin \theta' =$$

$$n' \frac{\omega}{c} \quad k'_y = 0 \quad = -\frac{\omega}{c} n \sin \theta$$

$$k'_z = \sqrt{k'^2 - k'_x^2} = \sqrt{n'^2 - n^2 \sin^2 \theta} = \frac{\omega}{c}$$

\Rightarrow for $\theta > \sin^{-1}\left(\frac{n'}{n}\right)$: k'_z becomes imaginary

$\Rightarrow e^{i k_z z} \sim e^{-|k_z| z}$ \sim exponential fall-off

effectively ^{the wave} gets reflected from a different surface \sim violation of geom. optics, Goos-Hänchen effect shift.