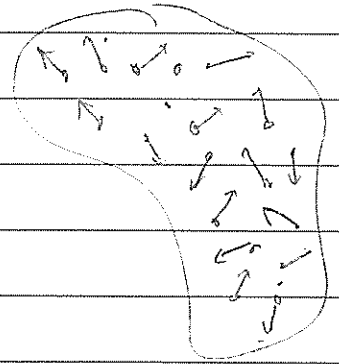


# Classical Ideal Gas

$N$  atoms (with volume  $v_0$ ) confined in volume  $V$

$$\log N \gg 1, \log \frac{V}{v_0} \gg 1$$

total energy  $E$



macrostate:  $(N, V, E)$

microstate:  $(\vec{r}_1, \vec{p}_1, \vec{r}_2, \vec{p}_2, \dots, \vec{r}_N, \vec{p}_N)$

each  $\vec{r}_n$  is inside volume  $V$

$$H = \sum_{n=1}^N \frac{1}{2m} \vec{p}_n^2$$

## Microcanonical Ensemble

number of microstates with energy  $E$

with energy within  $\frac{\Delta}{2}$  of  $E$ , where  $\log \frac{E}{\Delta} \ll N$

$$\Omega(E, V, N) = \frac{1}{\omega_0} \times \left( \text{volume of phase space with } E - \frac{\Delta}{2} < H < E + \frac{\Delta}{2} \right)$$

$$= \frac{1}{\omega_0} \int dw \delta(H(p, r) - E) \Delta$$

$\omega_0 =$  normalizing factor  
with dimensions of phase space

$$dw = d^3r_1 d^3p_1 \dots d^3r_N d^3p_N$$

entropy:  $S(E, V, N) = \log \Omega(E, V, N)$

temperature:  $T \quad \frac{1}{T} = \left( \frac{\partial S}{\partial E} \right)_{V, N}$

pressure:  $P \quad \frac{P}{T} = - \left( \frac{\partial S}{\partial V} \right)_{E, N}$

Hamiltonian:  $H(p) = \sum_{n=1}^N \frac{1}{2m} \vec{p}_n^2$  if each  $\vec{r}_n$  is inside volume  $V$

$$\begin{aligned} \Omega(E, V, N) &= \frac{1}{\omega_0} \int d^3r_1 d^3p_1 \cdots \int d^3r_N d^3p_N \delta(H(p) - E) \Delta \\ &= \frac{1}{\omega_0} \underbrace{\int d^3r_1 \cdots \int d^3r_N}_{V^N} \int d^3p_1 \cdots \int d^3p_N \delta(H(p) - E) \Delta \end{aligned}$$

$$\Omega(E, V, N) = V^N \times (\text{function of } E, N)$$

$$S(E, V, N) = N \log V + (\text{function of } E, N)$$

$$\left( \frac{\partial S}{\partial V} \right)_{E, N} = \frac{N}{V}$$

$$\frac{P}{T} = \frac{N}{V} \implies \boxed{PV = NT}$$

ideal gas law:  $PV = Nk_B T$  if temperature is measured in K and not energy units

$$\Omega(E, V, N) = \frac{1}{\omega_0} V^N \underbrace{\int d^3 p_1 \dots \int d^3 p_N}_{\text{dimension } P^{3N}} \underbrace{\delta(H(p) - E) \Delta}_{\text{dimensionless}}$$

deduce dependence on  $E$  using dimensional analysis

$$\sum_{n=1}^N \frac{1}{2m} \vec{p}_n^2 = E \implies 2mE \text{ has dimension } P^2$$

$$(2mE)^{3N/2} \quad \text{''} \quad P^{3N}$$

$$\Omega(E, V, N) = \frac{1}{\omega_0} V^N (2mE)^{3N/2} \times (\text{function of } N)$$

$$S(E, V, N) = N \log V + \frac{3N}{2} \log(2mE) + (\text{function of } N)$$

$$\left( \frac{\partial S}{\partial E} \right)_{N, V} = \frac{3N}{2E}$$

$$\frac{1}{T} = \frac{3N}{2E}$$

$$\implies \boxed{E = \frac{3}{2} NT}$$

energy of monatomic ideal gas:  $E = \frac{3}{2} N k_B T$

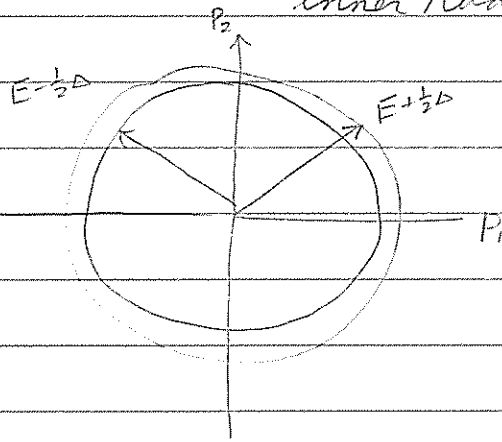
To complete calculation of  $\Omega(E, V, N)$ ,  
calculate momentum integral

$$\int d^3 p_1 \dots \int d^3 p_N \delta(H(p) - E) \Delta$$

$\approx$  volume of  $3N$ -dimensional momentum space

$$\text{satisfying } E - \frac{1}{2}\Delta < \sum_{n=1}^N \frac{1}{2m} \vec{p}_n^2 < E + \frac{1}{2}\Delta$$

= volume of shell in  $3N$ -dimensional momentum space  
 with outer radius  $2m(E + \frac{1}{2}\Delta)$   
 inner radius  $2m(E - \frac{1}{2}\Delta)$



hypervolume of ball of radius  $R$  in  $n$  dimensions

$$V_n = \frac{\pi^{n/2}}{(n/2)!} R^n \quad P+B, \text{ Appendix C}$$

where  $z! \equiv \Gamma(z+1) = z \Gamma(z)$ ,  $\Gamma(1) = 1$ ,  $\Gamma(1/2) = \sqrt{\pi}$

checks  $V_1 = \frac{\pi^{1/2}}{1/2!} R^1 = \frac{\sqrt{\pi}}{\Gamma(3/2)} R = \frac{\sqrt{\pi}}{1/2\sqrt{\pi}} R = 2$

$$V_2 = \frac{\pi}{1!} R^2 = \pi R^2$$

$$V_3 = \frac{\pi^{3/2}}{(3/2)!} R^3 = \frac{\pi\sqrt{\pi}}{\Gamma(5/2)} R^3 = \frac{\pi\sqrt{\pi}}{3/2 \cdot 1/2\sqrt{\pi}} R^3 = \frac{4}{3}\pi R^3$$

$$\int d^3 p_1 \dots \int d^3 p_{3N} S(H(p) - E) \Delta$$

$$\approx \frac{\pi^{3N/2}}{(3N/2)!} [2m(E + \frac{1}{2}\Delta)]^{3N/2} - \frac{\pi^{3N/2}}{(3N/2)!} [2m(E - \frac{1}{2}\Delta)]^{3N/2}$$

$$\approx \frac{\pi^{3N/2}}{(3N/2)!} \frac{d}{dE} [2mE]^{3N/2} \cdot \Delta = \frac{\pi^{3N/2}}{(\frac{3N}{2}-1)!} (2mE)^{3N/2-1} \cdot 2m\Delta$$

$$\Omega(E, V, N) = \frac{1}{\omega_0} V^N \frac{\pi^{3N/2}}{(3N/2)!} (2mE)^{3N/2} \frac{3NA}{2E}$$

$$S(E, V, N) = N \log V + \frac{3N}{2} \log(2mE) - \log\left(\left(\frac{3}{2}N\right)!\right) \\ - \log \omega_0 + \log N + \log \frac{3}{2} + \log \frac{\Delta}{E}$$

Stirling's formula for  $n!$

$$\log n! = n \log n - n + \frac{1}{2} \log n + \frac{1}{2} \log(2\pi) + O\left(\frac{1}{n}\right) \\ \approx n \log n - n$$

Keep only terms in  $S$  with explicit factor of  $N$

$$S(E, V, N) \approx N \log V + \frac{3N}{2} \log(2mE) \\ - \left[ \frac{3}{2}N \log\left(\frac{3}{2}N\right) - \frac{3}{2}N \right] - \log \omega_0 \\ = N \left[ \log V + \frac{3}{2} \log(2mE) - \frac{3}{2} \log\left(\frac{3}{2}N\right) + \frac{3}{2} \right] - \log \omega_0 \\ = N \left[ \log V + \frac{3}{2} \log \frac{4mE}{3N} + \frac{3}{2} \right] - \log \omega_0$$

Note:  $S$  is not exactly extensive because of  $\log V$  term

$$S(\lambda E, \lambda V, \lambda N) = \lambda N \left[ \log(\lambda V) + \frac{3}{2} \log \frac{4mE}{3N} + \frac{3}{2} \right] - \log \omega_0 \\ = \lambda S(E, V, N) + \lambda N \log \lambda$$

# Gibbs paradox

two systems consisting of same atom  
at same temperature  $T$ , pressure  $P$   
separated by barrier

$E_1, V_1, N_1$	$E_2, V_2, N_2$
$T, P$	$T, P$

if barrier is removed, macrostate is the same

$E_1 + E_2, N_1 + N_2, V_1 + V_2$
$P, T$

entropy before:  $S(E_1, V_1, N_1) + S(E_2, V_2, N_2)$

entropy after:  $S(E_1 + E_2, V_1 + V_2, N_1 + N_2)$

change in entropy:

$$\Delta S = (N_1 + N_2) \log(V_1 + V_2) - N_1 \log V_1 - N_2 \log V_2$$

$$= N_1 \log \frac{V_1 + V_2}{V_1} + N_2 \log \frac{V_1 + V_2}{V_2}$$

removal of barrier: same microstate, different entropy

Gibbs fudge factor: divide  $\Omega$  by  $N!$

$$\Omega \rightarrow \frac{1}{N!} \Omega$$

$$S \rightarrow S - \log N!$$

$$\simeq S - (N \log N - N)$$

$$S(E, V, N) = N \left[ \log \frac{V}{N} + \frac{3}{2} \log \frac{4mE}{3N} + \frac{5}{2} \right]$$

exactly extensive  $\Rightarrow$  no Gibbs paradox

### Quantum solution

atoms of same isotope are identical indistinguishable particles

interchange of any 2 atoms gives same quantum state  
permutation of  $N$  atoms " "

integrated over phase space of  $N$  atoms

overcounts quantum states by factor of  $N!$

correct counting of quantum states

$$\Omega(E, V, N) = \frac{1}{N!} \frac{1}{\omega_0} \int d\omega \delta(H(\mathbf{p}, \mathbf{q}) - E) \Delta$$