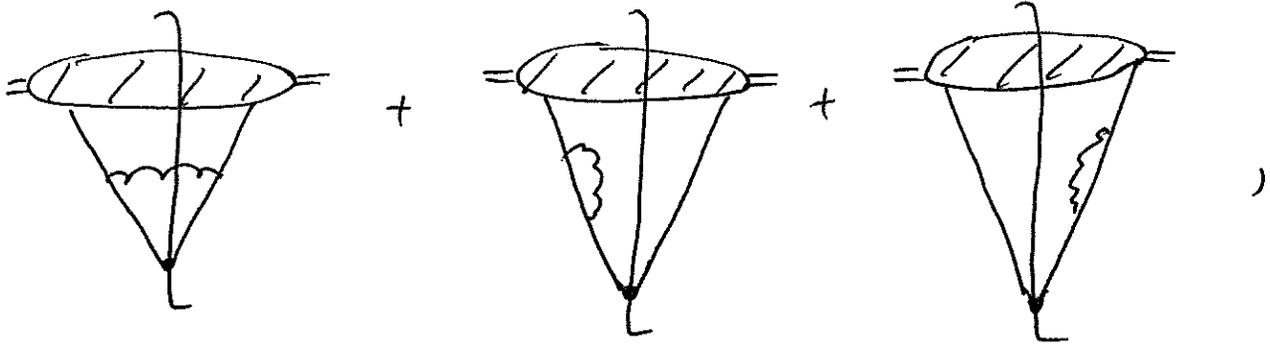


Last time

We completed the calculation of the 1-loop corrections to quark PDF,



and obtained an equation,

$$\frac{\partial}{\partial \ln Q^2} q^f(x, Q^2) = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dx_1}{x_1} P_{qq} \left( \frac{x_1}{x} \right) q^f(x_1, Q^2).$$

Here  $P_{qq}(z) = C_F \left( \frac{1+z^2}{1-z} \right)_+$  is the quark-quark splitting function.

$$\int_0^1 dz [h(z)]_+ f(z) = \int_0^1 dz h(z) [f(z) - f(1)]$$

"plus notation"



where  $x = \frac{k^+}{p^+}$ ,  $x_1 = \frac{k_1^+}{p^+} \Rightarrow z = \frac{k^+}{k_1^+} = \frac{x}{x_1}$

as  $z < 1 \Rightarrow x_1 > x$  in the integral.

Including the virtual terms (B and C)

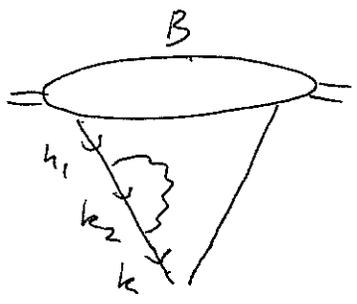
gives

$$P_{gg}(z) = C_F \left( \frac{1+z^2}{1-z} \right)_+ \sim \text{splitting function}$$

where

$$\int_0^1 dz [h(z)]_+ f(z) = \int_0^1 dz h(z) [f(z) - f(1)]$$

Easy to understand:



$$\propto \delta(k-k_1)$$

$$\Downarrow$$

$$x = x_1$$

$$\Rightarrow g^f(x_1, Q^2) = g^f(x, Q^2)$$

as  $x = x_1$

"real" part,  
 $\downarrow$  diagram A

$$Q^2 \frac{\partial}{\partial Q^2} g^f(x, Q^2) = \frac{\alpha C_F}{2\pi} \left[ \int_x^1 \frac{dz}{z} \frac{1+z^2}{1-z} \cdot g^f\left(\frac{x}{z}, Q^2\right) - \int_0^1 dz \frac{1+z^2}{1-z} g^f(x, Q^2) \right]$$

Virtual corrections, graphs B & C

$z \rightarrow 1$  divergence is cancelled between the real (A) and virtual (B+C) terms.

bare quark state  $|\psi_0\rangle = \text{---} \Rightarrow \langle \psi_0 | \psi_0 \rangle = 1$  (102)

(normalization)

dressed quark state  $|\psi\rangle = \underbrace{\text{---}}_{|\psi_0\rangle} + \underbrace{\text{---}}_{|\psi_1\rangle} + \underbrace{\text{---}}_{V|\psi_0\rangle}$

normalization:

$$\langle \psi | \psi \rangle = 1 = \langle \psi_0 | \psi_0 \rangle + \underbrace{\text{---}} + \underbrace{\text{---}} + \underbrace{\text{---}}$$

$$= 1 + \langle \psi_1 | \psi_1 \rangle + 2V \langle \psi_0 | \psi_0 \rangle = 1 + \langle \psi_1 | \psi_1 \rangle + 2V$$

$$\Rightarrow \boxed{V = -\frac{1}{2} \langle \psi_1 | \psi_1 \rangle}$$

$$\Rightarrow \text{graphs } B, C = -\frac{1}{2} A \Rightarrow \boxed{B + C = -A}$$

$\approx$  simply imposed probability conservation!

E

$$\frac{\partial}{\partial h Q^2} g^f(x, Q^2) = \frac{\alpha_C F}{2\pi} \int_x^1 \frac{dx_1}{x_1} \left( \frac{1 + \left(\frac{x}{x_1}\right)^2}{1 - \frac{x}{x_1}} \right)_+ g^f(x_1, Q^2) =$$

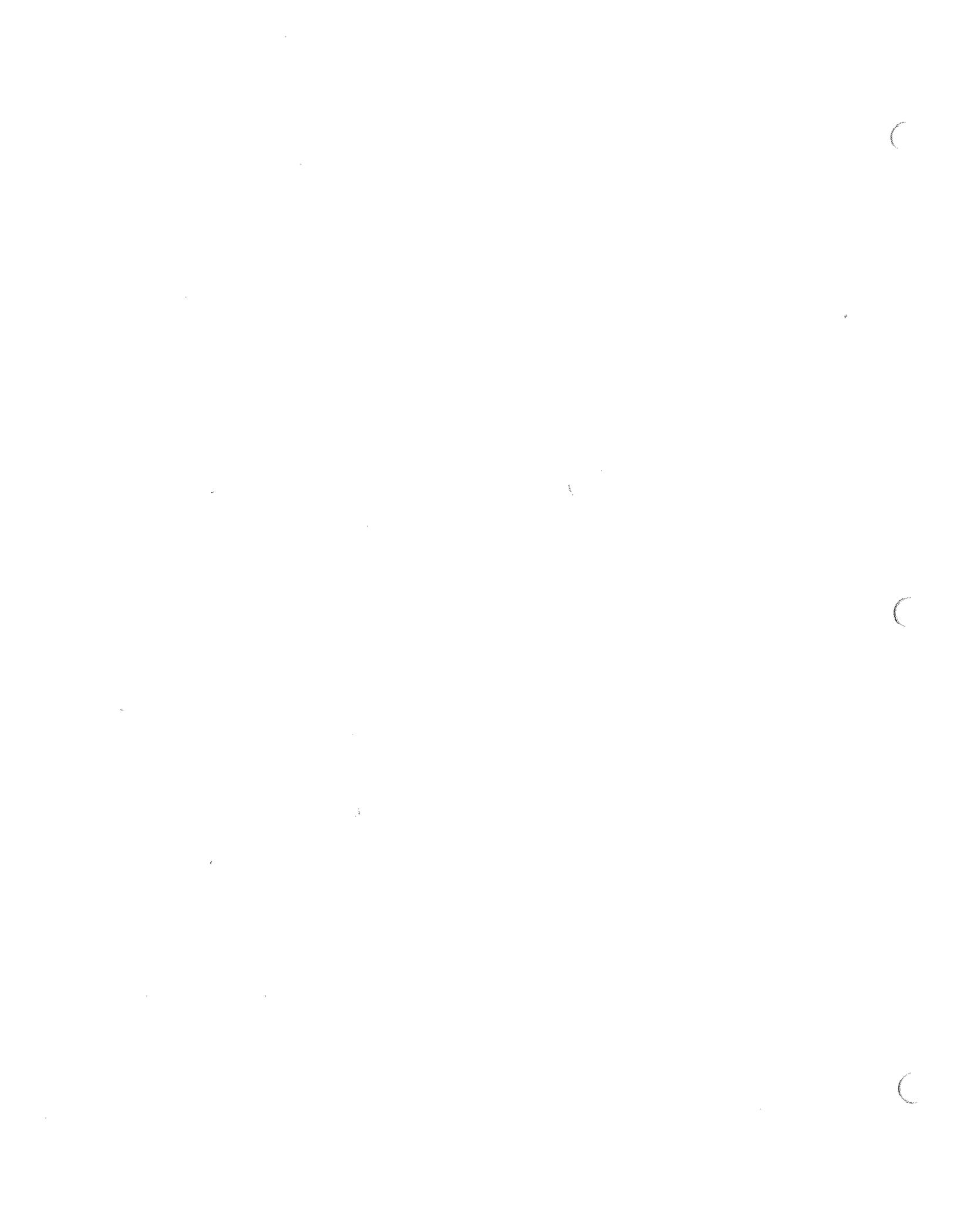
$$= \left\{ \begin{array}{l} z = \frac{x}{x_1} \\ \frac{dx_1}{x_1} = -\frac{dz}{z} \end{array} \right. = \frac{\alpha_C F}{2\pi} \int_x^1 \frac{dz}{z} \left( \frac{1+z^2}{1-z} \right)_+ g^f\left(\frac{x}{z}, Q^2\right)$$

$$= \frac{\alpha_C F}{2\pi} \int_0^1 dz \underbrace{\left( \frac{1+z^2}{1-z} \right)_+}_{h(z)} \underbrace{\frac{1}{z} g^f\left(\frac{x}{z}, Q^2\right) \Theta(z-x)}_{f(z)}$$

$$= \left\{ \begin{array}{l} \text{using} \\ \int_0^1 dz [h(z)]_+ f(z) = \int_0^1 dz h(z) [f(z) - f(1)] \end{array} \right.$$

$$= \frac{\alpha_C F}{2\pi} \int_0^1 dz \frac{1+z^2}{1-z} \left[ \frac{1}{z} g^f\left(\frac{x}{z}, Q^2\right) \Theta(z-x) - g^f(x) \Theta(1-x) \right]$$

$$= \frac{\alpha_S C_F}{2\pi} \left[ \int_x^1 \frac{dz}{z} \frac{1+z^2}{1-z} g^f\left(\frac{x}{z}, Q^2\right) - \int_0^1 dz \frac{1+z^2}{1-z} g^f(x) \right]$$



Def. Defining flavor singlet distribution

$$\Sigma(x, Q^2) \equiv \sum_f [q_f^+(x, Q^2) + q_f^-(x, Q^2)]$$

Def. and flavor non-singlet

$$\Delta^{f\bar{f}}(x, Q^2) \equiv q_f^+(x, Q^2) - q_{\bar{f}}^-(x, Q^2)$$

we write

$$Q^2 \frac{\partial}{\partial Q^2} \Delta^{f\bar{f}}(x, Q^2) = \frac{\alpha(Q^2)}{2\pi} \int_x^1 \frac{dx_1}{x_1} P_{gg}\left(\frac{x}{x_1}\right) \cdot \Delta^{f\bar{f}}(x_1, Q^2)$$

and

$$Q^2 \frac{\partial}{\partial Q^2} \begin{pmatrix} \Sigma(x, Q^2) \\ G(x, Q^2) \end{pmatrix} = \frac{\alpha(Q^2)}{2\pi} \int_x^1 \frac{dx_1}{x_1} \begin{pmatrix} P_{gg}\left(\frac{x}{x_1}\right) & P_{gq}\left(\frac{x}{x_1}\right) \\ P_{qg}\left(\frac{x}{x_1}\right) & P_{qq}\left(\frac{x}{x_1}\right) \end{pmatrix} \cdot \begin{pmatrix} \Sigma(x_1, Q^2) \\ G(x_1, Q^2) \end{pmatrix}$$

Dokshitzer, Gribov, Lipatov, Altarelli, Parisi

(DGLAP) Equations

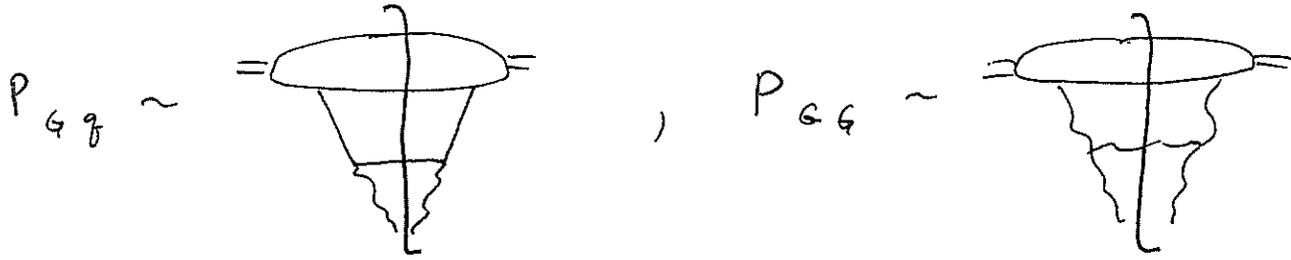
GL ~ QED case ~ '72

D, A & P ~ QCD case, '77

Def.

$$G(x, Q^2) = \text{[Diagram of a gluon distribution function: a circle with a horizontal line through its center and a jagged, downward-pointing shape below it]} \sim \langle A_i A_i \rangle \text{ in } A_+ = 0 \text{ gauge}$$

gluon distribution function



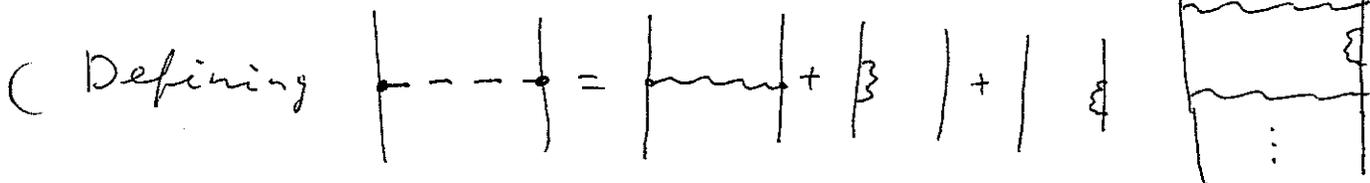
After explicit calculations one gets the splitting functions:

$$\left\{ \begin{aligned} P_{gg}(z) &= C_F \left( \frac{1+z^2}{1-z} \right)_+ \\ P_{gq}(z) &= C_F \frac{1+(1-z)^2}{z} \\ P_{qg}(z) &= N_F [z^2 + (1-z)^2] \\ P_{qq}(z) &= 2N_C \left[ \frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z) \right] + \frac{11N_C - 2N_F}{6} \delta(z-1) \end{aligned} \right.$$

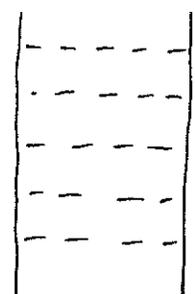
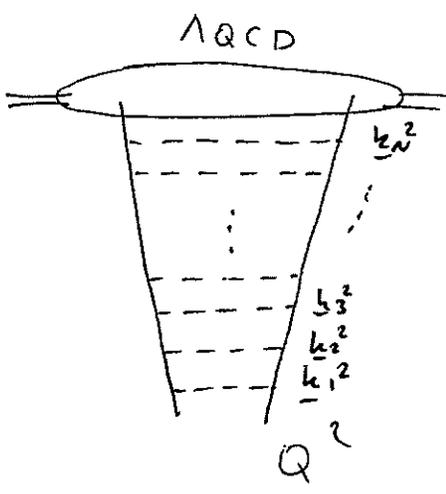
Note that  $P_{gq}(z)$  can be obtained from

$P_{qg}(z)$  by substituting  $z \rightarrow 1-z$  and dropping virtual corrections.

Iterate the evolution for  $f_i(x, Q^2)$ :



We get a ladder diagram:

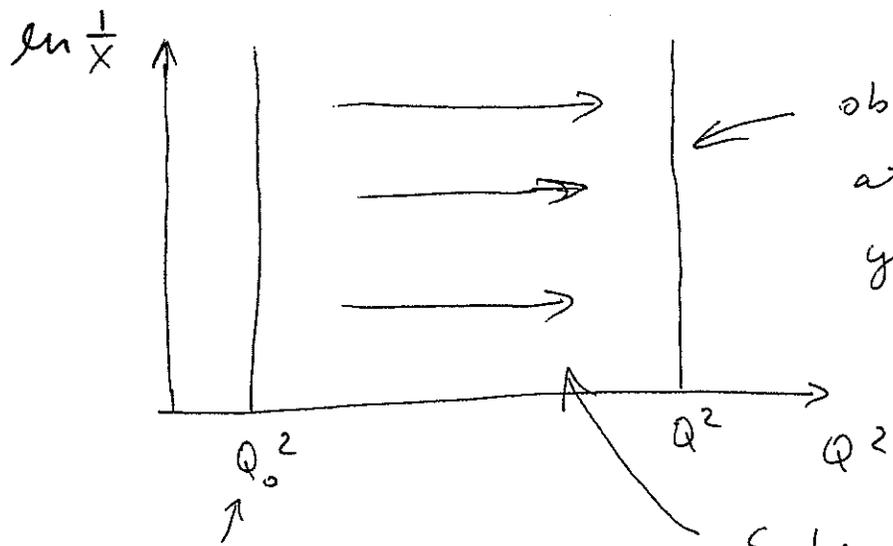


$$Q^2 \gg k_1^2 \gg k_2^2 \gg \dots \gg k_N^2 \gg \Lambda_{QCD}^2$$

DGLAP resums ladder graphs with the ladder connecting scales  $Q^2$  and  $\Lambda_{QCD}^2$ ,  $Q^2 \gg \Lambda_{QCD}^2$

such that  $\ln Q^2 / \Lambda_{QCD}^2 \gg 1$  and  $d_s \ln \frac{Q^2}{\Lambda_{QCD}^2} \sim 1$  is the resummation parameter.

How does DGLAP work?



obtain distribution at whatever  $Q^2$  you wanted.

start with some initial condition

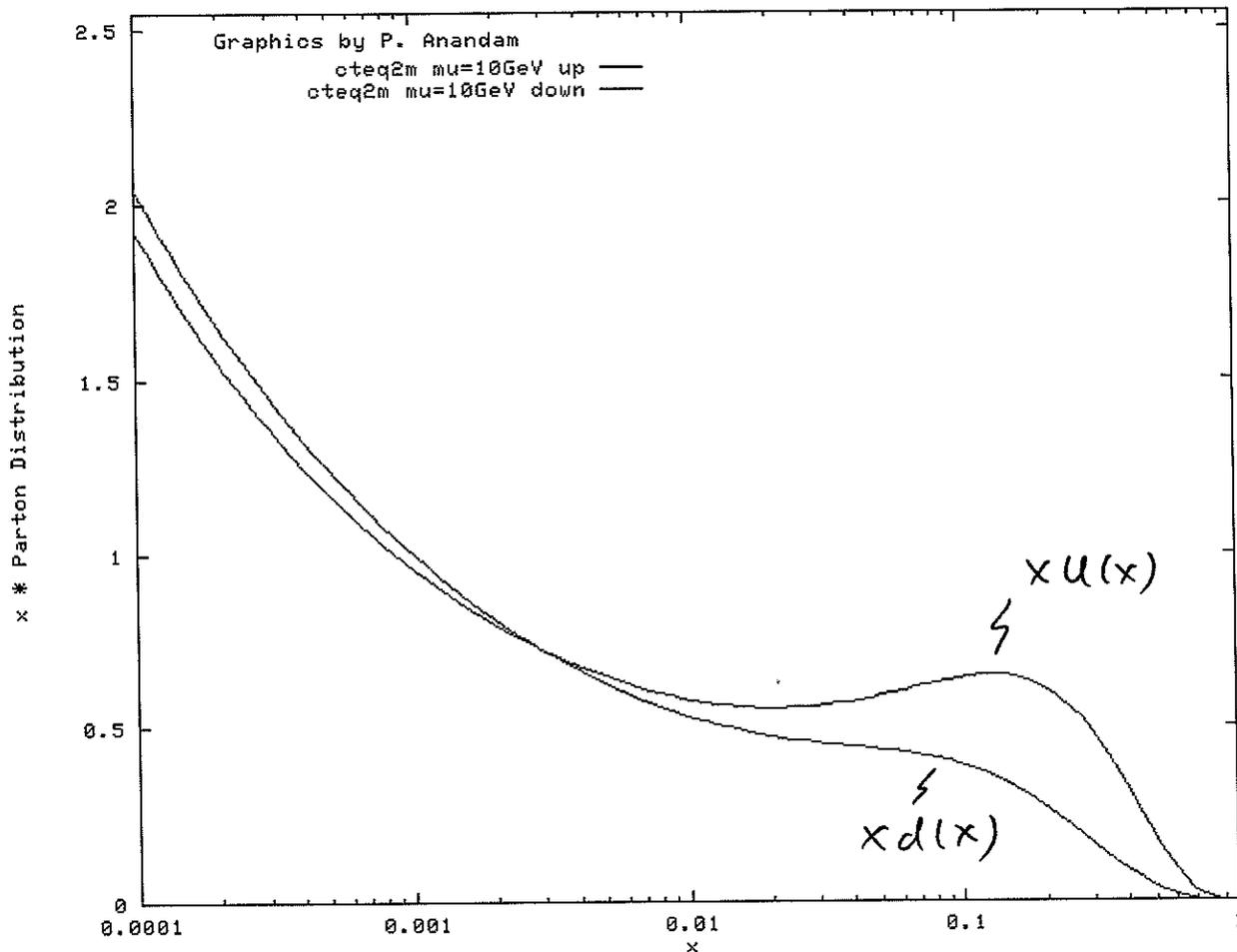
Solve the DGLAP equations ("evolve" the distribution function)

function  $q^f(x, Q_0^2)$

=> people calculate PDF's (Parton Distribution Functions) & fit the data. See attachments for PDF examples.

# Parton Distribution Graph

(Number of graphs plotted since 21 November 2000: 658)



$$Q = 10 \text{ GeV} \Rightarrow Q^2 = 100 \text{ GeV}^2$$

at large - x valence quarks dominate

$$\Rightarrow x u_v(x) = 2 x d_v(x)$$

$\Rightarrow$  not so at small - x

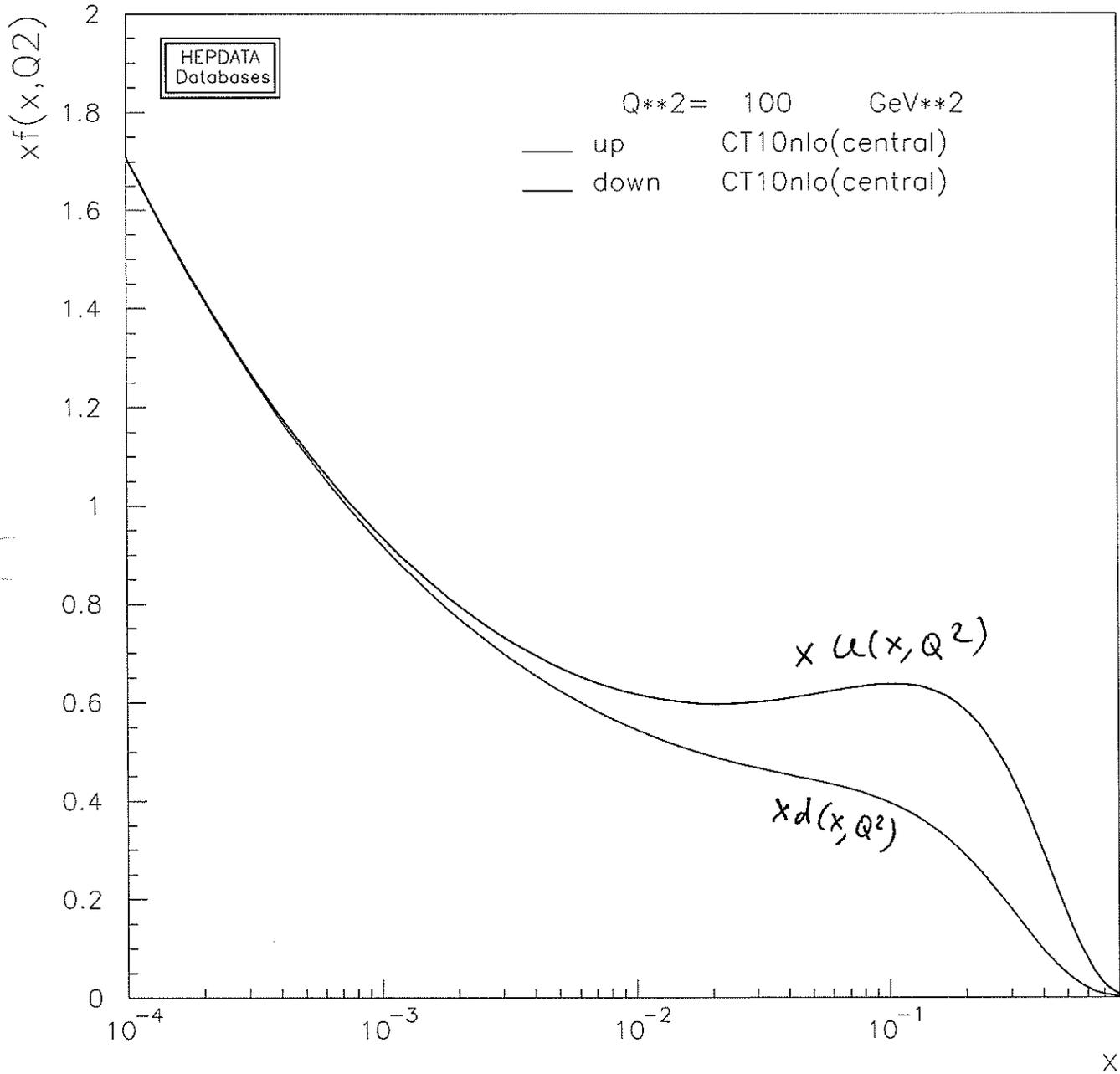
outdated  
site

go to <http://zebu.uoregon.edu/~parton/partongraf.html>

to plot more.

hepdata.cedar.ac.uk/pdf/pdf3.html  $\leftarrow$  can plot more if you wish

$$Q^2 = 100 \text{ GeV}^2 \quad (\Rightarrow \quad Q = 10 \text{ GeV})$$



$\sim$  at large  $x$  valence quarks dominate

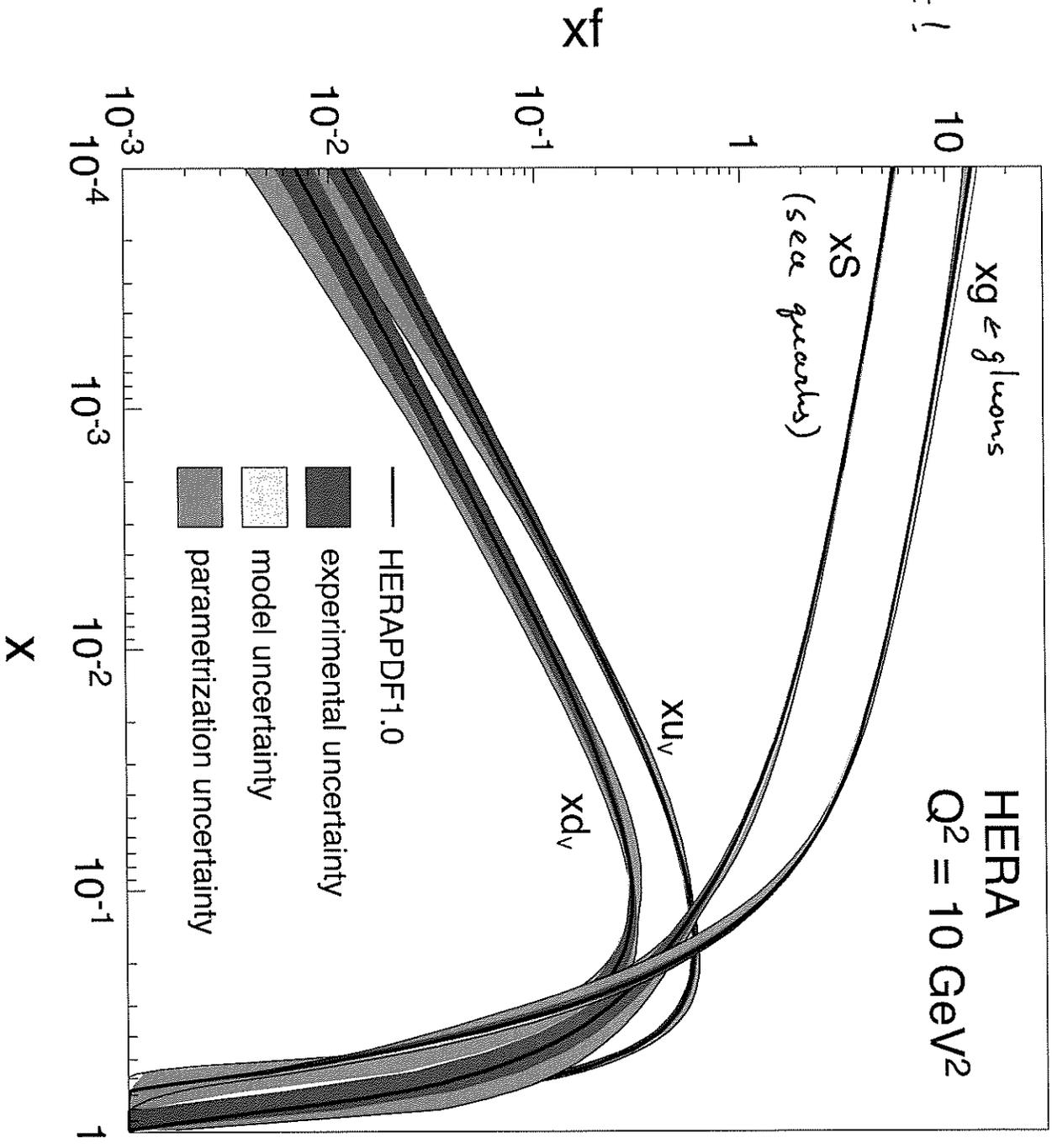
$$x u_v \approx 2 x d_v$$

$\sim$  at small  $x$  sea quarks dominate



Note: this is  
a log-log plot!

↙ at small  $x$ , gluons and sea quarks dominate!



C

C

C