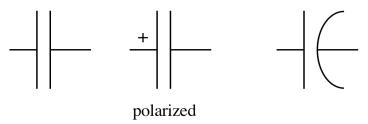
Lecture 2: Resistors and Capacitors

Capacitance:

- Capacitance (C) is defined as the ratio of charge (Q) to voltage (V) on an object:
 - C = Q/V = Coulombs/Volt = Farad
 - Capacitance of an object depends on geometry and its dielectric constant.
 - Symbol(s) for capacitors:



- A capacitor is a device that stores electric charge (memory devices).
- A capacitor is a device that stores energy

$$E = \frac{Q^2}{2C} = \frac{CV^2}{2}$$

- Capacitors are easy to fabricate in small sizes (μm), use in chips
- How to combine capacitance:
 - capacitors in parallel adds like resistors in series:

Total capacitance is *more* than individual capacitance!

• capacitors in series add like resistors in parallel:

Total capacitance is less than individual capacitance!

- Energy and Power in Capacitors
 - How much energy can a "typical" capacitor store?
 - Pick a 4 µF Cap (it would read 4 mF) rated at 3 kV $E = \frac{1}{2}CV^2 = \frac{1}{2}4 \times 10^{-6} \cdot 3000^2 = 18 \text{ J}$
 - This is the same as dropping a 2 kg weight (about 4 pounds) 1 meter
 - How much power is dissipated in a capacitor?

$$Power = \frac{dE}{dt}$$
$$= \frac{d}{dt} \left(\frac{CV^{2}}{2} \right)$$
$$P = CV \frac{dV}{dt}$$

dV/dt must be finite otherwise we source (or sink) an infinite amount of power!

THIS WOULD BE UNPHYSICAL.

- the voltage across a capacitor cannot change instantaneously
 - a useful fact when trying to guess the transient (short term) behavior of a circuit
- the voltage across a resistor can change instantaneously
 - \Box the power dissipated in a resistor does not depend on dV/dt:

$$P = I^2 R$$
 or V^2/R

- Why do capacitors come in such small values?
 - **Example:** Calculate the size of a 1 Farad parallel capacitor with air between the plates.

$$C = \frac{k\varepsilon_0 A}{d}$$

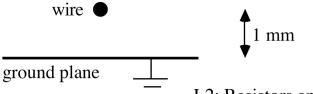
k = dielectric constant (= 1 for air)

$$\varepsilon_0 = 8.85 \times 10^{-12} N^{-1} m^{-2}$$

d = distance between plates (assumed 1 mm)

$$A = \text{area of plates} = 1.1 \times 10^8 \, m^2!!!$$

- square plate of 6.5 miles per side
- breakthroughs in capacitor technologies (driven by the computer industries)
 - fabrication of 0.5-5 F capacitors of small size (1-2 cm high) and low cost (< \$5)
- How small can we make capacitors?
 - A wire near a ground plane has $C \sim 0.1$ pf = 10^{-13} F.

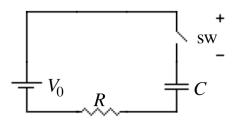


K.K. Gan

L2: Resistors and Capacitors

- Some words to the wise on capacitors and their labeling.
 - Typical capacitors are multiples of micro Farads (10⁻⁶ F) or picoFarads (10⁻¹² F).
 - Whenever you see mF it almost always is micro, *not* milli F and *never* mega F.
 - picoFarad (10⁻¹² F) is sometimes written as pf and pronounced *puff*.
 - no *single* convention for labeling capacitors
 - Many manufacturers have their own labeling scheme (See Horowitz and Hill lab manual).
- Resistors and Capacitors
 - Examine voltage and current vs. time for a circuit with one R and one C.
 - Assume that at t < 0 all voltages are zero, $V_R = V_C = 0$.
 - At $t \ge 0$ the switch is closed and the battery (V_0) is connected.
 - Apply Kirchhoff's voltage rule:

$$\begin{aligned} V_0 &= V_R + V_C \\ &= IR + \frac{Q}{C} \\ &= R \frac{dQ}{dt} + \frac{Q}{C} \end{aligned}$$



Solve the differential equation by differentiating both sides of above equation:

$$\frac{dV_0}{dt} = \frac{1}{C} \frac{dQ}{dt} + R \frac{d^2Q}{dt^2}$$

$$0 = \frac{I}{C} + R \frac{dI}{dt}$$

$$\frac{dI}{dt} = -\frac{I}{RC}$$

- This is just an exponential decay equation with time constant RC (sec).
- The current as a function of time through the resistor and capacitor is:

$$I(t) = I_0 e^{-t/RC}$$

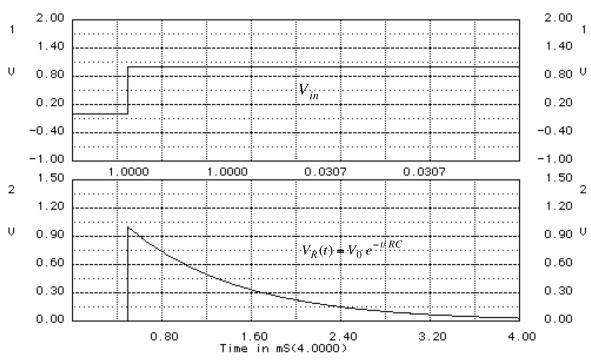
- What's $V_R(t)$?
 - By Ohm's law:

$$V_R(t) = I_R \cdot R$$

$$= I_0 R e^{-t/RC}$$

$$= V_0 e^{-t/RC}$$

- At t = 0 all the voltage appears across the resistor, $V_R(0) = V_0$.



- What's $V_C(t)$?
 - Easiest way to answer is to use the fact that $V_0 = V_R + V_C$ is valid for all t.

$$V_C = V_0 - V_R$$

$$V_C = V_0 \left(1 - e^{-t/RC} \right)$$

- At t = 0 all the voltage appears across the resistor so $V_C(0) = 0$.
- At $t = \infty$, $V_C(\infty) = V_0$.
- Suppose we wait until I = 0 and then short out the battery.

$$0 = V_R + V_C$$

$$V_R = -V_C$$

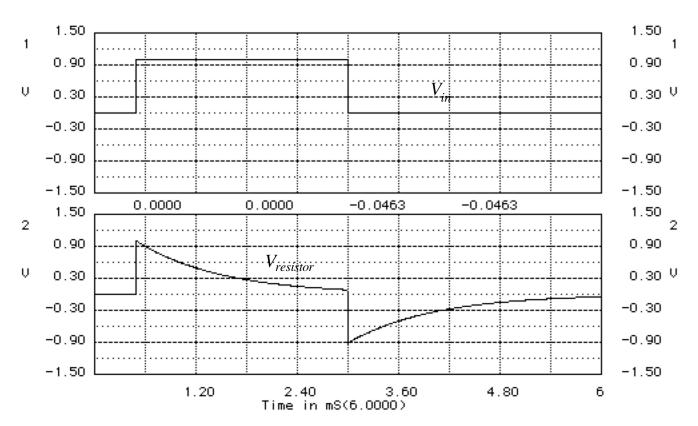
$$R\frac{dQ}{dt} = -\frac{Q}{C}$$

$$\frac{dQ}{dt} = -\frac{Q}{RC}$$

Solving the exponential equation yields,

$$Q(t) = Q_0 e^{-t/RC}$$

- We can find V_C using V = Q/C, $V_C(t) = V_0 e^{-t/RC}$
- Finally we can find the voltage across the resistor using $V_R = -V_{C_1}$ $V_R(t) = -V_0 e^{-t/RC}$



- Suppose $V_C(t) = V_0 \sin \omega t$ instead of DC
 - What happens to V_C and I_C ?

$$Q(t) = CV(t)$$
$$= CV_0 \sin \omega t$$

$$I_C = dQ/dt$$

$$=\omega CV_0\cos\omega t$$

$$=\omega CV_0\sin(\omega t+\pi/2)$$

current in capacitor varies like a sine wave too, but 90° out of phase with voltage.

We can write an equation that looks like Ohm's law by defining V^* :

$$V^* = V_0 \sin(\omega t + \pi/2)$$

the relationship between the voltage and current in C looks like:

$$V^* = I_C / \omega C$$
$$= I_C R^*$$

 $1/\omega C$ can be identified as a kind of resistance, <u>capacitive reactance</u>:

$$X_C = 1/\omega C \text{ (Ohms)}$$

- $X_C = 0$ for $\omega = \infty$
 - high frequencies: a capacitor looks like a short circuit
- $X_C = \infty \text{ for } \omega = 0$
 - low frequencies: a capacitor looks like an open circuit (high resistance).

Inductance:

- Define inductance by: V = LdI/dt
 - Unit: Henry
 - Symbol:

- Inductors are usually made from a coil of wire
 - tend to be bulky and are hard to fabricate in small sizes (μm), not used in chips.
- Two inductors next to each other (transformer) can step up or down a voltage
 - no change in the frequency of the voltage
 - provide isolation from the rest of the circuit

• How much energy is stored in an inductor?

$$dE = VdQ$$

$$I = \frac{dQ}{dt}$$

$$dE = VIdt$$

$$V = L \frac{dI}{dt}$$

$$dE = LIdI$$

$$E = L \int_0^I I dI$$

$$E = \frac{1}{2}LI^2$$

• How much power is dissipated in an inductor?

$$Power = \frac{dE}{dt}$$
$$= \frac{d}{dt} \left(\frac{LI^2}{2} \right)$$

$$P = LI \frac{dI}{dt}$$

 \Box dI/dt must be finite as we can't source (or sink) an infinite amount of power in an inductor!

THIS WOULD BE UNPHYSICAL.

the *current* across an inductor *cannot* change *instantaneously*.

• Two inductors in series:

Apply Kirchhoff's Laws,

$$V = V_1 + V_2$$

$$= L_1 \frac{dI}{dt} + L_2 \frac{dI}{dt}$$

$$= L_{tot} \frac{dI}{dt}$$

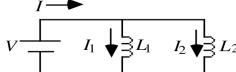
$$L_{tot} = L_1 + L_2$$

$$= \sum L_i$$

Inductors in series add like resistors in series.

The *total* inductance is *greater* than the individual inductances.

Two inductors in parallel:



Since the inductors are in parallel,

$$V_1 = V_2 = V$$

The total current in the circuit is

$$I = I_1 + I_2$$

$$\frac{dI}{dt} = \frac{dI_1}{dt} + \frac{dI_2}{dt}$$

$$= \frac{V}{L_1} + \frac{V}{L_2}$$

$$= \frac{V}{L_{tot}}$$

$$\frac{1}{L_{tot}} = \frac{1}{L_1} + \frac{1}{L_2}$$

$$L_{tot} = \frac{L_1 L_2}{L_2}$$

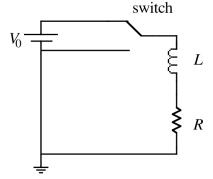
If we have more than 2 inductors in parallel, they combine like: $\frac{1}{L_{tot}} = \sum_{i=1}^{L} \frac{1}{L_{i}}$

$$\frac{1}{L_{tot}} = \sum \frac{1}{L_i}$$

Inductors in parallel add like resistors in parallel.

The *total* inductance is *less* than the individual inductances.

- Resistors and Inductors
 - Examine voltage and current versus time for a circuit with one R and one L.
 - Assume that at t < 0 all voltages are zero, $V_R = V_L = 0$.
 - At $t \ge 0$ the switch is closed and the battery (V_0) is connected.



Like the capacitor case, apply Kirchhoff's voltage rule:

$$V_0 = V_R + V_L$$
$$= IR + L \frac{dI}{dt}$$

Solving the differential equation, assuming at t = 0, I = 0:

$$I(t) = \frac{V_0}{R} \left(1 - e^{-tR/L} \right)$$

- This is just an exponential decay equation with time constant L/R (seconds).
- What's $V_R(t)$?
 - By Ohm's law $V_R = I_R R$ at any time:

$$V_R = I(t)R = V_0 (1 - e^{-tR/L})$$

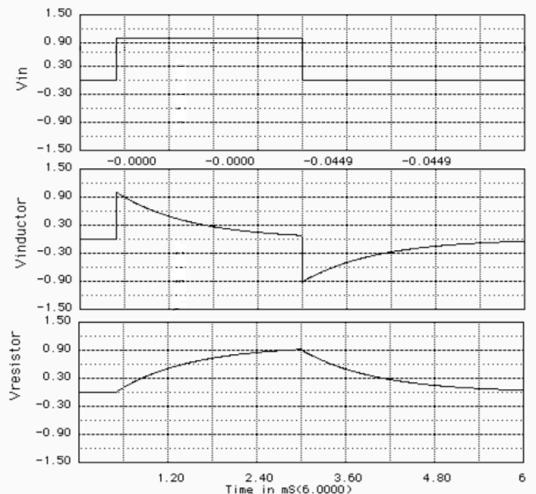
- At t = 0, none of the voltage appears across the resistor, $V_R(0) = 0$.
- $At t = \infty, V_R(\infty) = V_0.$

- What's $V_L(t)$?
 - Easiest way to answer is to use the fact that $V_0 = V_R + V_L$ is valid for all t.

$$V_L = V_0 - V_R$$

$$V_L(t) = V_0 e^{-tR/L}$$

- At t = 0, all the voltage appears across the inductor so $V_L(0) = V_0$.
- At $t = \infty$, $V_L(\infty) = 0$.



Pick L/R = 1 ms

L2: Resistors and Capacitors

• Suppose $V_C(t) = V_0 \sin \omega t$ instead of DC, what happens to V_L and I_L ?

$$V = L \frac{dI_L}{dt}$$

$$I_L = \frac{1}{L} \int_0^t V dt$$

$$= -\frac{V_0}{\omega L} \cos \omega t$$

$$I_L(t) = \frac{V_0}{\omega L} \sin(\omega t - \pi/2)$$

- The current in an inductor varies like a sine wave too, but 90° out of phase with the voltage.
- We can write an equation that looks like Ohm's law by defining V^* :

$$V^* = V_0 \sin(\omega t - \pi/2)$$

$$V^* = I_L \omega L = I_L R^*$$

 \square ωL can be identified as a kind of resistance, <u>inductive reactance</u>:

$$X_L = \omega L \text{ (Ohms)}$$

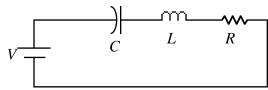
$$X_L = 0 \quad \text{if } \omega = 0$$

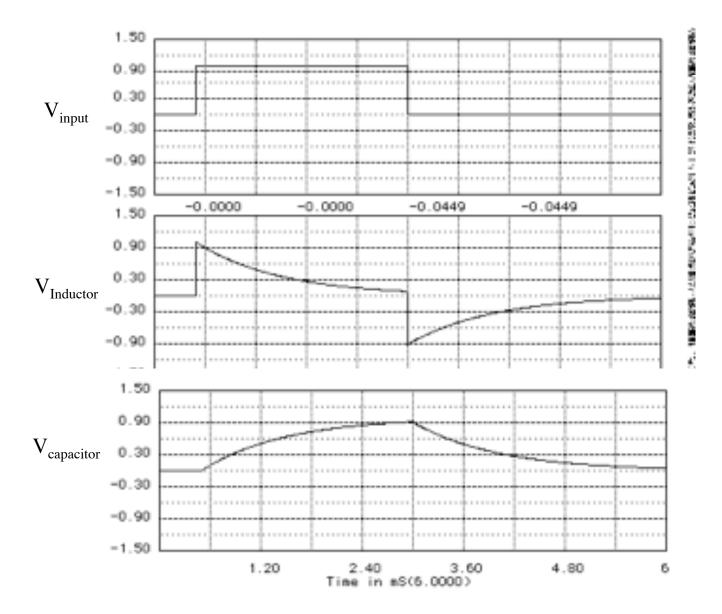
low frequencies: an inductor looks like a short circuit (low resistance).

$$X_L = \infty$$
 if $\omega = \infty$

high frequencies: an inductor looks like an open circuit.

- Some things to remember about *R*, *L*, and *C*'s.
 - For DC circuits, after many time constants (L/R or RC):
 - Inductor acts like a wire (0Ω) .
 - Capacitor acts like an open circuit ($\infty \Omega$).
 - For circuits where the voltage changes very rapidly or transient behavior:
 - Arr Capacitor acts like a wire (0Ω) .
 - Inductor acts like an open circuit ($\propto \Omega$).
 - Example, RLC circuit with DC supply:
 - At t = 0, voltages on R, C are zero and $V_L = V_{0.}$
 - At $t = \infty$, voltages on R, L are zero and $V_C = V_0$.





L2: Resistors and Capacitors