

Last time

Axial Anomaly (cont'd)

massless QED : $\mathcal{L} = \bar{\psi} i \not{\partial} \psi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$

$U(1)$: $\psi(x) \rightarrow e^{i\alpha} \psi(x)$
global

$U(1)_A$: $\psi(x) \rightarrow e^{i\gamma_5 \alpha} \psi(x)$
global

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

$$j^\mu_5 = \bar{\psi} \gamma^\mu \gamma^5 \psi$$

$$\partial_\mu j^\mu = 0$$

$$\partial_\mu j^\mu_5 = 0 \quad (m=0)$$

current conservation is true classically.

Is it true when quantum loops are included?

$$T_{\mu\nu\rho} (k_1, k_2) \equiv i \int d^4x_1 d^4x_2 e^{ik_1 \cdot x_1 + ik_2 \cdot x_2}$$

$$\cdot \langle 0 | T [j^\mu(x_1) j^\nu(x_2) j^\rho_5(0)] | 0 \rangle$$

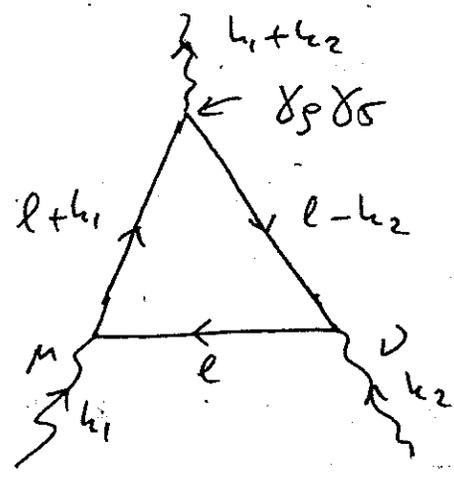
if $\partial_\mu j^\mu_5 = 0 \Rightarrow (k_1 + k_2)^\rho T_{\mu\nu\rho} (k_1, k_2) = 0$.

(can prove)

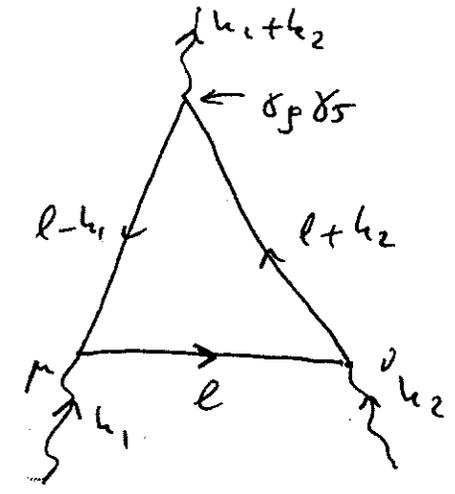
Check this statement:

$T_{\mu\nu\rho}(k_1, k_2) =$

arrow indicates both momentum & fermion #.



graph A



graph B

(Can write $T_{\mu\nu\rho}(k_1, k_2) = i \int d^4x_1 d^4x_2 e^{-ik_1 \cdot x_1 + ik_2 \cdot (x_2 - x_1)}$

$\langle 0 | T [j_\mu(0) j_\nu(x_2) j_\rho^S(x_1)] | 0 \rangle \Rightarrow (k_1 + k_2)^\rho T_{\mu\nu\rho} =$

$= i \int d^4x_1 d^4x_2 i \partial_{x_1}^\rho \left(e^{ik_2 \cdot x_2 - i(k_1 + k_2) \cdot x_1} \langle 0 | T [j_\mu(0) j_\nu(x_2) j_\rho^S(x_1)] | 0 \rangle \right)$

$= (\text{parts}) = \int d^4x_1 d^4x_2 e^{ik_2 \cdot x_2 - i(k_1 + k_2) \cdot x_1} \langle 0 | T [j_\mu(0) j_\nu(x_2) \cdot$

+ more work (equal-time commut. relations, etc) ~ see attached pages

$\cdot \partial^\rho j_\rho^S(x_1)] | 0 \rangle = 0 \quad \text{if} \quad \partial^\rho j_\rho^S = 0$

fermion loop, $m=0$

$-i T_{\mu\nu\rho} = -(-ie) \int \frac{d^4l}{(2\pi)^4} \text{Tr} \left[\delta_\rho \delta_\sigma \frac{i}{l+k_1} \delta_\mu \frac{i}{l} \delta_\nu \frac{i}{l-k_2} \right]$

We put $j_\rho = -ie \bar{\psi} \gamma_\rho \psi$ to have a real QED vertex.

$-(-ie)^2 \int \frac{d^4l}{(2\pi)^4} \text{Tr} \left[\delta_\rho \delta_\sigma \frac{i}{l+k_2} \delta_\nu \frac{i}{l} \delta_\mu \frac{i}{l-k_1} \right] =$

$= -ie^2 \int \frac{d^4l}{(2\pi)^4} \frac{\text{Tr} [\delta_\rho \delta_\sigma (l+k_1) \delta_\mu l \delta_\nu (l-k_2)]}{(l^2 + i\epsilon)((l+k_1)^2 + i\epsilon)((l-k_2)^2 + i\epsilon)}$

$$-ie^2 \int \frac{d^4 \ell}{(2\pi)^4} \frac{\text{Tr} [\gamma_\rho \gamma_5 (\ell + k_2) \gamma_\nu \not{\ell} \gamma_\mu (\ell - k_1)]}{(\ell^2 + i\varepsilon)((\ell + k_2)^2 + i\varepsilon)((\ell - k_1)^2 + i\varepsilon)} \quad (118)$$

$$\Rightarrow (k_1 + k_2)^\rho T_{\mu\nu\rho} = e^2 \int \frac{d^4 \ell}{(2\pi)^4} \left\{ \frac{\text{Tr} [(k_1 + k_2) \gamma_5 (\ell + k_1) \gamma_\mu \not{\ell} \gamma_\nu (\ell - k_2)]}{(\ell^2 + i\varepsilon)((\ell + k_1)^2 + i\varepsilon)((\ell - k_2)^2 + i\varepsilon)} \right. \\ \left. + \frac{\text{Tr} [(k_1 + k_2) \gamma_5 (\ell + k_2) \gamma_\nu \not{\ell} \gamma_\mu (\ell - k_1)]}{(\ell^2 + i\varepsilon)((\ell + k_2)^2 + i\varepsilon)((\ell - k_1)^2 + i\varepsilon)} \right\} \quad \begin{matrix} \text{"A"} \\ \text{"B"} \end{matrix}$$

$$\text{Numerator of A} = \text{Tr} [(k_1 + \not{\ell} - (\not{\ell} - k_2)) \gamma_5 (\ell + k_1) \gamma_\mu \not{\ell} \gamma_\nu (\ell - k_2)] \\ = -(\ell + k_1)^2 \text{Tr} [\gamma_5 \gamma_\mu \not{\ell} \gamma_\nu (\ell - k_2)] - \\ - (\ell - k_2)^2 \text{Tr} [\gamma_5 (\ell + k_1) \gamma_\mu \not{\ell} \gamma_\nu]$$

$$\text{Numerator of B} = \text{Tr} [((k_2 + \not{\ell}) - (\not{\ell} - k_1)) \gamma_5 (\ell + k_2) \gamma_\nu \not{\ell} \gamma_\mu (\ell - k_1)] \\ = -(\ell + k_2)^2 \text{Tr} [\gamma_5 \gamma_\nu \not{\ell} \gamma_\mu (\ell - k_1)] - \\ - (\ell - k_1)^2 \text{Tr} [\gamma_5 (\ell + k_2) \gamma_\nu \not{\ell} \gamma_\mu] \Rightarrow$$

$$(k_1 + k_2)^\rho T_{\mu\nu\rho} = -e^2 \int \frac{d^4 \ell}{(2\pi)^4} \frac{1}{\ell^2 + i\varepsilon} \left\{ \frac{\text{Tr} [\gamma_5 \gamma_\mu \not{\ell} \gamma_\nu (\ell - k_2)]}{(\ell - k_2)^2 + i\varepsilon} \right. \\ + \frac{\text{Tr} [\gamma_5 (\ell + k_1) \gamma_\mu \not{\ell} \gamma_\nu]}{(\ell + k_1)^2 + i\varepsilon} + \frac{\text{Tr} [\gamma_5 \gamma_\nu \not{\ell} \gamma_\mu (\ell - k_1)]}{(\ell - k_1)^2 + i\varepsilon} \\ \left. + \frac{\text{Tr} [\gamma_5 (\ell + k_2) \gamma_\nu \not{\ell} \gamma_\mu]}{(\ell + k_2)^2 + i\varepsilon} \right\}$$

Now, if in (4) we shift $l \rightarrow l - k_2 \Rightarrow$ it

would cancel (1) as (1) + (4) $\propto \{\delta_5, \delta_6\} = 0$.

In (3) shift $l \rightarrow l + k_1 \Rightarrow$ cancel (2).

\Rightarrow seems to get $(k_1 + k_2)^\rho T_{\mu\nu\rho} = 0$ in

expectation with $\partial^\rho \int_S^\rho = 0 \dots$

Problem at large l all integrals are quadratically divergent!

We get (1) \sim (2) \sim (3) \sim (4) $\sim \int d^4l \frac{1}{l^2} \sim \int dl \cdot l \sim \infty^2$.

\Rightarrow can't shift variables in divergent integrals!

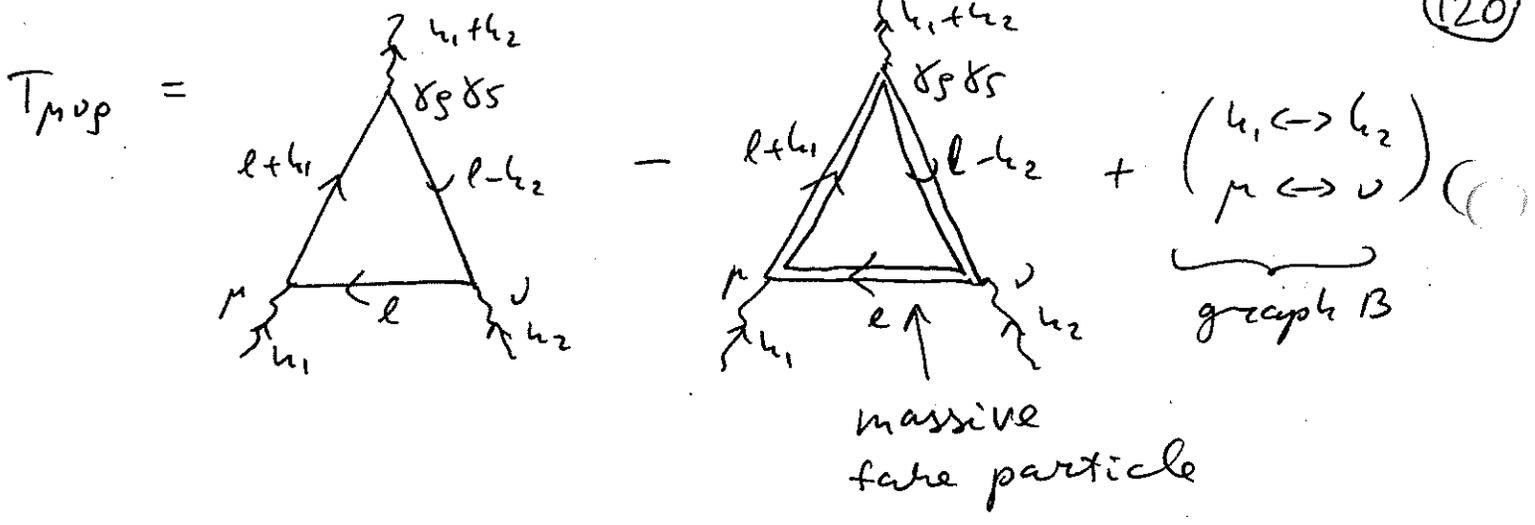
$$\int_0^\infty dl \cdot l \xrightarrow[\substack{\text{shift } l \rightarrow l+a}]{-a} \int_0^\infty dl \cdot (l+a) = \int_0^\infty dl \cdot (l+a) + \int_{-a}^0 dl \cdot (l+a)$$

$$= \int_0^\infty dl \cdot l + a \cdot \int_0^\infty dl + \left(\frac{l^2}{2} + ale \right) \Big|_{-a}^0 = \underbrace{\int_0^\infty dl \cdot l}_{\text{old integral}} + a \underbrace{\int_0^\infty dl + \frac{a^2}{2}}$$

\Rightarrow did not survive the shift, got corrections?

\Rightarrow ill-defined procedure \Rightarrow need to make integrals finite, need to regulate them!

(We'll use Pauli-Villars regularization: introduce ^(subtract) a new particle with mass m , which is then taken to ∞ to eliminate the particle.



$$T_{\mu\nu\rho} = -e^2 \int \frac{d^4 l}{(2\pi)^4} \left\{ \frac{\text{Tr} [\gamma_\rho \gamma_5 (\not{l} + \not{k}_1) \not{l} \gamma_\nu (\not{l} - \not{k}_2)]}{(l^2 + i\epsilon) ((l+k_1)^2 + i\epsilon) ((l-k_2)^2 + i\epsilon)} \right.$$

$$\left. - \frac{\text{Tr} [\gamma_\rho \gamma_5 (\not{l} + \not{k}_1 + \not{m}) \not{l} + \not{m} \gamma_\nu (\not{l} - \not{k}_2 + \not{m})]}{(l^2 - m^2 + i\epsilon) ((l+k_1)^2 - m^2 + i\epsilon) ((l-k_2)^2 - m^2 + i\epsilon)} \right\} + \left(\begin{matrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{matrix} \right)$$

The second Tr has only even powers of m in its expansion. (Tr of odd # of γ 's is zero.) Write:

$$(k_1 + k_2)^\rho T_{\mu\nu\rho} = -e^2 \int \frac{d^4 l}{(2\pi)^4} \left\{ \text{Tr} [(\not{k}_1 + \not{k}_2) \gamma_5 (\not{l} + \not{k}_1) \not{l} \gamma_\nu \cdot (\not{l} - \not{k}_2)] \left[\frac{1}{l^2 (l+k_1)^2 (l-k_2)^2} - \frac{1}{(l^2 - m^2) [(l+k_1)^2 - m^2] [(l-k_2)^2 - m^2]} \right] \right.$$

$$\left. - \frac{m^2 \text{-term in 2nd trace (O(l))}}{[l^2 - m^2] [(l+k_1)^2 - m^2] [(l-k_2)^2 - m^2]} \right\} + \left(\begin{matrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{matrix} \right)$$

Now the integral is convergent & shifts are allowed!

\Rightarrow the $m=0$ term in [...] vanishes like (121) before. (first)

For the term in [...] containing m^2 write:

$$\begin{aligned} & \text{Tr} \left[(\not{k}_1 + \not{k} - (\not{k} - \not{k}_2)) \gamma_5 (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu (\not{k} - \not{k}_2) \right] = \\ & = -(\not{k} + \not{k}_1)^2 \text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu (\not{k} - \not{k}_2)] - (\not{k} - \not{k}_2)^2 \text{Tr} [\gamma_5 (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu] \\ & \cdot \text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu] = -[(\not{k} + \not{k}_1)^2 - m^2] \text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu (\not{k} - \not{k}_2)] \\ & - [(\not{k} - \not{k}_2)^2 - m^2] \text{Tr} [\gamma_5 (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu] - m^2 (\text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu \cdot \\ & \cdot (\not{k} - \not{k}_2)] + \text{Tr} [\gamma_5 (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu]) \end{aligned}$$

First two terms also cancel after shifts.

We get:

$$\begin{aligned} (\not{k}_1 + \not{k}_2)^{\rho} T_{\text{loop}} &= -e^2 \int \frac{d^4 \ell}{(2\pi)^4} m^2 \int_0^1 \frac{1}{[e^2 - m^2][(\not{k} + \not{k}_1)^2 - m^2][(\not{k} - \not{k}_2)^2 - m^2]} \\ & \cdot \left\{ \text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu (\not{k} - \not{k}_2)] + \text{Tr} [\gamma_5 (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu] - \right. \\ & \left. - \text{Tr} [(\not{k}_1 + \not{k}_2) \gamma_5 \not{\epsilon} \gamma_\nu (\not{k} - \not{k}_2)] + \text{Tr} [(\not{k}_1 + \not{k}_2) \gamma_5 \right. \\ & \left. \cdot (\not{k} + \not{k}_1) \not{\epsilon} \gamma_\nu] - \text{Tr} [(\not{k}_1 + \not{k}_2) \gamma_5 \not{\epsilon} \gamma_\nu] \right\} + \left. \begin{matrix} m^2 \\ \text{terms} \\ \text{in} \\ \text{Tr.} \end{matrix} \right. \\ & \left. = \left(\text{as } \text{Tr} [\gamma_5 \not{\epsilon} \gamma_\nu \not{\alpha} \not{\beta}] = -4i \epsilon^{\mu\nu\alpha\beta} \right) = \dots \right. \end{aligned}$$

$$= 4i e^2 \varepsilon^{\mu\nu\alpha\beta} \int \frac{d^4 l}{(2\pi)^4} \frac{m^2}{(l^2 - m^2) [(l+k_1)^2 - m^2] [(l-k_2)^2 - m^2]} \quad (122)$$

$$\left\{ \begin{aligned} & -l_\alpha (l-k_2)_\beta + l_\alpha (l+k_1)_\beta - (l-k_2)_\alpha (k_1+k_2)_\beta + (k_1+k_2)_\alpha (l+k_1)_\beta \\ & - (k_1+k_2)_\alpha l_\beta \end{aligned} \right\}_\Lambda = 4i e^2 \varepsilon^{\mu\nu\alpha\beta} \int \frac{d^4 l}{(2\pi)^4}$$

+ $\begin{pmatrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{pmatrix}$

$$\frac{m^2}{[(l^2 - m^2) [(l+k_1)^2 - m^2] [(l-k_2)^2 - m^2]} \left\{ \cancel{l_\alpha k_{2\beta}} + \cancel{l_\alpha k_{1\beta}} - \right.$$

$$\left. \begin{aligned} & - \cancel{l_\alpha (k_1+k_2)_\beta} + \cancel{(k_1+k_2)_\alpha l_\beta} - \cancel{(k_1+k_2)_\alpha l_\beta} + k_{2\alpha} (k_1+k_2)_\beta \\ & + k_{1\beta} (k_1+k_2)_\alpha \end{aligned} \right\}_L = 8i e^2 \varepsilon^{\mu\nu\alpha\beta} k_{1\beta} k_{2\alpha} \int \frac{d^4 l}{(2\pi)^4}$$

+ $\begin{pmatrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{pmatrix}$

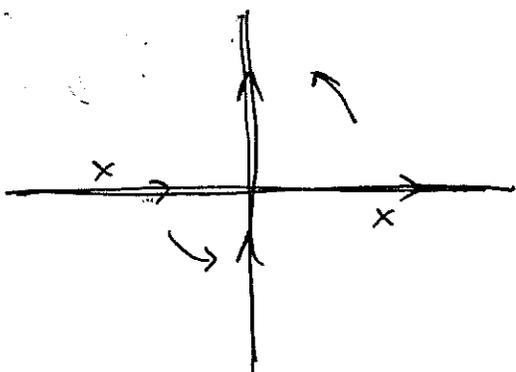
$$\frac{m^2}{(l^2 - m^2) [(l+k_1)^2 - m^2] [(l-k_2)^2 - m^2]} \quad \text{Im}$$

+ $\begin{pmatrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{pmatrix}$

$k_1, k_2 \ll m$ always
if $l \ll m \Rightarrow I_m \sim \frac{1}{m^3}$
 \nearrow if $l \sim m \Rightarrow I_m \gg k_1, k_2$

Approximate the integral by: ($l, m \sim$ large)

$$I_m \approx \int \frac{d^4 l}{(2\pi)^4} \frac{m^2}{[l^2 - m^2 + i\epsilon]^3} = \left| \begin{array}{l} \text{Wick rotation} \\ l_0 = +i l_0^E \end{array} \right.$$



$$l^2 - m^2 + i\epsilon = (l_0 - \sqrt{\vec{l}^2 + m^2} + i\epsilon)$$

$$\cdot (l_0 + \sqrt{\vec{l}^2 + m^2} - i\epsilon)$$