



# Ground-State Orbital Analysis Predicts S<sub>1</sub> Charge Transfer in Donor–Acceptor Materials

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ABSTRACT:	Donor–acceptor (D–A) materials	can	exhibit a wide range of unique	2000

Abstroact: Donor-acceptor (D-A) materials can exhibit a wide range of unique photophysical properties with applications in next-generation optoelectronics. Electronic structure calculations of D–A dimers are often employed to predict the properties of D–A materials. One of the most important D–A dimer quantities is the degree of charge transfer (DCT) in the S<sub>1</sub> state, which correlates with properties such as fluorescence lifetimes and intersystem crossing rates in D–A materials. While predictive metrics of the S<sub>1</sub> DCT generally require an excited-state quantum chemistry calculation, presented here is a novel metric that predicts S<sub>1</sub> DCT solely with ground-state orbital analysis. This metric quantifies the similarity of the orbitals between a dimer complex and its monomer components. A linear relationship is found between this similarity metric and the S<sub>1</sub> DCT, calculated using a data set of 31 D–A dimers. Best practices for integrating this novel orbital structure–function relationship into high-throughput screening methods are discussed.



In the search for next-generation optoelectronic devices, there has been a growing interest in donor-acceptor (D-A) materials, including D-A co-crystals, for their application in organic solar cells and organic light-emitting diodes.<sup>1-3</sup> D-A complexes, by definition, exhibit charge transfer (CT) in their ground state and select excited states.<sup>4</sup> In the search for structure-function relationships to integrate into highthroughput screening and machine learning protocols,<sup>5-</sup> several studies have assessed ways to predict the degree of CT (DCT), or ionicity parameter, in the  $S_0$  state of D-A materials from molecular quantities such as orbital energies, vibrational frequencies, and geometric parameters.<sup>10-14</sup> Early evidence shows a relationship between DCT in S<sub>0</sub> to the magnitude of effective CT integrals,<sup>11</sup> commonly used in models of charge transport.<sup>13–18</sup> The DCT in the first electronically excited state of D-A dimers, S1, has emerged as a key quantity for predicting radiative and nonradiative lifetimes in D-A materials, including intersystem crossing rates and fluorescence lifetimes.<sup>2,19-23</sup> These lifetimes are particularly difficult to compute directly using D-A dimer models, as energy levels and transition dipoles often differ substantially between the molecular cluster and material.<sup>24,25</sup>

A variety of DCT metrics for excited-state calculations have been put forward, as reviewed recently.<sup>26,27</sup> A widely used example is the " $\Lambda$  metric" introduced by Peach et al.<sup>28</sup> as a diagnostic for time-dependent density functional theory calculations. The definition of  $\Lambda$  is based on spatial overlaps of occupied and virtual orbitals, evaluated by numerical quadrature and weighted by excitation amplitudes, but importantly, this and other standard DCT metrics require an excited-state calculation. Here, we consider whether groundstate orbital overlaps are sufficient for indicating CT character in the  $S_0 \to S_1$  transition.

We construct a similarity metric,  $\eta$ , as follows. First, we define a molecular orbital (MO)

$$|\phi_i\rangle = \sum_{\mu} C_{\mu i} |\chi_{\mu}\rangle \tag{1}$$

where  $\phi_i$  is the sum of atomic basis functions  $\chi_{\mu}$  with MO coefficients  $C_{\mu i}$ . We compute the overlap  $O_{ij}$  between the MOs  $\phi_i$  and  $\tilde{\phi}_j$ , where  $\tilde{\phi}_j$  uses the same atomic basis functions as  $\phi_i$  at a displaced geometry. This overlap is

$$O_{ij} = \langle \phi_i | \tilde{\phi}_j \rangle = \sum_{\mu\nu} C_{\mu i} S_{\mu\nu} \tilde{C}_{\nu j}$$
(2)

where  $S_{\mu\nu} = \langle \chi_{\mu} | \tilde{\chi}_{\nu} \rangle$  is the overlap matrix involving displaced basis functions. We compute  $O_{ij}$  twice: once between the HOMO (highest occupied molecular orbital) of the isolated donor molecule (HOMO<sup>*i*</sup>) and the HOMO of the donor molecule within the complex (HOMO<sup>*c*</sup>), using ghost functions to ensure that both calculations have the same basis functions; and second, between the LUMO (lowest unoccupied molecular orbital) of the isolated acceptor molecule (LUMO<sup>*i*</sup>) and the LUMO of the acceptor molecule within

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the complex (LUMO<sup>c</sup>). We define  $\eta$  as the average of these two quantities:

$$\eta = \frac{1}{2} (O_{\text{HOMO}^{i}, \text{HOMO}^{c}} + O_{\text{LUMO}^{i}, \text{LUMO}^{c}})$$
(3)

The two geometries are oriented using the Kabsch algorithm<sup>29</sup> in order to maximize the overlap. Each computation of  $\eta$  requires three ground-state calculations: to obtain the orbitals of the isolated donor molecule, the orbitals of the isolated acceptor molecule, and the orbitals of the D–A complex. The calculation of  $O_{ij}$  is performed in a locally modified version of Q-Chem.<sup>30</sup>

To assess the correlation between  $\eta$  and the S<sub>1</sub> DCT of D–A dimers, we screened 31 D–A complexes with the donor and acceptor molecules shown in Figure 1, whose S<sub>1</sub> states are



Figure 1. Chemical structures of the donor and acceptor molecules.

dominated by a HOMO  $\rightarrow$  LUMO transition. The chosen donor and acceptor molecules are augmented from a data set recently chosen in a screening of S<sub>0</sub> DCT.<sup>14</sup> The donor molecules exhibit a diversity of molecular structures, while the acceptor molecules are 7,7,8,8-tetracyanoquinodimethane (TCNQ) and its fluorinated derivatives, F<sub>x</sub>TCNQ. Geometries are optimized with Gaussian G16 at the CAM-B3LYP/6-31+G(d,p) level of theory with Grimme D3 dispersion.<sup>28,31-33</sup> Excited-state calculations were also carried out using CAM-B3LYP-D3/6-31+G(d,p) within the linear response timedependent density functional theory (TDDFT) formalism.<sup>26,34-36</sup> We calculated the DCT using natural bond orbital (NBO) population analysis implemented in Gaussian G16 and transition density matrix (TDM) analysis in TheoDORE.<sup>37,38</sup>

We investigate the impact of geometries used to compute  $\eta$  by plotting  $\eta$  versus  $S_1$  DCT in two ways. First, we optimize both the isolated monomers and their dimer complex and compute  $\eta$  versus  $S_1$  DCT (Figure 2, top). Second, we optimize the dimer complex and take the geometries of the isolated monomers to be the same as those in their dimer complex (Figure 2, bottom). In each case, the  $S_1$  DCT is computed by using TDM analysis. Notably, we observed significant variations in the  $\eta$  values when the monomers were optimized, as evidenced by the outliers shown in the top plot of Figure 2. The  $R^2$  value from linear regression is 0.66 when



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Figure 2. Plots of  $\eta$  versus  $S_1$  DCT. Top: values computed at optimized monomer geometries and optimized dimer geometries. Bottom: values computed at the optimized dimer geometries only.

comparing optimized monomers versus optimized dimers, while the  $R^2$  value is 0.96 when monomer geometries are unrelaxed from those found in the dimer complex. In both cases, our analysis reveals a positive linear correlation between the S<sub>1</sub> DCT and  $\eta_i$  indicating that the S<sub>1</sub> DCT is large when the HOMO (LUMO) orbital of the donor (acceptor) retains its character from the isolated molecule. To assess the underlying reasons for the impact of the geometry choice on the evaluation of  $\eta$ , we plot the change in the average atomic root mean square deviation (RMSD) between geometries alongside the absolute difference in  $\eta$ , as shown in the Supporting Information, Figure S1. We show a correlation between the average atomic RMSD and absolute difference in  $\eta_{1}$ , indicating that the orbital character is likely to change most when the relaxed monomer geometries are significantly different from the geometries they assume in the dimer complex.

There are certain practical advantages that arise with the finding that the unrelaxed monomer geometries provide superior performance, including the elimination of the computational cost associated with optimizing the monomer geometries. Moreover, computing  $\eta$  at a single geometry eliminates the necessity of evaluating the atomic overlap integrals at displaced geometries. This simplifies the calculation of  $\eta$ , as  $S_{\mu\nu}$  becomes the atomic basis self-overlap matrix, which is commonly printed in the output of electronic structure programs.<sup>31</sup>

To assess the sensitivity of  $\eta$  to the dimer geometry, thus bridging the gap between D–A dimers and D–A co-crystals, we perform an analysis of the S<sub>1</sub> DCT versus  $\eta$  for geometries from experimental crystallographic data (where available) and compare these results to those found in Figure 2. S<sub>1</sub> DCT is again computed using TDM analysis. Figure 3 shows the



**Figure 3.** Plot of  $\eta$  versus S<sub>1</sub> DCT at the experimental crystal structures (orange) and optimized D–A dimer geometries (blue).

results obtained using the geometries from the experimental data in comparison to the dimer geometries, which indicate that the positive linear correlation still holds when the D–A geometries are taken from experimental crystal structures. In fact, when linear regressions are performed separately, the trend lines are almost indistinguishable. In Table 1, we report the value of  $\eta$  computed at each geometry (dimer complex versus experiment) and the percent deviation between the two. The values of  $\eta$  obtained from these two geometries have very small percent deviations, from 0.1% to 6% with an average of 1.3%. This shows that  $\eta$  can bridge between different types of molecular structures (isolated dimer and experimental crystal), indicating the reliability of the metric for different data sources. This is challenging for predictive metrics that rely on orbital energies.

We next assessed two methodologies for calculating the correlation of S<sub>1</sub> DCT with  $\eta$ . Using unrelaxed molecular geometries taken from the dimer complexes, Figure 4 plots  $\eta$  versus S<sub>1</sub> DCT using TDM analysis (as in Figure 2) as well as NBO analysis. The TDM-analyzed data ( $R^2 = 0.96$ ) exhibits stronger correlation than the NBO-analyzed data ( $R^2 = 0.57$ ). The poor fit in the latter case is due to the method's inability to treat delocalized electron transfer.<sup>37</sup> Conversely, TDM analysis provides a spatial mapping of the electron—hole pair associated with an electronic transition between two states and can successfully treat such delocalized electron transfer.<sup>38,39</sup> While both methods are commonly used to calculate the charges in molecules, we recommend TDM analysis over NBO analysis.

In the left panel of Figure 5, the HOMO for DMeO-BTBT, the LUMO for TCNQ, and the HOMO and LUMO of their

Table 1.  $\eta$  Values Using the Experimental Crystal (exp) and Dimer Geometries and the Percent Deviation (%D) of  $\eta_{dimer}$ from  $\eta_{exp}$ 

system	$\eta_{ m exp}$	$\eta_{ m dimer}$	%D
4T:F <sub>2</sub> TCNQ	0.932	0.945	1.4
4T:F <sub>4</sub> TCNQ	0.917	0.931	1.5
4T:TCNQ	0.946	0.957	1.2
anthracene:F <sub>4</sub> TCNQ	0.983	0.981	0.2
DMeO-BTBT:F2TCNQ	0.975	0.980	0.5
DMeO-BTBT:F4TCNQ	0.971	0.974	0.3
DMeO-BTBT:TCNQ <sup>a</sup>	0.957	0.948	1.0
DPTTA:25F2TCNQ <sup>a</sup>	0.929	0.915	1.5
DPTTA:26F2TCNQ	0.935	0.879	6.1
DPTTA:F4TCNQ	0.853	0.828	3.0
DTBDT:F2TCNQ	0.978	0.980	0.1
DTBDT:F4TCNQ	0.972	0.975	0.3
DTBDT:TCNQ	0.980	0.981	0.1
naphthalene:TCNQ	0.997	0.975	2.1
Npe:TCNQ	0.993	0.985	0.9
STB:F4TCNQ	0.978	0.964	1.4
STB:TCNQ	0.984	0.974	1.0
tetracene:F <sub>4</sub> TCNQ	0.985	0.973	1.2
average %D			1.3

<sup>*a*</sup>Acceptor and donor molecules are perpendicular and do not interact via  $\pi - \pi$  stacking.

dimer complex are visualized; this D-A dimer has the greatest S1 DCT in the data set. The visual similarity between the monomer orbitals and those in the complex is apparent, and the localization onto the donor and acceptor moieties in the complex is striking. In the right panel of Figure 5, the HOMO of DPTTA, the LUMO of F4TCNQ, and the HOMO and LUMO of their dimer complex are visualized; this D-A dimer has the smallest S1 DCT in the data set. While there is significant visual similarity between the monomer MOs and those in the complex, there is also substantial delocalization of the HOMO onto the acceptor molecule and similar delocalization of the LUMO onto the donor. To maximize  $S_1$  DCT, the electron density must be localized on the electron donor in the HOMO and transferred completely to the acceptor LUMO. The value of  $\eta$  quantifies the degree to which the isolated HOMO (LUMO) of the donor (acceptor) correlates with that in the complex, which predicts  $S_1$  DCT.

We have shown that the  $S_1$  (HOMO  $\rightarrow$  LUMO) DCT can be predicted by a novel metric,  $\eta$ , that computes the average similarity between the donor (acceptor) molecule's HOMO (LUMO) and that of the corresponding orbital in the D-Acomplex. We found that  $\eta$  exhibits a positive linear correlation with the S<sub>1</sub> DCT for a set of 31 D-A pairs. In choosing molecular geometries to compute  $\eta$ , we found that by comparing the orbitals between isolated donor and acceptor molecules and their corresponding D-A complexes, one should use the same molecular geometries in the isolated molecules as in the D-A complex. Alternatively, experimental crystal structure data for the D-A complex can be used instead of optimized D-A dimer geometries. This allows flexibility in input data for inclusion in high-throughput screening and machine learning protocols. Lastly, we compared two methods for determining DCT, NBO analysis and TDM analysis, and found that TDM analysis is more reliable due to its ability to treat electron delocalization. Future work will aim to generalize the  $\eta$  metric to characterize the DCT of other electronic states



Figure 4. Plot of  $\eta$  versus S<sub>1</sub> DCT using NBO (top) and TDM (bottom) analyses at the D–A dimer geometries.

with the goal of uncovering further orbital structure-function relationships.

### ASSOCIATED CONTENT

#### Supporting Information

A data repository of the calculations and code used in this work can be accessed free of charge at https://doi.org/10.6084/m9.figshare.24681261. The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.jpclett.3c02787.

Plot of absolute difference in  $\eta$  along with average RMSD between monomer-optimized and dimer-optimized geometries (PDF)

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#### Notes

The authors declare the following competing financial interest(s): J.M.H. serves on the board of directors of Q-Chem Inc.

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	DMeO-BTBT:TCNQ		DPTTA:F4TCNQ		
	НОМО	LUMO	НОМО	LUMO	
Dimer	en e				
D, A	<u> <u></u></u>	Sa afer 6		<del>24005</del> 5	
DCT		0.95 e	0.57	e e	
η	0.978		0.828		

Figure 5. HOMO and LUMO for the D–A complex (top) and for the isolated molecules (bottom, donor at left and acceptor at right) using an isosurface value of 0.03 au. Associated  $S_1$  DCT and  $\eta$  values are also shown.

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