## Supporting Information for: "Reinterpreting $\pi$ -Stacking"

Kevin Carter-Fenk and John M. Herbert<sup>\*</sup> Department of Chemistry & Biochemistry, The Ohio State University

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## S1 Normalized Planarity Index (NPI)

Here we describe the NPI that was introduced in the context of Fig. 6. An MBD eigenvector  $\phi$  consists of  $3 \times N_{\text{atoms}}$  atomic displacement vectors  $\vec{\phi}_i$   $(i = 1, ..., N_{\text{atoms}})$ , where each  $\vec{\phi}_i$  is an atomcentered vector in three-dimensional space. Let  $\vec{u}$  and  $\vec{v}$  denote vectors that define the plane of interest, with

$$\vec{n} = \|\vec{u} \times \vec{v}\| \tag{S1}$$

being a vector that is perpendicular to the  $\vec{u} \cdot \vec{v}$  plane. The definition of this plane is arbitrary but the choice of a plane that bisects the center-to-center vector of the two monomers is convenient and intuitive. Mapping of the MBD eigenvectors onto the  $\vec{u} \cdot \vec{v}$  plane is accomplished by subtracting the components of each eigenvector that lie orthogonal to the plane, to obtain new vectors

$$\vec{p}_i = \vec{\phi}_i - (\vec{n} \cdot \vec{\phi}_i) \, \vec{n} \,. \tag{S2}$$

The norm  $\|\vec{p}_i\| \in [0, 1]$ , with the value zero suggesting that the original eigenvector  $\vec{\phi}_i$  is orthogonal to the reference plane and  $\|\vec{p}_i\| = 1$  indicating that it is parallel to the reference plane. The total magnitude P for the the projected eigenvectors is a sum over the magnitudes of the projected eigenvectors that describe each oscillator,

$$P = \sum_{i}^{N_{\text{atoms}}} \left(\sum_{\alpha=x,y,z} p_{i\alpha}^2\right)^{1/2}.$$
(S3)

The NPI is taken to be the magnitude of the projected eigenvectors normalized to the number of atoms,

$$NPI = P/N_{atoms} .$$
 (S4)

## S2 Additional Data

<sup>\*</sup>herbert@chemistry.ohio-state.edu

Molecule	$\omega / a_0^{-1}$
benzene	0.340
naphthalene	0.303
anthracene	0.288
tetracene	0.278
pentacene	0.268
cyclohexane	0.321
perhydronaphthalene	0.300
perhydroanthracene	0.287
perhydrotetracene	0.275
perhydropentacene	0.264
$C_{96}H_{24}$	0.230
corannulene	0.277

Table S1: Tuned range separation parameters computed using  $LRC-\omega PBE/def2-TZVPPD$ .



Figure S1: Two-dimensional potential surfaces for cofacial sliding of benzene along the surface of a  $C_{96}H_{24}$ graphene nanoflake, using (a) the Hunter-Sanders electrostatic + dispersion model potential (Ref. 1) and (b) the vdW model potential from c The centers of the two molecules are aligned at the coordinate origin, which is a local minimum in (a) but a saddle point in (b). Note that the absolute energy scale for the vdW model is not meaningful, as the very simple parameterization in Ref. 2 does not include a short-range damping function.



Figure S2: One-dimensional cross sections of the two-dimensional surfaces in Fig. S1, further elucidating the topography near the minimum. (a) The Hunter-Sanders model places benzene at the center of  $C_{96}H_{24}$ , whereas (b) the vdW model predicts that this is geometry is a saddle point, and instead predicts a slightly off-center structure as the minimum. Note that the absolute energy scale for the vdW model is not meaningful, as the very simple parameterization in Ref. 2 does not include a short-range damping function.



Figure S3: HOMO/LUMO gaps for linear acenes and their saturated analogues, computed at the LRC-  $\omega \rm PBE/def2\text{-}TZVPPD$  level.



Figure S4: Two-dimensional scan of  $E_{\rm elst+ind}$  for (a) naphthalene dimer and (b) perhydronaphthalene dimer. The two-dimensional contours illustrate the topography of the potential surface and the cofacial sliding coordinates of the acene dimer and the analogues of the same displacements for perhydronaphthalene dimer.

## References

- [1] Hunter, C. A.; Sanders, J. K. M. The nature of  $\pi$ - $\pi$  interactions. J. Am. Chem. Soc. 1990, 112, 5525–5534.
- [2] Carter-Fenk, K.; Herbert, J. M. Electrostatics does not dictate the slip-stacked arrangement of aromatic  $\pi$ - $\pi$  interactions. *Chem. Sci.* **2020**, *11*, 6758–6765.