## Homework Set No. 2, Physics 880.08 Deadline – Monday, November 2, 2009

1. (10 pts) In class we showed that Dirac spinors transform as

$$\psi_D(x) \to \psi'_D(x') = \begin{pmatrix} e^{\frac{i}{2}\vec{\sigma} \cdot (\vec{\theta} - i\vec{\xi})} & 0\\ 0 & e^{\frac{i}{2}\vec{\sigma} \cdot (\vec{\theta} + i\vec{\xi})} \end{pmatrix} \psi_D(x)$$
(1)

under Lorentz transformations. Show that this transformation rule is equivalent to

$$\psi_D(x) \to \psi'_D(x') = e^{-\frac{i}{4}\,\omega^{\mu\,\nu}\,\sigma_{\mu\,\nu}}\,\psi_D(x)$$

where

$$\sigma_{\mu\nu} = \frac{i}{2} [\gamma_{\mu}, \gamma_{\nu}].$$

As usual  $\xi^i = \omega^{0i}$  and  $\theta_i = \frac{1}{2} \epsilon_{ijk} \omega_{jk}$ . You can consider boosts and rotations separately for the full credit.

2. (a) (5 pts) Complete the proof started in class that

 $\bar{\psi} \gamma^{\mu} \psi$ 

is a 4-vector by showing that it transforms like one under infinitesimal boosts.  $\psi$  is the Dirac spinor which transforms according to Eq. (1).

(b) (5 pts) Prove that

 $\bar\psi\,\gamma^\mu\,\gamma^5\,\psi$ 

is a 4-vector under both boosts and rotations. What happens to it under parity?

**3.** (10 pts) Writing Dirac spinor in terms of two-component Weyl spinors

$$\psi_D(x) = \left(\begin{array}{c} \chi_L \\ \chi_R \end{array}\right)$$

one can see from Eq. (1) above that Weyl spinors transform under Lorentz transformations as

$$\chi_L(x) \to \chi'_L(x') = e^{\frac{i}{2}\vec{\sigma} \cdot (\vec{\theta} - i\vec{\xi})} \chi_L(x)$$
  
$$\chi_R(x) \to \chi'_L(x') = e^{\frac{i}{2}\vec{\sigma} \cdot (\vec{\theta} + i\vec{\xi})} \chi_R(x).$$

Show that  $\sigma^2 \chi_L^*$  transforms as a right-handed Weyl spinor. Here \* denotes complex conjugation and  $\sigma^2$  is a Pauli matrix.

- 4. Consider a massive Dirac field  $\psi$  with mass m.
- **a.** (2 pts) Starting with Dirac equation

$$[i\gamma^{\mu}\partial_{\mu} - m]\psi = 0$$

derive an equation for  $\bar{\psi}$ .

**b.** (3 pts) Using the result of part **a** show that the electromagnetic current

$$j_{\mu} = \bar{\psi} \gamma_{\mu} \psi$$

is conserved at the classical level, i.e., show that  $\partial_{\mu}j^{\mu} = 0$ . **c.** (5 pts) Defining  $\gamma^5 = i \gamma^0 \gamma^1 \gamma^2 \gamma^3$  use the anti-commutation relations for  $\gamma$ -matrices  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2g^{\mu\nu}$  to show that

$$\left\{\gamma^{\mu},\gamma^{5}\right\} = 0.$$

Use this result to show that the divergence of the axial vector current

$$j^{5\,\mu} = \bar{\psi}\gamma^{\mu}\gamma^{5}\psi$$

is

$$\partial_{\mu} j^{5\,\mu} = 2\,i\,m\,\bar{\psi}\,\gamma^5\,\psi. \tag{2}$$

That is the axial current is conserved for massless particles (in this classical theory).