

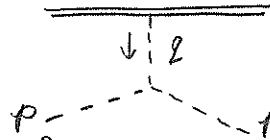
# Wavefunction Renormalization for a Scalar Field

The Lagrangian for a real scalar field with a cubic interaction is

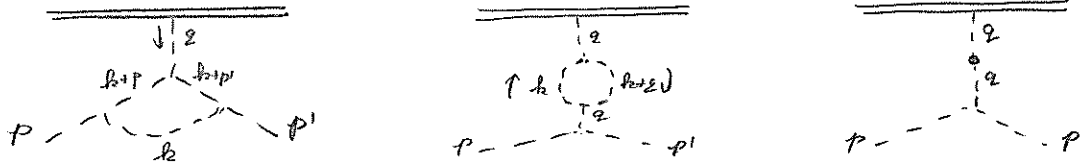
$$\mathcal{L} = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{2}m^2\phi^2 - \frac{1}{6}g\phi^3 - \frac{1}{2}\delta m^2\phi^2,$$

where  $\delta m^2$  is order  $g^2$ . The Feynman rules are  $i/(p^2 - m^2 + i\epsilon)$  for the scalar propagator,  $-ig$  for the 3-scalar vertex, and  $-i\delta m^2$  for the 2-scalar vertex. The scalar also couples to a heavy particle with a vertex  $-iG$ .

A. Draw the tree-level Feynman diagram of order  $g$  for the scattering of the scalar from the heavy particle, with momentum  $q = p' - p$  transferred to the scalar. Write down the matrix element  $i\mathcal{M}$ .



B. Draw the three Feynman diagrams of order  $g^3$  for the scattering of the scalar from the heavy particle, labeling the momenta of the scalar lines.



C. The contribution to  $i\mathcal{M}$  from the vertex-correction diagram has the form  $(-iG)(-ig)^3[i/(q^2 - m^2)]F(q^2)$ . Express the form factor  $F(q^2)$  as an integral over a loop momentum.

$$F(q^2) = \int \frac{d^4k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon} \frac{i}{(k+p)^2 - m^2 + i\epsilon} \frac{i}{(k+p')^2 - m^2 + i\epsilon}$$

The contribution to  $i\mathcal{M}$  from the two propagator-correction diagrams is

$$(-iG)(-ig) \frac{i}{q^2 - m^2} \left[ -i\Pi(q^2) - i\delta m^2 \right] \frac{i}{q^2 - m^2},$$

where the self-energy function is

$$\Pi(p^2) = -\frac{g^2}{32\pi^2} \left( \log \frac{\Lambda^2}{m^2} - 1 + \int_0^1 dx \log \frac{m^2 - x(1-x)p^2 - i\epsilon}{m^2} \right).$$

D. How must  $\delta m^2$  depend on the ultraviolet cutoff  $\Lambda$  so that  $\mathcal{M}$  does not depend on  $\Lambda$ ?  $\delta m^2 = \frac{g^2}{32\pi^2} \left( \log \frac{\Lambda^2}{m^2} + \text{constant} \right)$

E. What must  $\delta m^2$  be to avoid an unphysical double pole  $1/(q^2 - m^2)^2$ ?

$$\delta m^2 = -\mathcal{P}(m^2) = -\frac{g^2}{32\pi^2} \left[ \log \frac{\Lambda^2}{m^2} - 1 + \int_0^1 dx \log(1-x+x^2) \right]$$

The self-energy function and its derivative at  $p^2 = m^2$  are

$$\Pi(m^2) = -\frac{g^2}{32\pi^2} \left( \log \frac{\Lambda^2}{m^2} + \frac{\pi}{\sqrt{3}} - 3 \right), \quad \Pi'(m^2) = \frac{g^2}{32\pi^2 m^2} \left( \frac{2\pi}{3\sqrt{3}} - 1 \right).$$

The complete propagator obtained by summing the geometric series of order- $g^2$  self-energy diagrams to all orders is

$$D(p^2) = \frac{i}{p^2 - m^2 - \Pi(p^2) - \delta m^2 + i\epsilon}.$$

The pole in  $D(p^2)$  is at  $p^2 = m_{\text{phys}}^2$ , where  $m_{\text{phys}}$  is the physical mass.

F. Write down the equation that must be satisfied by  $m_{\text{phys}}^2$ .

$$m_{\text{phys}}^2 - m^2 - \Pi(m_{\text{phys}}^2) - \delta m^2 = 0$$

G. Suppose  $\delta m^2 = (g^2/32\pi^2) \log(\Lambda^2/m^2)$ ? Solve the equation for  $m_{\text{phys}}^2$  to order  $g^2$ . (In terms with an explicit factor of  $g^2$ , you can replace  $m_{\text{phys}}^2$  by  $m^2$ .)

$$\begin{aligned} m_{\text{phys}}^2 &= m^2 + \Pi(m^2) + \delta m^2 = m^2 - \frac{g^2}{32\pi^2} \left( \log \frac{\Lambda^2}{m^2} + \frac{\pi}{\sqrt{3}} - 3 \right) + \frac{g^2}{32\pi^2} \log \frac{\Lambda^2}{m^2} \\ &= m^2 - \frac{g^2}{32\pi^2} \left( \frac{\pi}{\sqrt{3}} - 3 \right) \end{aligned}$$

H. Suppose  $\delta m^2 = -\Pi(m^2)$ . What is the value of  $m_{\text{phys}}$ ?

$$m_{\text{phys}}^2 = m^2$$

I. Having chosen  $\delta m^2$  so  $m$  is the physical mass, determine the residue  $iZ_s$  of the pole in  $D(p^2)$  at  $p^2 = m^2$ , and expand  $Z_s$  to order  $g^2$ .

$$Z_s = \frac{1}{1 - \Pi'(m^2)} = \frac{1}{1 - \frac{g^2}{32\pi^2 m^2} \left( \frac{2\pi}{3\sqrt{3}} - 1 \right)} = 1 + \frac{g^2}{32\pi^2 m^2} \left( \frac{2\pi}{3\sqrt{3}} - 1 \right)$$

J. Draw the ~~two~~<sup>four</sup> Feynman diagrams of order  $g^3$  for scattering of the scalar from the heavy particle that have propagator corrections on external scalar lines.



The complete matrix element for scattering at next-to-leading order in  $g$  is

$$\mathcal{M} = \left( \sqrt{Z_s} \right)^2 \frac{Gg}{q^2 - m^2} \left[ 1 + g^2 F(q^2) + \frac{\Pi(q^2) - \Pi(m^2)}{q^2 - m^2} + \mathcal{O}(g^4) \right].$$

K. Expand this out to obtain the complete order  $g^3$  contribution to  $\mathcal{M}$ .

$$\mathcal{M} = \frac{Gg}{q^2 - m^2} + \frac{Gg^3}{q^2 - m^2} \left[ F(q^2) + \frac{\Pi(q^2) - \Pi(m^2)}{q^2 - m^2} + \frac{1}{32\pi^2 m^2} \left( \frac{2\pi}{3\sqrt{3}} - 1 \right) \right]$$