

The CKM Matrix

- We need to connect the weak eigenstates of quarks with the mass eigenstates

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

\uparrow \uparrow \uparrow

weak eigenstates V_{CKM} mass eigenstates

◆ If ν 's have mass there will be analogous matrix

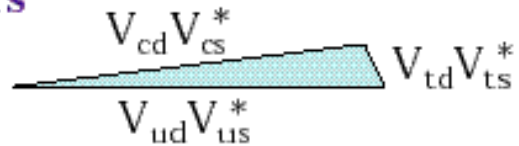
Proper Formulation of CKM Matrix

$$V = \begin{array}{c} \mathbf{u} \\ \mathbf{c} \\ \mathbf{t} \end{array} \begin{array}{ccc} \mathbf{d} & \mathbf{s} & \mathbf{b} \\ \left(1 - \frac{1}{2}\lambda^2 \right) & \lambda & A\lambda^3 \left(\rho - i\eta \left(1 - \frac{1}{2}\lambda^2 \right) \right) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 - i\eta A^2 \lambda^4 & A\lambda^2 (1 + i\eta \lambda^2) \\ A\lambda^2 (1 - \rho - i\eta) & -A\lambda^2 & 1 \end{array}$$

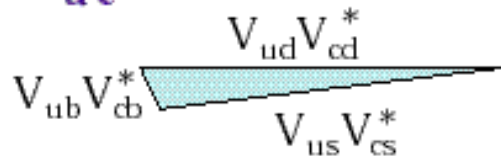
- **Good λ^3 in real part & λ^5 in imaginary part**
- **We know $\lambda = 0.22$, $A \sim 0.8$; constraints on ρ & η**
- **Due unitarity there are 6 CKM triangles**

The 6 CKM triangles

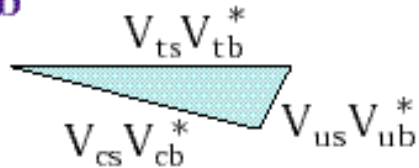
ds



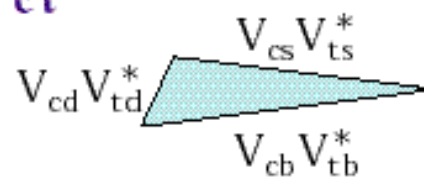
uc



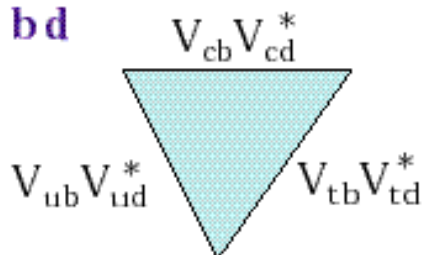
sb



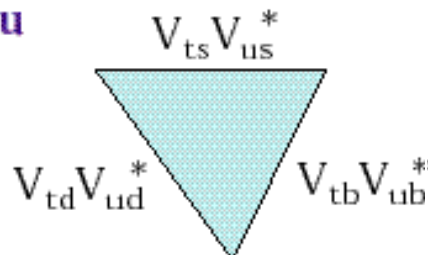
ct



bd



tu



- “ds” - indicates rows or columns used
- There are 4 independent phases, which can be used to construct entire CKM matrix

The 4 CKM Phases

$$\beta = \arg\left(-\frac{V_{tb} V_{td}^*}{V_{cb} V_{cd}^*}\right) \quad \gamma = \arg\left(-\frac{V_{ub}^* V_{ud}}{V_{cb}^* V_{cd}}\right)$$
$$\chi = \arg\left(-\frac{V_{cs}^* V_{cb}}{V_{ts}^* V_{tb}}\right) \quad \chi' = \arg\left(-\frac{V_{ud}^* V_{us}}{V_{cd}^* V_{cs}}\right)$$

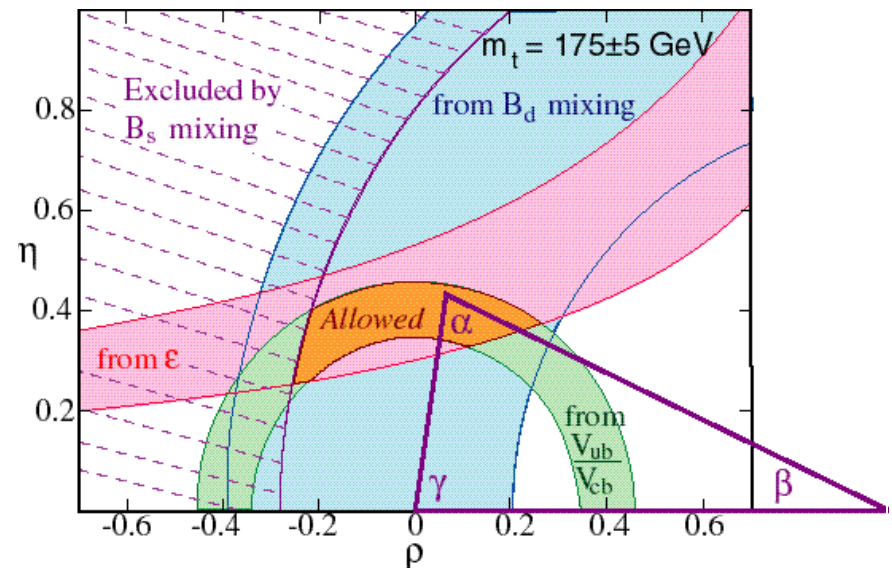
β & γ probably large, χ small ~ 0.02 , χ' smaller



Derive reference triangle

Usual Triangle (for ref.)

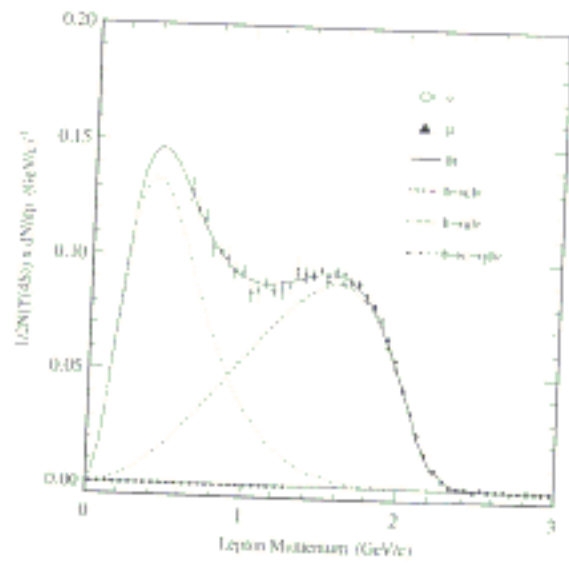
- Real a , b & g must sum to 180°
- Therefore, any two will do IF we are really measuring the intrinsic angles
- New physics can hide, only these angle measurements not sufficient.
- Ex: Suppose there is new physics in B^0 - B^0 mixing (q) & we measure CP in γK_S and p^+p^- , then $2b\phi = 2b+q$, $2a\phi = 2a-q$, & $2b\phi + 2a\phi = 2b+2a$, new physics but $b\phi + a\phi + g = 180^\circ$



◆ Two sides $|V_{ub}/V_{cb}|$ & $|V_{td}/V_{ts}|$ important.

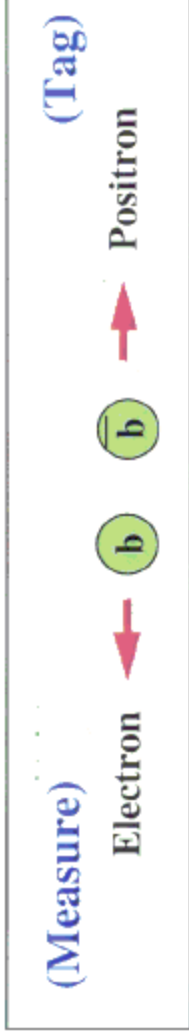
Ambiguities

- Suppose we measure $\sin(2b)$ using yK_s , what does that tell us about b ?
- **Ans: 4 fold ambiguity- b , $p/2 - b$, $p + b$, $3p/2 - b$**
- Only reason $h > 0$, is $B_k > 0$ from theory, and related theoretical interpretation of $e\dot{c}$



toward a model independent measurement

Trick: b quarks are produced in pairs

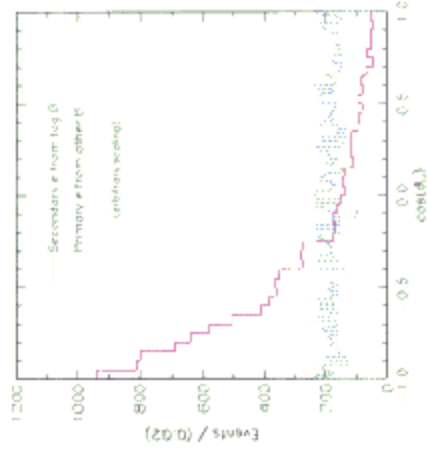


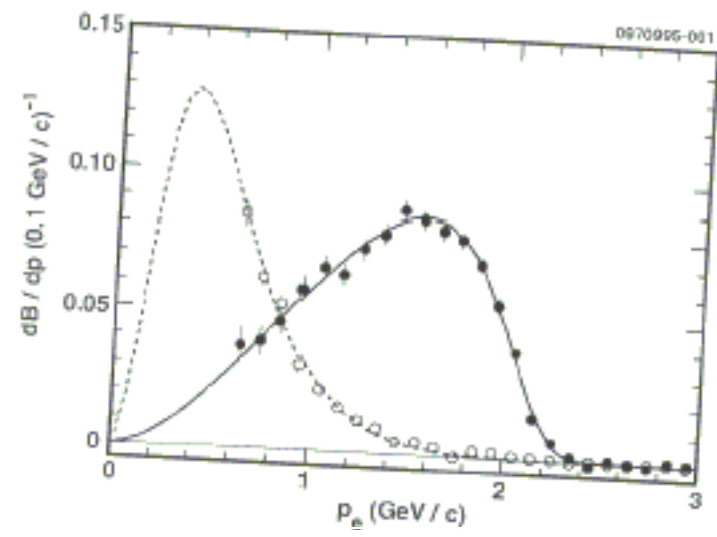
But



... there is much more to this analysis

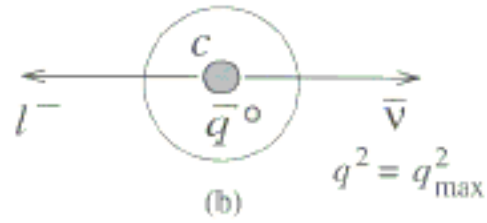
e.g. cuts to
suppress more
background



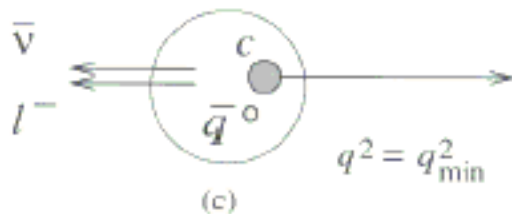




(a)



(b)



(c)

Heavy Quark Effective Theory

- HQET tells us that in first order when a b quark transforms to a c quark with the c going at the same velocity as the b, the form factor is 1 in first order **AND** the corrections to 1 can be calculated
- The form-factor therefore known to be 1-correction, at maximum q^2 , called $w=1$, where

$$\omega = \frac{M_B^2 + M_{D^*}^2 - q^2}{2M_B M_{D^*}}$$

Heavy Quark Effective Theory

B meson



D meson



$$m_Q \rightarrow \infty$$

(Heavy) flavor symmetry
spin symmetry

Universal
Formfactor



Absolute Normalization
at zero recoil

$$\xi(1) = 1$$

$$m_Q < \infty$$

$$\xi(1) = \eta_A (1 + \epsilon_1 \Lambda_{\text{QCD}}/m_Q + \epsilon_2 (\Lambda_{\text{QCD}}/m_Q)^2 \dots) = \eta_A (1 + \delta)$$

for $B \rightarrow D^* l \nu$
(Luke)

$$\xi(1) = 0.91 \pm 0.03 \text{ (Neubert, Shifman ...)}$$

Experiments

$|V_{cb}|$ from $B \rightarrow D^* \ell \bar{\nu}$

- Use $B \rightarrow D^* \ell \bar{\nu}$ because the decay rate is largest for and the corrections are better determined.
- In HQET there is one “universal” form-factor function, so we don’t have to deal with 3 form-factors
- To find V_{cb} measure value at $w=1$, here D^* is at rest in B rest frame

$$\frac{d\Gamma(B \rightarrow D^* \ell \bar{\nu})}{d\omega} = \frac{G_F^2}{48\pi^3} |V_{cb}|^2 F^2(\omega) (m_B - m_{D^*})^2 m_{D^*}^3 \sqrt{\omega^2 - 1} \times \left[4\omega(\omega+1) \frac{1 - 2\omega \frac{m_{D^*}}{m_B} + \frac{m_{D^*}^2}{m_B^2}}{\left(1 - \frac{m_{D^*}}{m_B}\right)^2} \right]$$

V_{cb} results, an example

- To get results fit using shape proposed by Caprini et al, or Boyd & Grinstein, or in CLEO case by Stone
- Use $F(1)=0.91\pm 0.03$, from Caprini, Uraltsev.....
- **Results**
 - DELPHI: $(41.2\pm 1.5\pm 1.8\pm 1.4)\times 10^{-3}$
 - ALEPH: $(34.4\pm 1.6\pm 2.3\pm 1.4)\times 10^{-3}$
 - OPAL: $(36.0\pm 2.1\pm 2.1\pm 1.2)\times 10^{-3}$
 - CLEO: $(39.4\pm 2.1\pm 2.2\pm 1.3)\times 10^{-3}$
 - **World Average 0.0381 ± 0.0021 by adding theoretical error in quadrature with exp error.**

Theoretical Value of $F(1)$

- $\lim F(1) = 1$ as $m_b \rightarrow \infty$,
- $F(1) = 1 + O(a_s/p) + d_{1/m^2} + d_{1/m^3}$ (no $d_{1/m}$, Lukes thrm)
 - $F(1) = 0.91 \pm 0.03$, from Caprini, Uraltsev.....
 - $F(1) = 0.89 \pm 0.06$, from Bigi
 - Can we get an accurate non-quenched value from the Lattice?
 - The errors are not consistent. What do the errors mean?
- **Bigi:** "In stating a theoretical error, I mean that the real value can lie almost anywhere in this range with basically equal probability rather than follow a Gaussian distribution. Furthermore, my message is that I would be quite surprised if the real value would fall outside this range. Maybe one could call that a 90% confidence level, but I do not see any way to be more quantitative."

QCD Sum Rules for $|V_{cb}|$

- Using Operator Product Expansion & Heavy Quark Expansion, in terms of $a_s(m_b)$, Λ , and the matrix elements l_1 and l_2 , we can accurately determine V_{cb} .

- These quantities arise from the differences

$$m_B - m_b = \bar{\Lambda} - \frac{\lambda_1 + 3\lambda_2}{2m_b} + K, \quad m_{B^*} - m_b = \bar{\Lambda} - \frac{\lambda_1 - \lambda_2}{2m_b} + K,$$

- From $B^* - B$ mass difference, $l_2 = 0.12 \text{ GeV}^2$

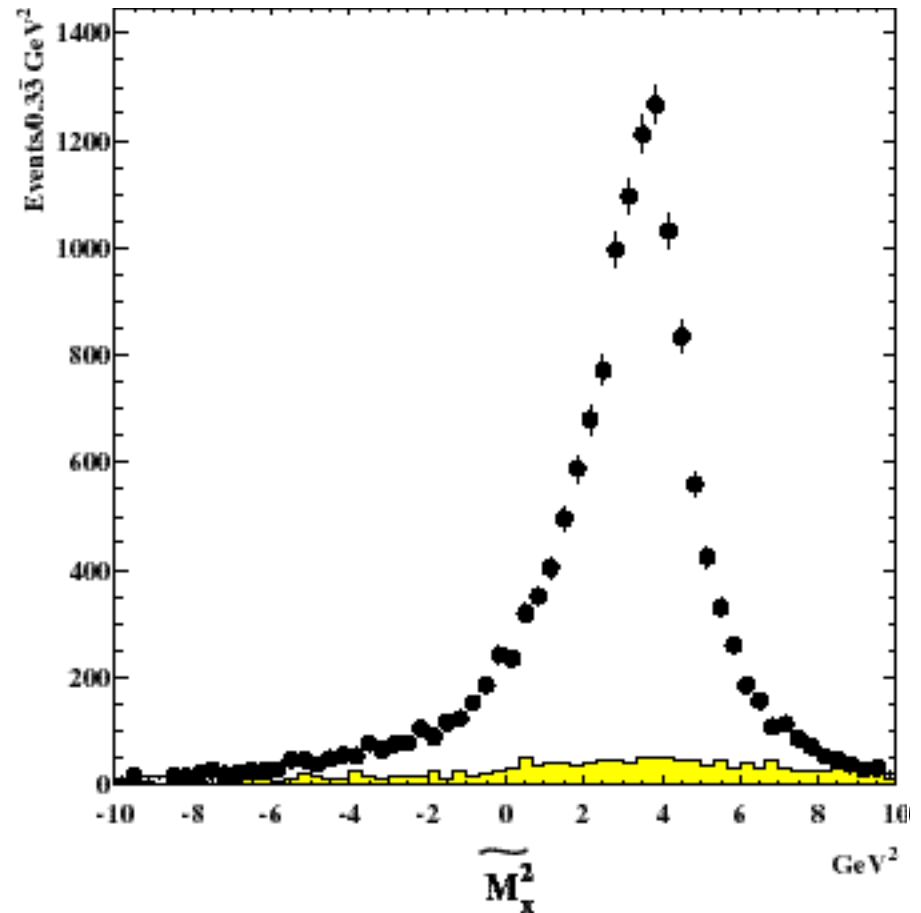
- $$\Gamma_{sl} = \frac{G_F^2 |V_{cb}|^2 m_B^5}{192\pi^3} 0.369 \left[\begin{array}{l} 1 - 1.54 \frac{\alpha_s}{\pi} - 1.65 \frac{\bar{\Lambda}}{m_B} \left(1 - 0.87 \frac{\alpha_s}{\pi} \right) \\ -0.95 \frac{\bar{\Lambda}^2}{m_B^2} - 3.18 \frac{\lambda_1}{m_B^2} + 0.02 \frac{\lambda_2}{m_B^2} \end{array} \right]$$

Measurement of G_{sl} & moment analysis

- Use total Branching Ratio Measurement

- CLEO using lepton tags $(10.49 \pm 0.17 \pm 0.43)\%$
- Lifetime $1.613 \pm 0.020 \text{ ps}$ $\Gamma_{sl} = 65.0 \pm 3.0 \text{ ns}^{-1}$
- (note LEP $68.6 \pm 1.6 \text{ ns}^{-1}$)

- “Moment Analysis” of $B \rightarrow X l n$, M_x and E_l



Result for $|V_{cb}|$ Using Moment Analysis

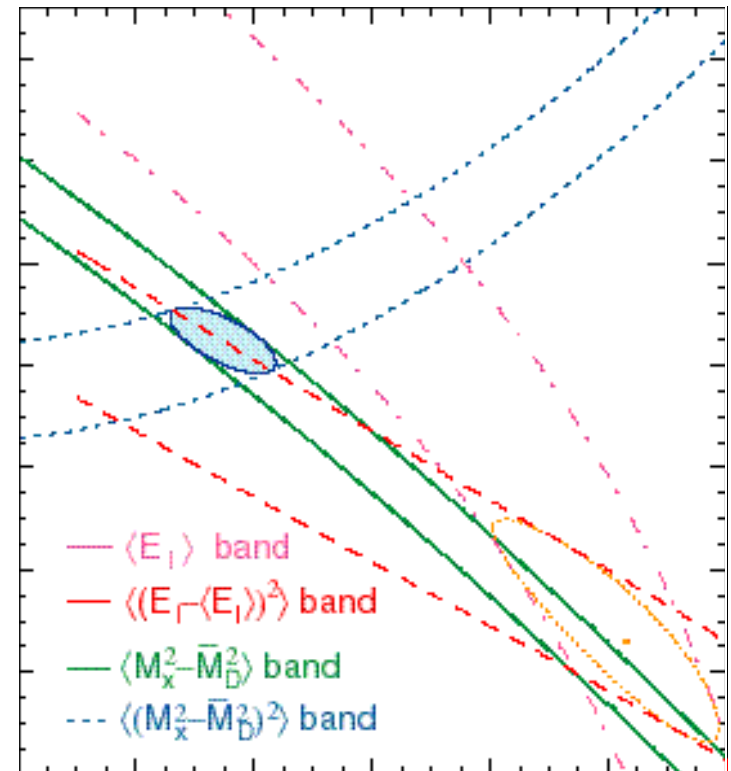
- **Discrepancy between hadronic mass moments and E_1 moments**
- **Theoretical estimates favor M_x moments**
- **Taking M_x estimates only:**

$$\Lambda = 0.33 \pm 0.02 \pm 0.08 \text{ GeV}$$

$$\lambda_1 = -0.13 \pm 0.01 \pm 0.06 \text{ GeV}^2$$

- **Ligeti claims:**

$|V_{cb}| = 0.0415 \pm 0.0012$, but
it would be foolish to use it



*Is this an experimental problem
or an inherent problem in OPE?*

A new measurement of V_{cb}

The differential decay rate, for $B \rightarrow D^* \ell \nu$ is:

$$\frac{d\Gamma}{dw} = \frac{G_F^2}{48\pi^3} |V_{cb}|^2 |F(w)|^2 G(w)$$

w is the Lorentz γ factor of the recoiling D^* .

It is a function of masses and q^2 . ($1 \leq w \leq 1.5$)

$w = 1$ is the "zero recoil" point.

$G(w)$ is a known kinematic function.

$F(w)$ is the form factor. HQET constrains it.

As $m_{b,c} \rightarrow \infty$, $F(1) \rightarrow 1$.

For finite m , the corrections are $O(1/m^2)$.

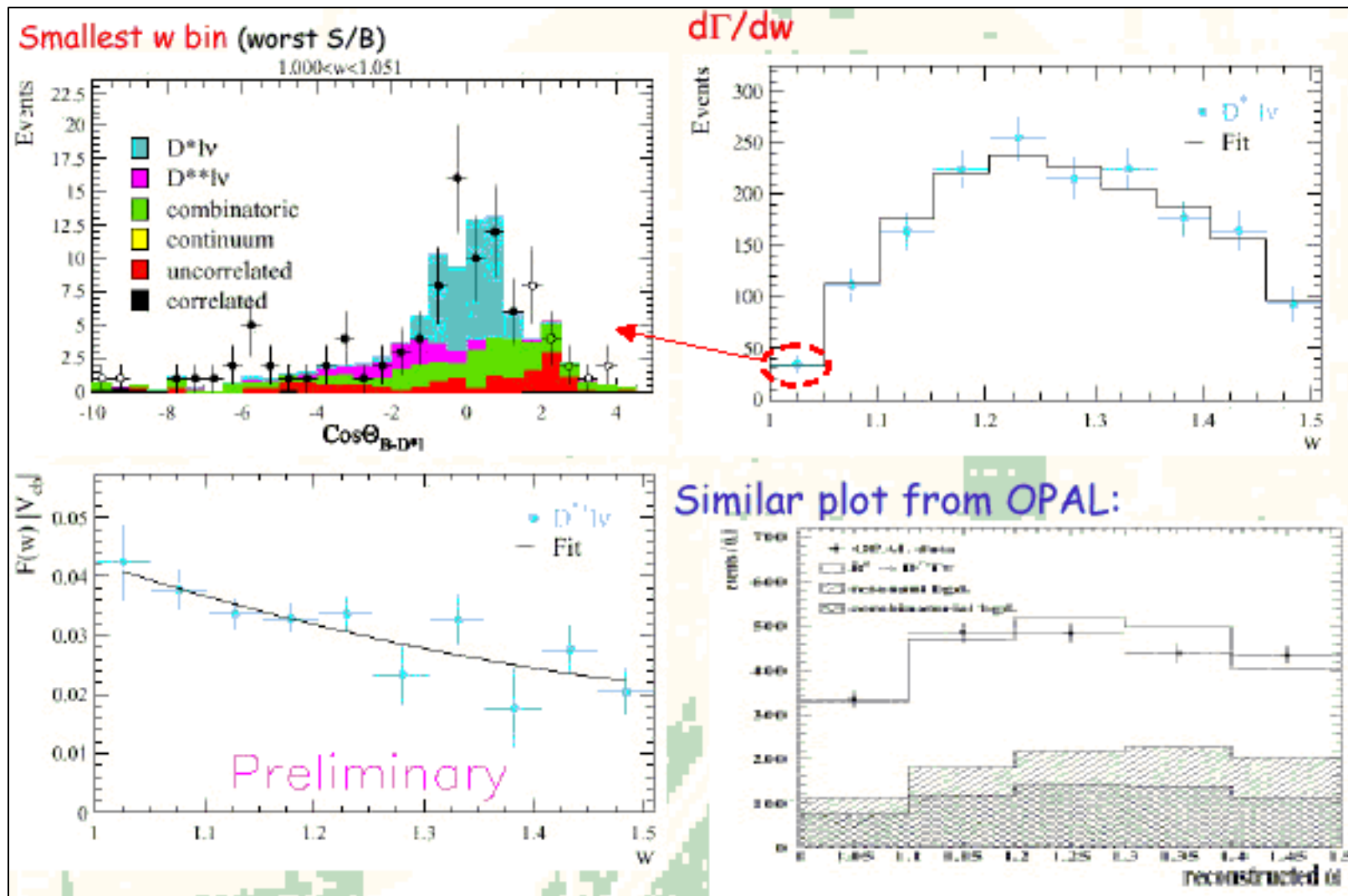
Two methods:

✖ Extrapolate differential rate to $w = 1$.

✖ Integrate the total $B(b \rightarrow X_c \ell \nu)$.

The first method is less sensitive to $F(w)$ systematics, but has worse statistics.

Results



Is there a problem?

$$F(1)|V_{cb}| = (42.4 \pm 1.8 \pm 1.9) \times 10^{-3}$$

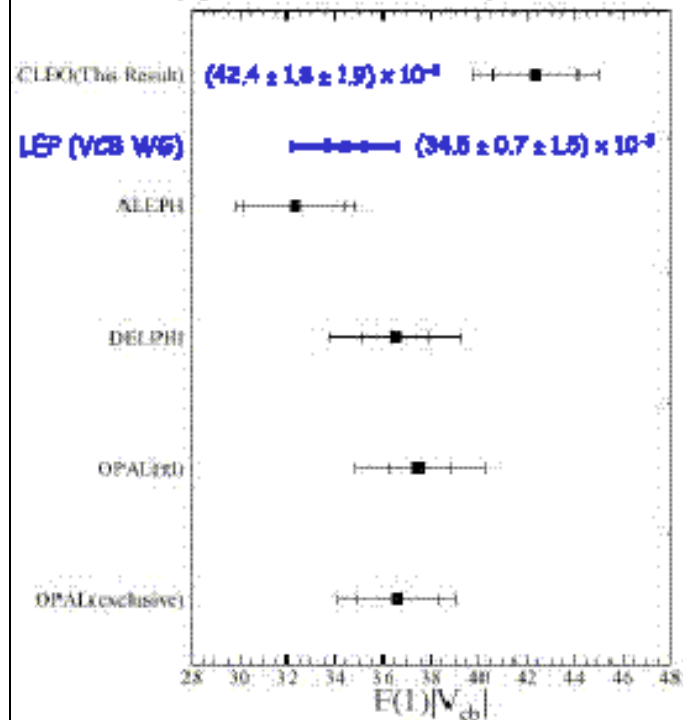
$$\text{BR}(D^{*+} \ln) = (5.66 \pm 0.29 \pm 0.33)\%$$

using: $F(1) = 0.913 \pm 0.042$

yields: $|V_{cb}| = (46.4 \pm 2.0 \pm 2.1 \pm 2.1) \times 10^{-3}$

↑
F(1) uncertainty

There appears to be a problem:



This result is strongly correlated with the measured slope, r^2 , of $F(w)$:

CLEO: $r^2 = 1.67 \pm 0.11 \pm 0.22$

LEP: $r^2 = 1.01 \pm 0.08 \pm 0.16$

As we saw, LEP and CLEO now agree about the total semileptonic BR.

Preliminary:

Belle: $\text{BR}(D^{*+} \ln) = (4.74 \pm 0.25 \pm 0.51)\%$

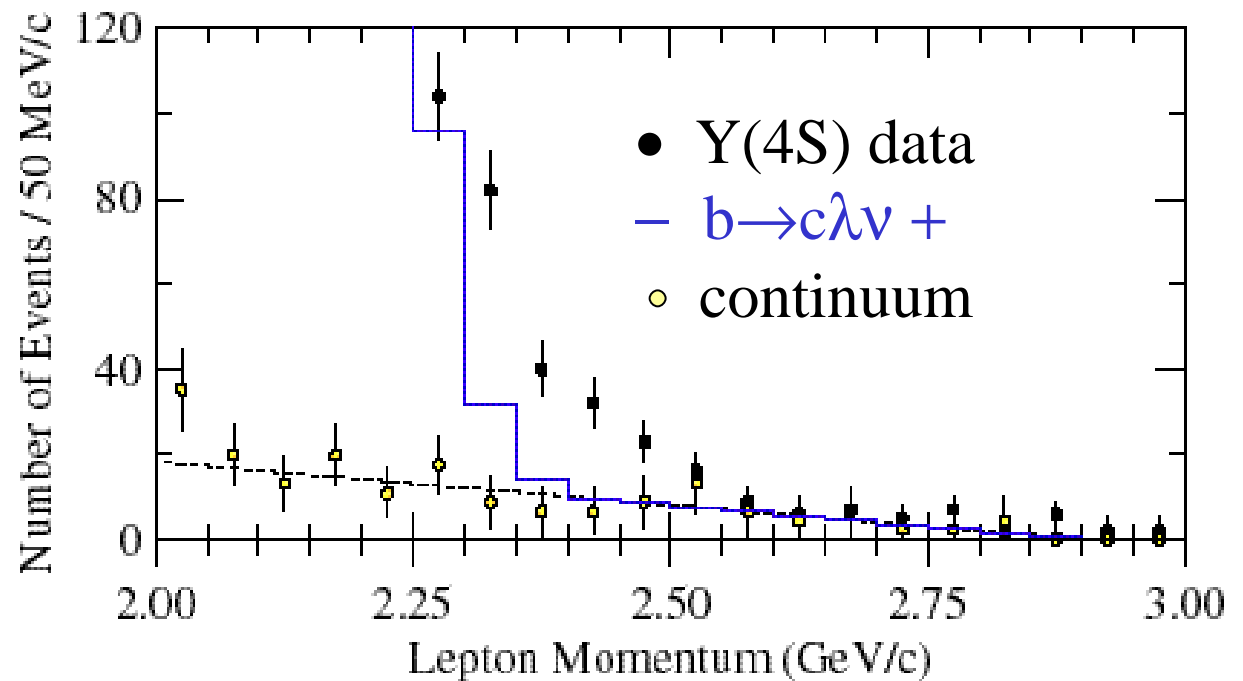
$\text{BR}(D \ln) = (2.07 \pm 0.21 \pm 0.31)\%$

Waiting for a V_{ub} update...

- Exclusive ($B \rightarrow \ell \ln$ or $p \ln$) [CLEO]
 $|V_{ub}| = (3.25 \pm 0.14^{+0.21}_{-0.29} \pm 0.55) \times 10^{-3}$
- Inclusive ($b \rightarrow X_u \ln$) [LEP, CLEO]
 $|V_{ub}| / |V_{cb}| = 0.8 \pm 0.2$
- Rare hadronic decays (e.g. $B \rightarrow p \bar{p}$) (seen)
- $B \rightarrow D_s^{(*)} p^{+(0)}$ (not seen)
- $B \rightarrow t \bar{n}$

V_{ub} from lepton endpoint

- BR & value of V_{ub} depends on model, since fraction of leptons in signal region depends on model!



Neutrino Reconstruction



Resolution

$$\sigma_{\text{Emiss}} \sim 260 \text{ MeV} \quad \sigma_{|\vec{p}_{\text{miss}}|} \sim 110 \text{ MeV}$$

σ_{MM}^2 dominated by $2E_{\text{Emiss}}\sigma_{\text{Emiss}}$ term

Require: $\text{MM}^2 / 2E_{\text{miss}} < 300 \text{ MeV}$

→ reconstructed $\nu = (|\vec{p}_{\text{miss}}|, \vec{p}_{\text{miss}})$

traditional full B reconstruction

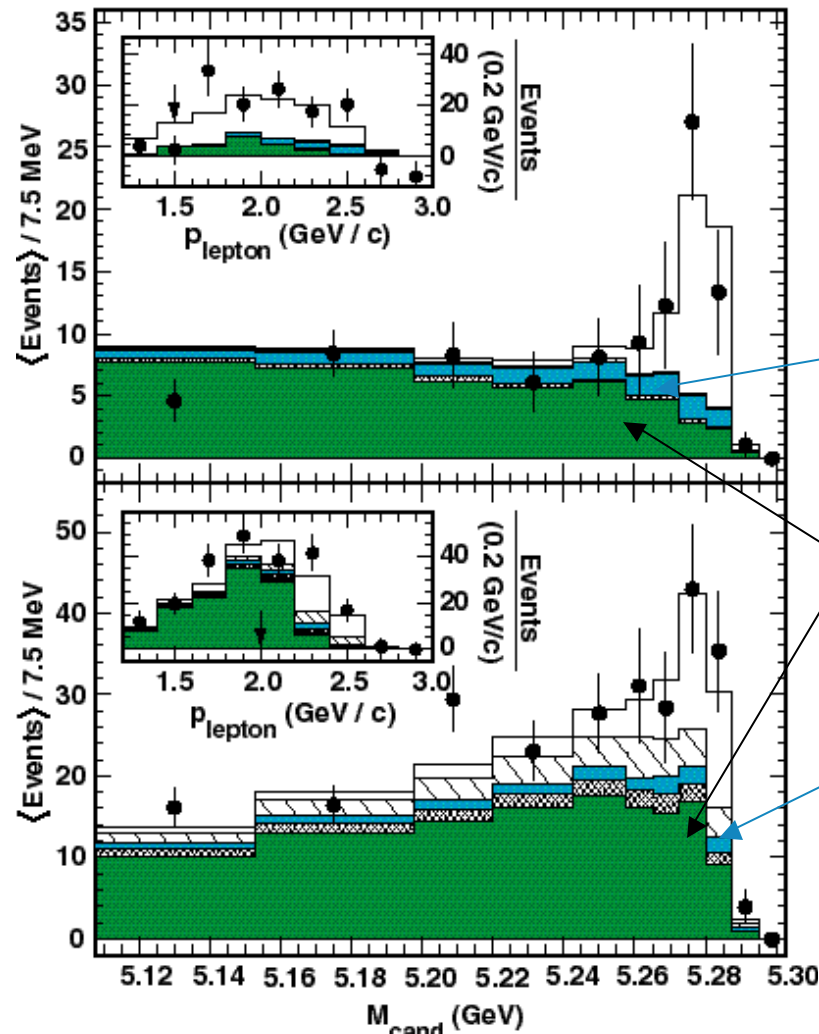
- lepton momentum $p_l > 1.5 \text{ GeV}/c$
- $\Delta E = E_{\text{beam}} - (E_\nu + E_l + E_\pi)$
- $M = \sqrt{E_{\text{beam}}^2 - (\vec{p}_\nu + \vec{p}_l + \vec{p}_\pi)^2}$

Efficiency (ISGW) = 2.9%

V_{ub} from $p l n$ and $r l n$

$p^+ l n + p^0 l n$

$r^+ l n + r^0 l n + w^0 l n$



$b \rightarrow u$ backgrounds & cross-feeds

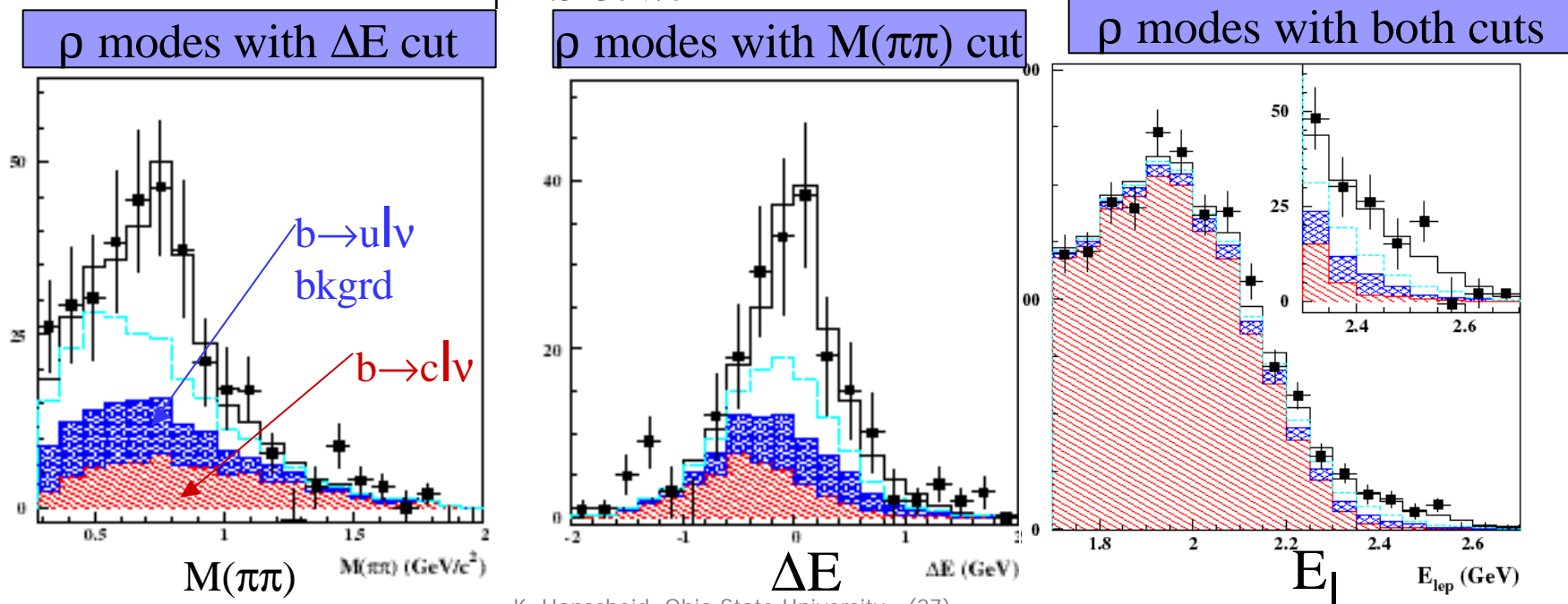
$b \rightarrow c$ backgrounds

$b \rightarrow u$ backgrounds & cross-feeds

New CLEO form-factor Analysis

Find $\rho^+\lambda\nu + \rho^0\lambda\nu + \omega^0\lambda\nu$ as function of E_λ using likelihood method to fit $M(\pi\pi)$ & ΔE distributions, where $\Delta E = E_\rho + E_l + |\mathbf{p}_{\text{miss}}| - E_{\text{beam}}$

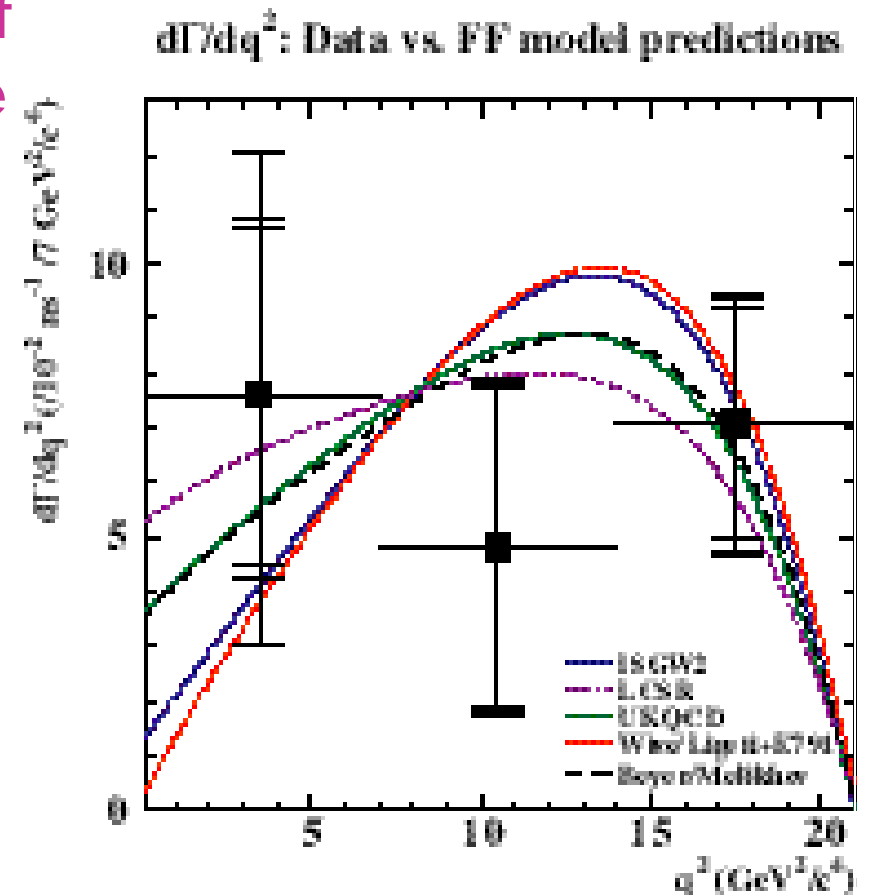
$E_l > 2.3 \text{ GeV}/c$



Form-factor Results

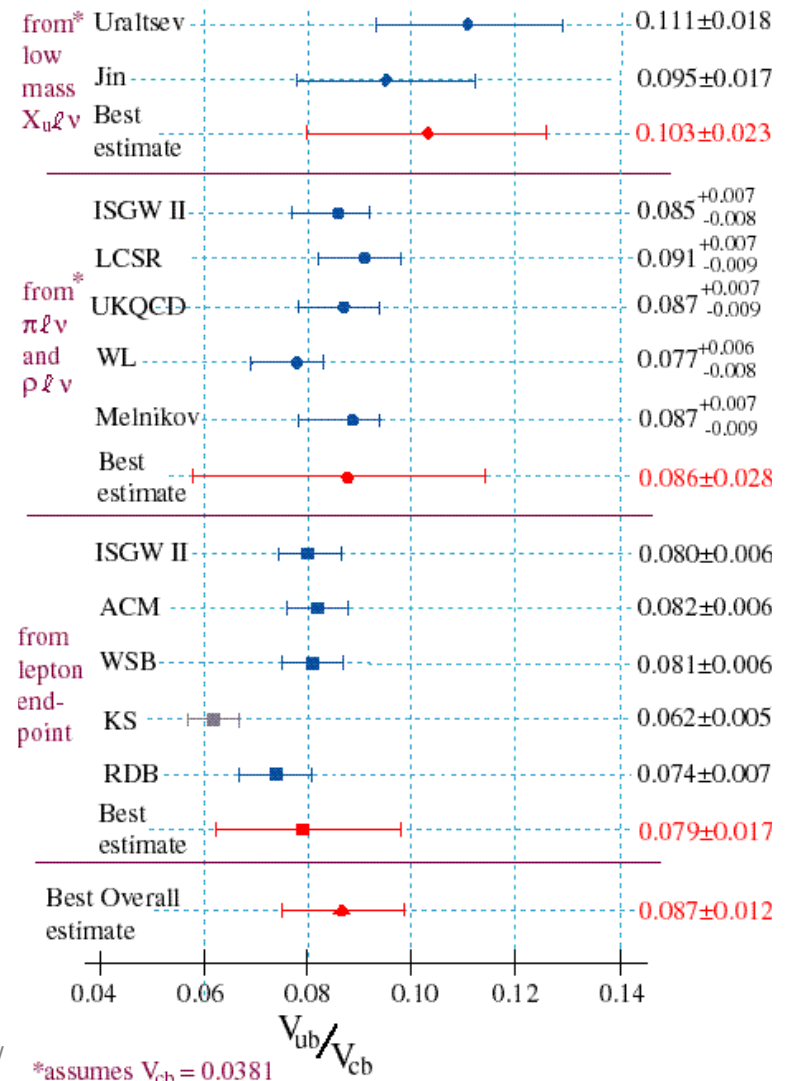
- In general 3 form-factors f_0^- @ 1^- transitions, but we do not have enough precision to disentangle them
- Data shows the need for more data
- Combining with older result:

$$|V_{ub}| = (3.25 \pm 0.14^{+0.21}_{-0.29} \pm 0.55) \times 10^{-3}$$



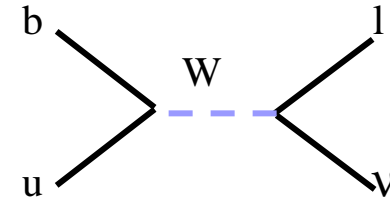
Summary of $|V_{ub}|$ Results

- LEP measurements use 8% theoretical error as given by Uraltsev. However Jin's similar calculation claims a 10% error but differs by 14%. I use a 14% theory error here
- Since the LEP Monte-Carlo calculations are highly correlated, I take a common 14% systematic uncertainty
- The exclusive channels rule out the Korner & Schuler (KS) model (gets the wrong V/P ratio), but have large errors
- The CLEO endpoint results have the best statistical error. Hard to estimate the theoretical error. I take 14%.



Pure leptonic decays: $B \rightarrow \tau \nu$

$$BR(B^+ \rightarrow l^+ \nu) = \frac{G_F^2 m_B m_l^2}{8\pi} |V_{ub}|^2 f_B^2 \tau_B$$



Standard Model: $B \rightarrow \tau \nu \sim (0.2-1) \times 10^{-4}$

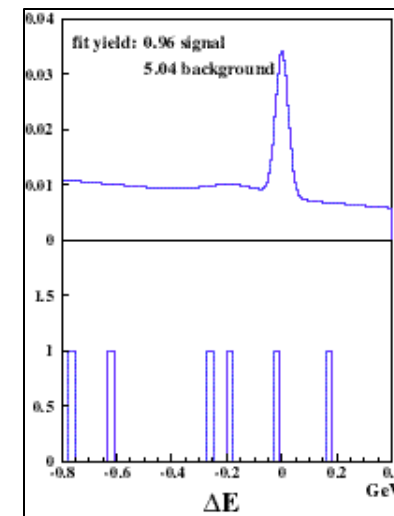
Idee: fully reconstruct on B, look for τ + missing energy

Using 9.7×10^6 BB events we find

$$BR(B \rightarrow \tau \nu) < 8.4 \times 10^{-4}$$

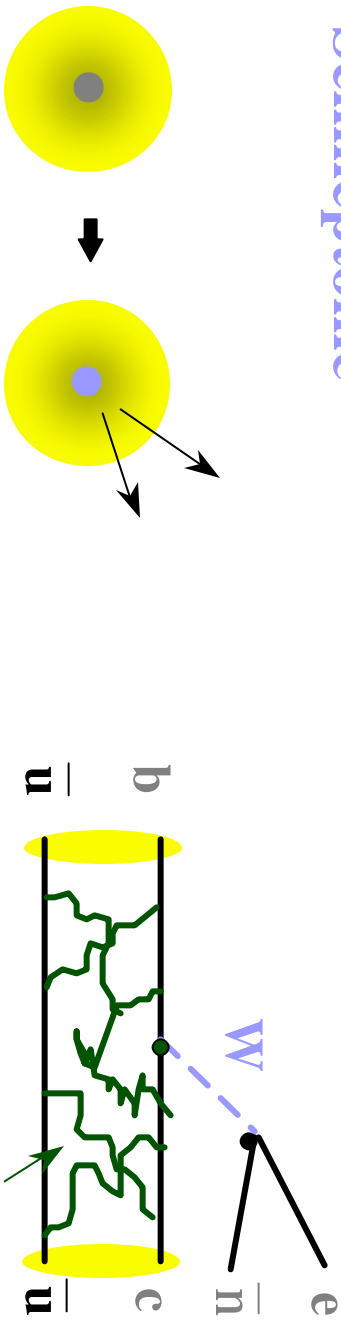
Bonus:

$$BR(B \rightarrow K \nu \nu) < 2.4 \times 10^{-4}$$



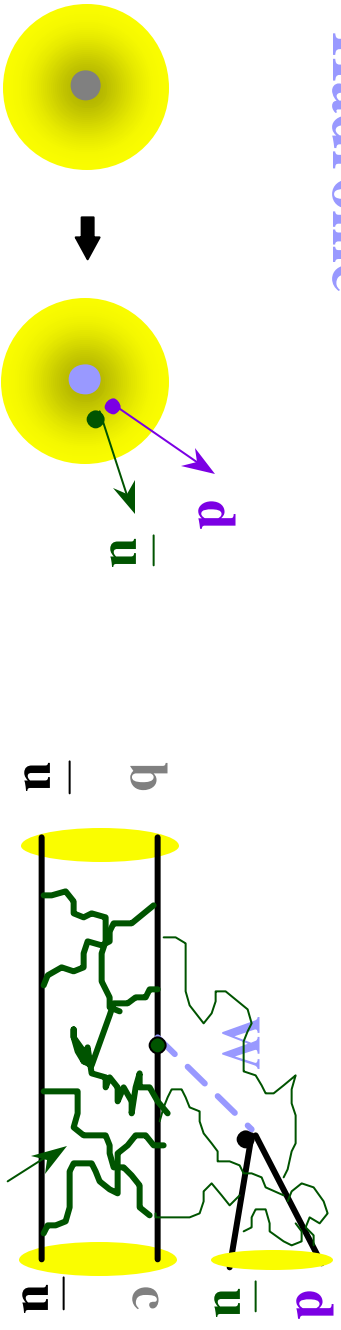
Understanding Hadronic Decays

— Semileptonic



Strong Interaction

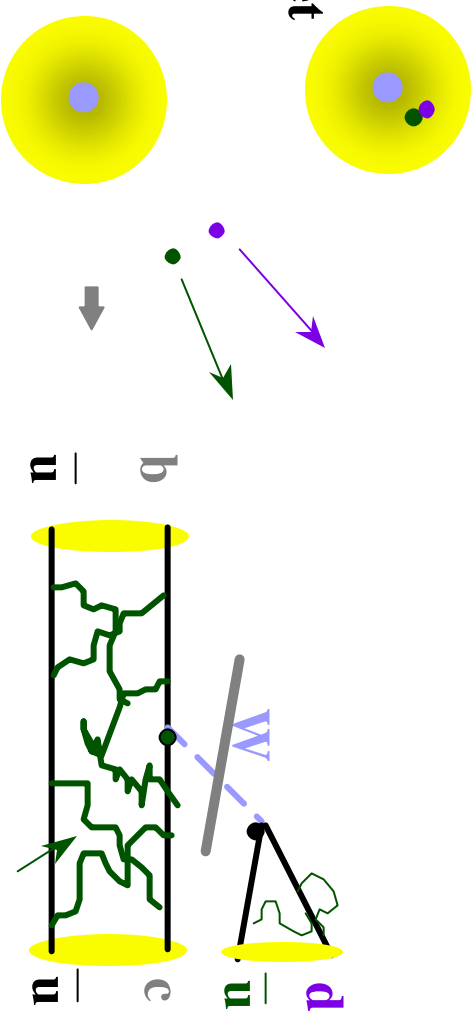
— Hadronic



Strong Interaction

But if ...

W creates $u\bar{d}$
pointlike
→ color singlet



Strong Interaction

if they get out
fast enough ...

Hadronic B Decays

- **Semileptonic Decay** $A = \frac{G_F}{\sqrt{2}} V_{cb} V_{ub}^* \langle n | g_m (1 - g_5) | l \rangle \langle D^{*-} | (cb) | B^0 \rangle$
- **Hadronic + Factorization** $A = \frac{G_F}{\sqrt{2}} V_{cb} V_{ub}^* \langle p | (du) | 0 \rangle \langle D^{*-} | (cb) | B^0 \rangle$

Factorization Tests:

- **Branching Ratios**

$$\frac{G(B \otimes D^{*+} h^-)}{\frac{dG}{dq^2}(B \otimes D^{*+} ln) |_{q^2=m_h^2}} = 6p^2 c_1^2 f_h^2 |V_{ud}|^2$$

- **Polarization**

$$G_L/G(B \otimes D^{*+} h^-) = G_L/G(B \otimes D^{*+} ln) |_{q^2=m_h^2}$$

Hadronic Decays and Factorization

(1997)

Semileptonic (e.g. $B \rightarrow D^* l^+ \nu$)

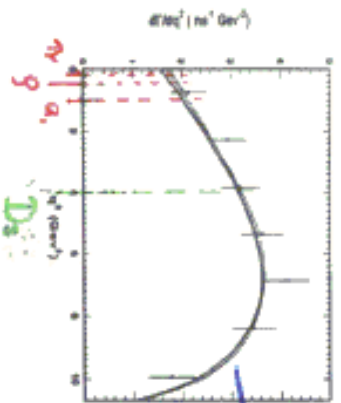
$$A = G_F/1.4 V_{cb} \langle \nu | \gamma_\mu (1 - \gamma_5) | l \rangle \langle D^* | | B^0 \rangle$$

Hadronic (+Factorization) (e.g. $B \rightarrow D^* \bar{s}^+$)

$$A = G_F/1.4 V_{cb} \langle \pi | (du) | 0 \rangle \langle D^* | | B^0 \rangle$$

I. Branching Ratio Tests

Input: π decay constant \checkmark
semileptonic decay rate \checkmark



$$\frac{\Gamma(B^0 \rightarrow D^* \pi^+)}{q^2 \int_{q^2_{min}}^{q^2_{max}} d\Gamma/dq^2(B^0 \rightarrow D^* l^+ \nu)} = 6 \pi^2 c_1^2 F_\pi^2 |V_{ud}|^2$$

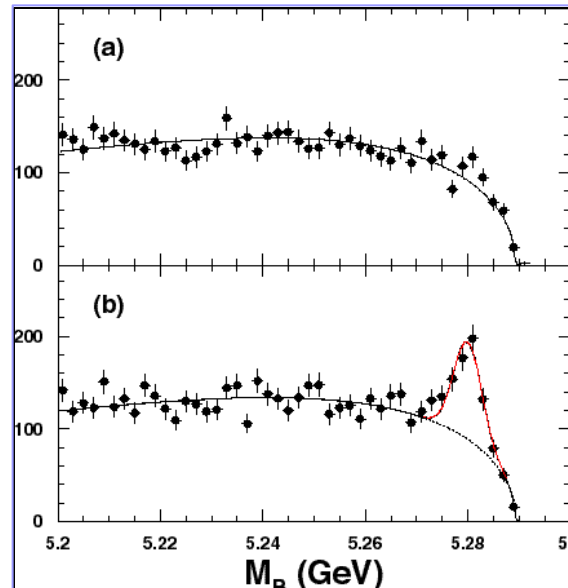
	$B \rightarrow D^*$	$B \rightarrow D$	“Theory”
$B^0 \rightarrow D^{(*)} \pi^+$	1.18 +/- 0.21	0.94 +/- 0.30	1.22 +/- 0.15
$B^0 \rightarrow D^{(*)} \rho^+$	2.92 +/- 0.70	2.63 +/- 0.88	3.26 +/- 0.42
$B^0 \rightarrow D^{(*)} a_1^+$	3.8 +/- 1.0		3.0 +/- 0.5

II. Polarization Tests

$$\Gamma_L / \Gamma(B^0 \rightarrow D^* \rho^+) = 90 +/- 7 +/- 5\% \rightarrow \Gamma_L / \Gamma(B^0 \rightarrow D^* l^+ \nu) = 88\%$$

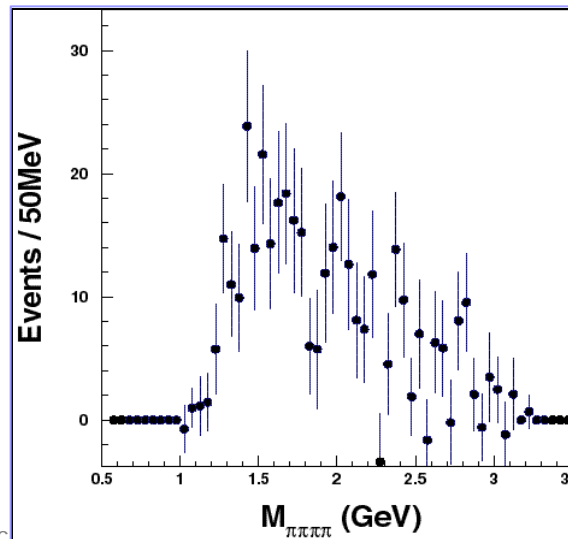
The $D^{*+}p^+p^-p^-p^0$ Final State

- (a) DE sidebands
| 3.0 – 5.0 s |
- (b) DE around 0 \pm 2.0s fit with sideband shape fixed & norm allowed to float
- Also signals in $D^0 \textcircled{R} K^- p^+ p^0$ and $D^0 \textcircled{R} K^- p^+ p^+ p^-$ (not shown)
- Fit B yield in bins of $M(4p)$



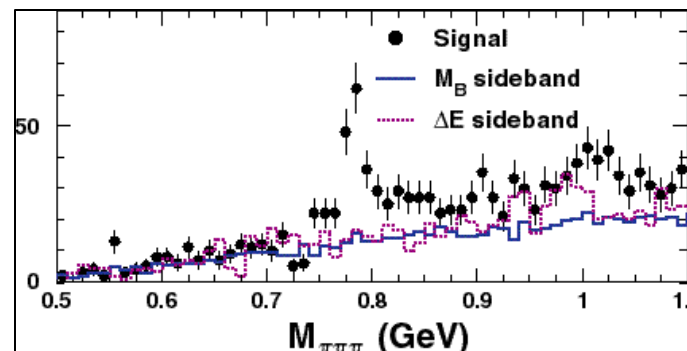
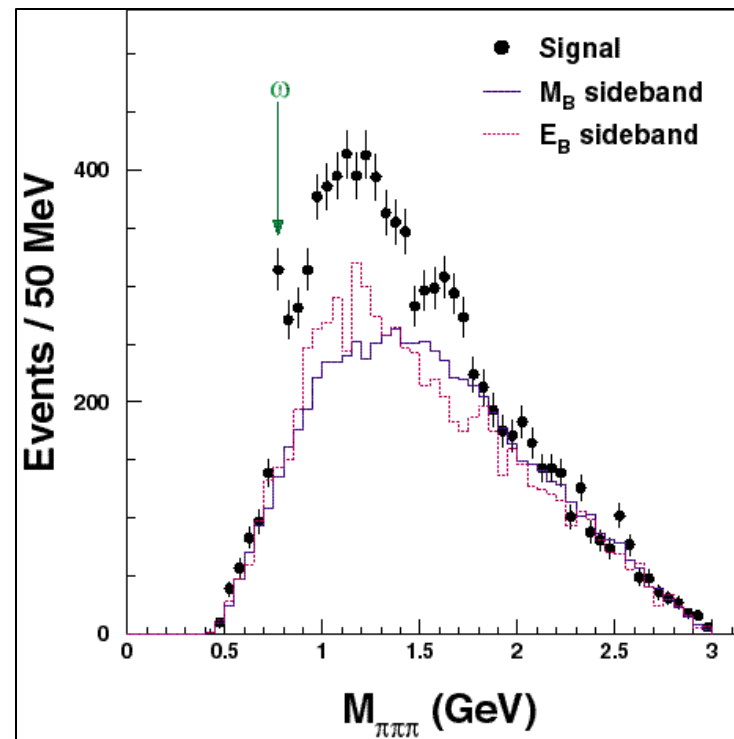
$(D^0 \rightarrow K^- \pi^+)$

358 ± 29



The $p^+p^-p^0$ Mass Distribution

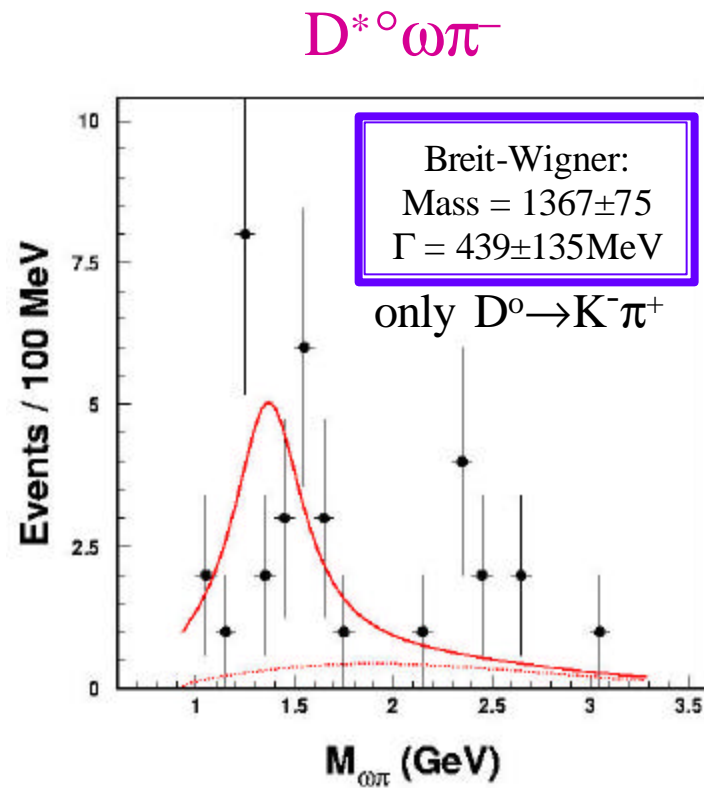
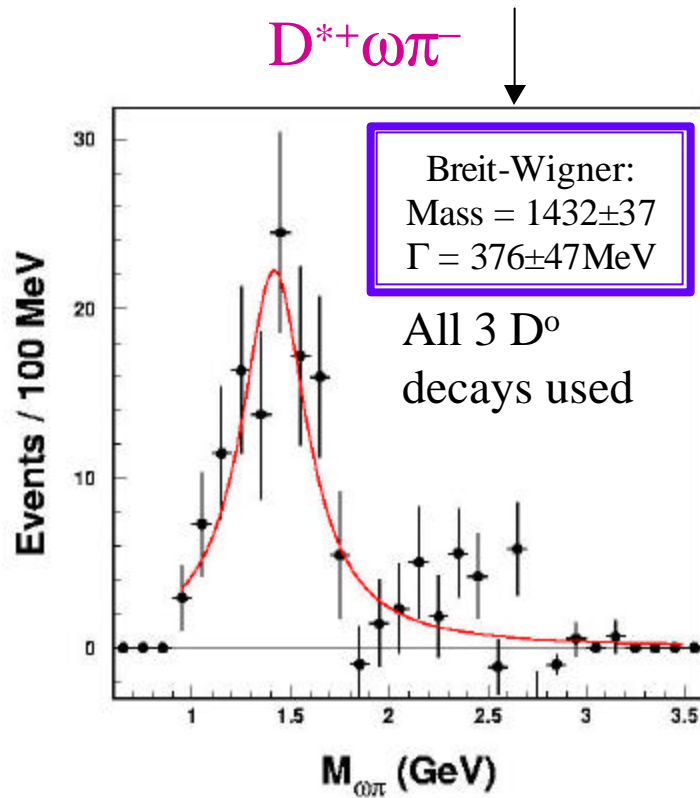
- What are the decay mechanisms for the $(4p)^-$ final state?
- We examine the $p^+p^-p^0$ mass spectrum (2 combinations/event). All 3 D^0 decay modes summed



Enlarged & Dalitz
plot exterior removed

The $\omega\pi^-$ Mass Distribution

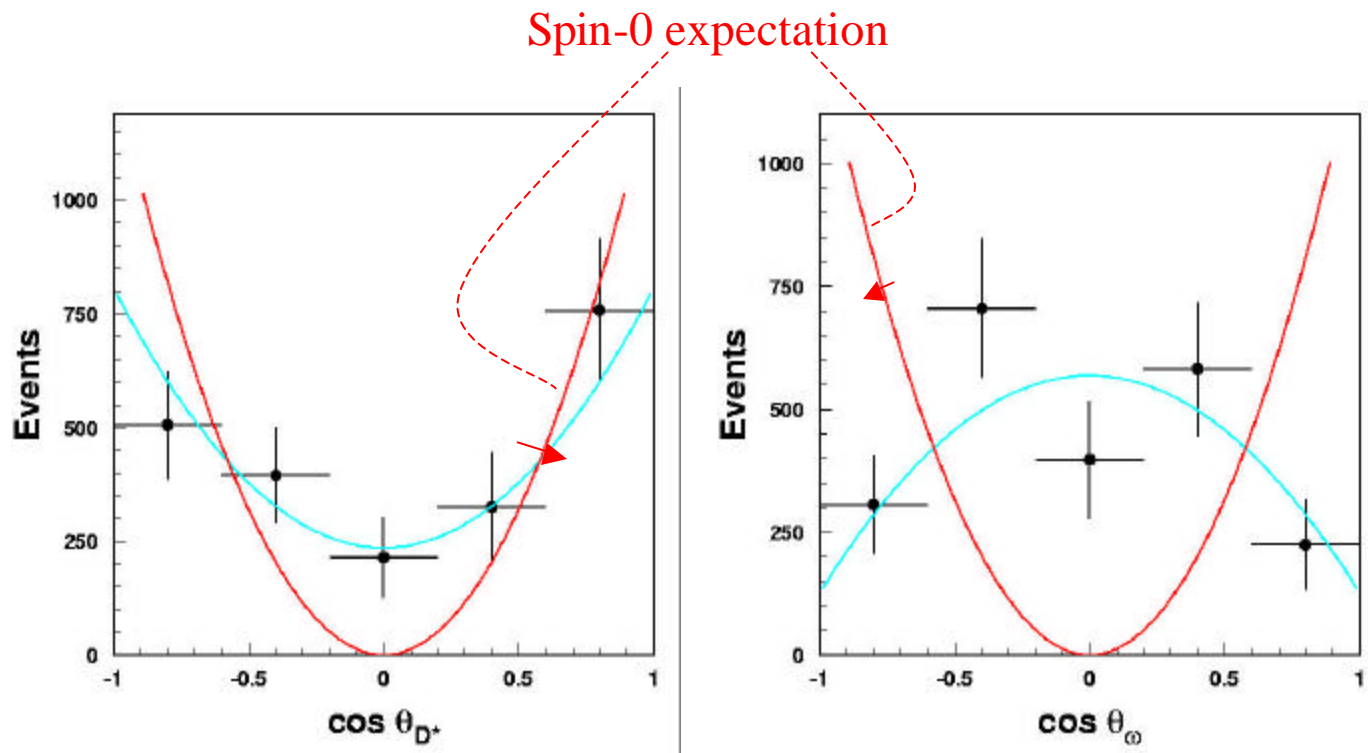
Fit M_B distribution in $\omega\pi$ mass bins



Possible resonance (A) at $M=1419 \pm 33$ MeV, $G=382 \pm 44$ MeV

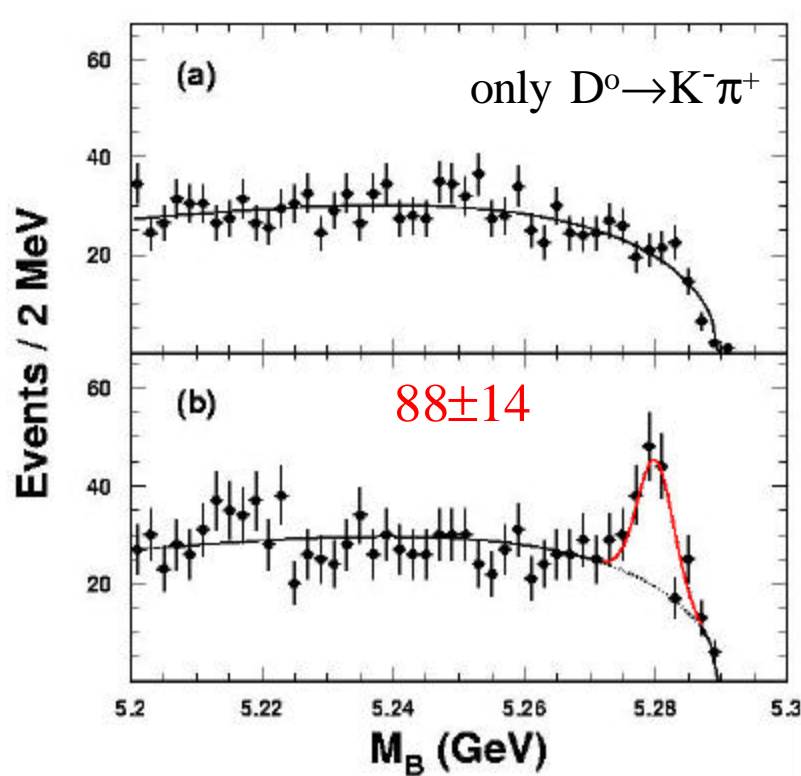
$D^{*+}(w)^-$ Angular Distributions

- For a spin-0 A the D^* & w would be fully polarized
- **Spin 0** $\chi^2/dof = 3.5$ ($\cos\theta_{D^*}$), 22 ($\cos\theta_w$) \Rightarrow Ruled out
- **Best fit** $\Rightarrow G_L/G = 0.63 \pm 0.09$ (D^{*+}), 0.10 ± 0.09 (w)

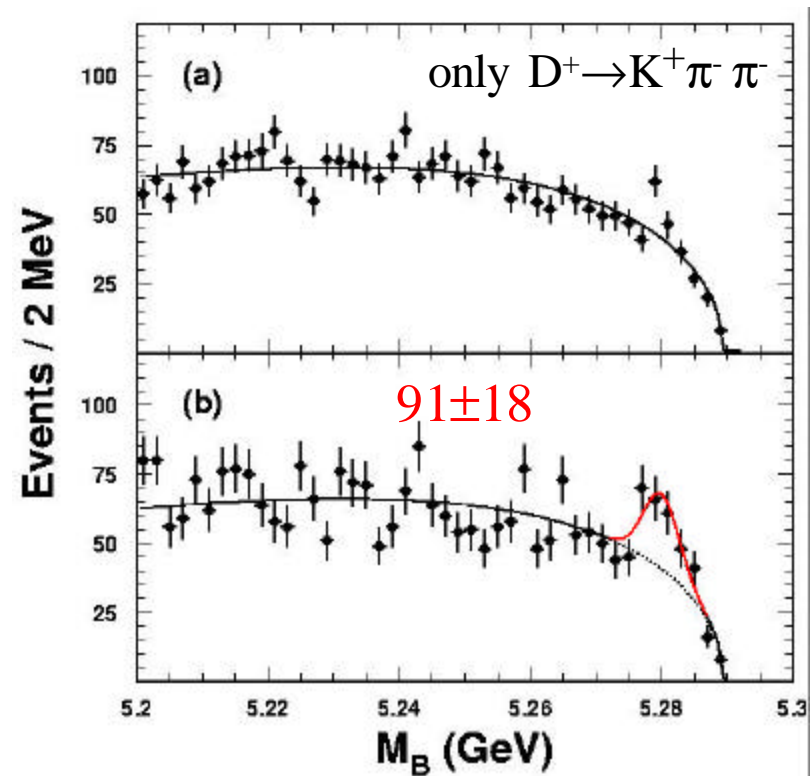


The $D\omega\pi^-$ Final State

$D^0\omega\pi^-$

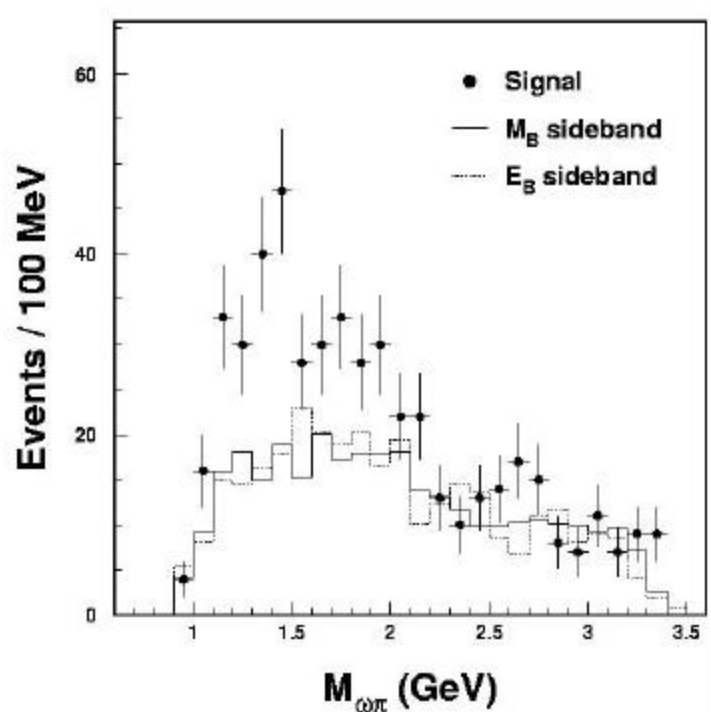


$D^+\omega\pi^-$

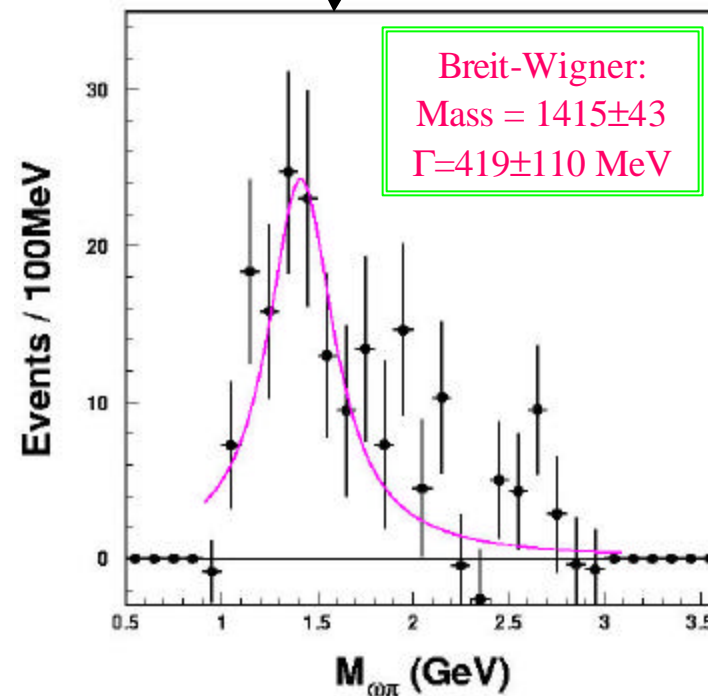


- ◆ Signal: $|\Delta E| < 2\sigma$ (18 MeV) Sideband: $3\sigma < |\Delta E| < 7\sigma$
- ◆ No signal in ω sidebands

The $\omega\pi^-$ Mass Distribution



Fit M_B distribution in $\omega\pi$ mass bins



- ◆ Combined $D^0\omega\pi^-$ and $D^+\omega\pi^-$ modes (179 events)
- ◆ Consistent with $D^*\omega\pi$ result
- ◆ Select (1.1–1.7 GeV) for angular study (104 events)

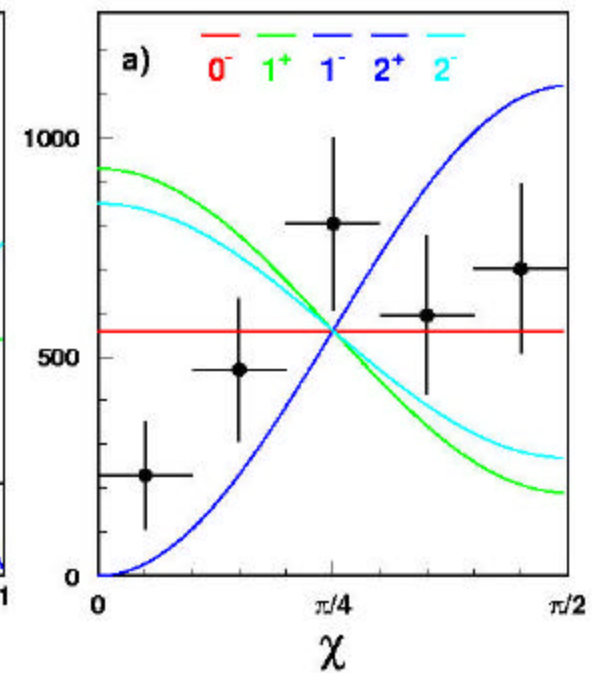
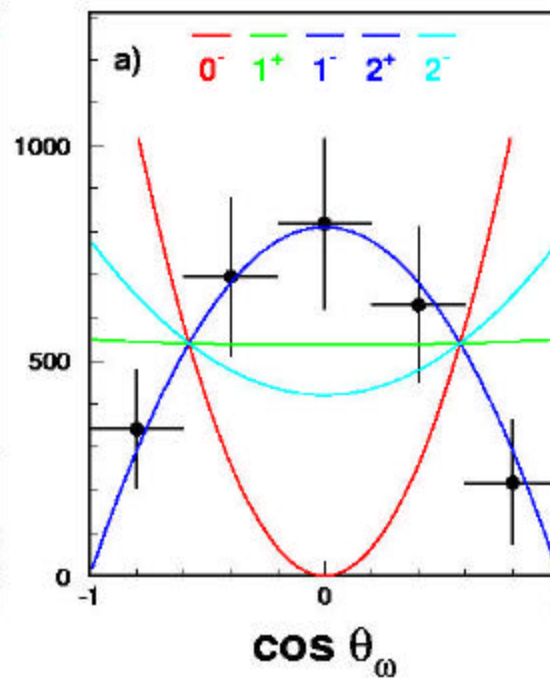
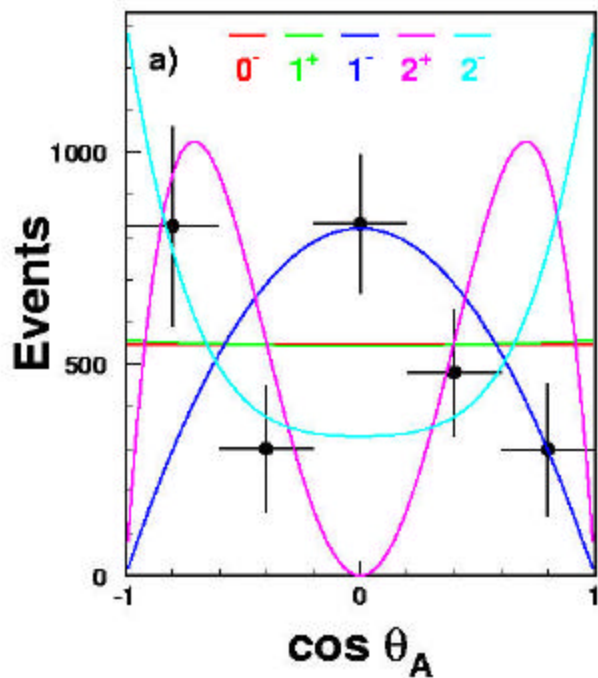
The Angular Distributions in

$B \rightarrow D A^- : A^- \rightarrow \omega p^-, \omega \rightarrow p^+ p^-$

\angle between ω in A frame
& A boost direction

\angle between normal of ω
decay plane & ω boost

\angle between A &
 ω decay planes



- ◆ Small efficiency corrections applied
- ◆ For 1^+ and 2^- , the longitudinal ratio (Γ_L/Γ) floats
- ◆ 1^- preferred, $\chi^2/\text{dof} (1^-) = 1.7, (2^+) = 3.2$
- ◆ A^- properties: mass = $1418 \pm 26 \pm 19$ MeV, $\Gamma = 388 \pm 41 \pm 32$ MeV

Identifying the A^- with the $r\bar{c}$

- Clegg & Donnachie: ($t\bar{c}(4p)n$, $e^+e^-\bar{c}p^+p^-$, $p^+p^+p^-p^-$) find two 1^- states with $(M, G) = (1463 \pm 25, 311 \pm 62)$ MeV & $(1730 \pm 30, 400 \pm 100)$ MeV, mixed with non- qq states, only the lighter one decays to $w\bar{p}$
- Godfrey & Isgur: Predict first radial excited $r\bar{c}$ at 1450 MeV, $G=320$ MeV, $B(r\bar{c} \rightarrow w\bar{p}^-) = 39\%$

Summary & Discussion of Rates

Mode	Br (%)	# of events
$\underline{B}^{\circ} \rightarrow D^{*+} \pi^{\circ} \pi^{+} \pi^{-} \pi^{-}$	$1.72 \pm 0.14 \pm 0.24$	1230 ± 70
$\underline{B}^{\circ} \rightarrow D^{*+} \omega \pi^{-}$	$0.29 \pm 0.03 \pm 0.04$	136 ± 15
$\underline{B}^{\circ} \rightarrow D^{+} \omega \pi^{-}$	$0.28 \pm 0.05 \pm 0.03$	91 ± 18
$B^{-} \rightarrow D^{*\circ} \pi^{\circ} \pi^{+} \pi^{-} \pi^{-}$	$1.80 \pm 0.24 \pm 0.25$	195 ± 26
$B^{-} \rightarrow D^{*\circ} \omega \pi^{-}$	$0.45 \pm 0.10 \pm 0.07$	26 ± 6
$B^{-} \rightarrow D^{\circ} \omega \pi^{-}$	$0.41 \pm 0.07 \pm 0.04$	88 ± 14

$r\bar{c}$ dominates the wp^{-} final state

$$G(\underline{B}^{\circ} \otimes D^{*+} r\bar{c}) / G(\underline{B}^{\circ} \otimes D^{+} r\bar{c}) = 1.04 \pm 0.21 \pm 0.06$$

$$G(B^{-} \otimes D^{*\circ} r\bar{c}) / G(B^{-} \otimes D^{\circ} r\bar{c}) = 1.10 \pm 0.31 \pm 0.06$$

$$G(B \otimes D^{*} r\bar{c}) / G(B \otimes D r\bar{c}) = 1.06 \pm 0.17 \pm 0.04$$

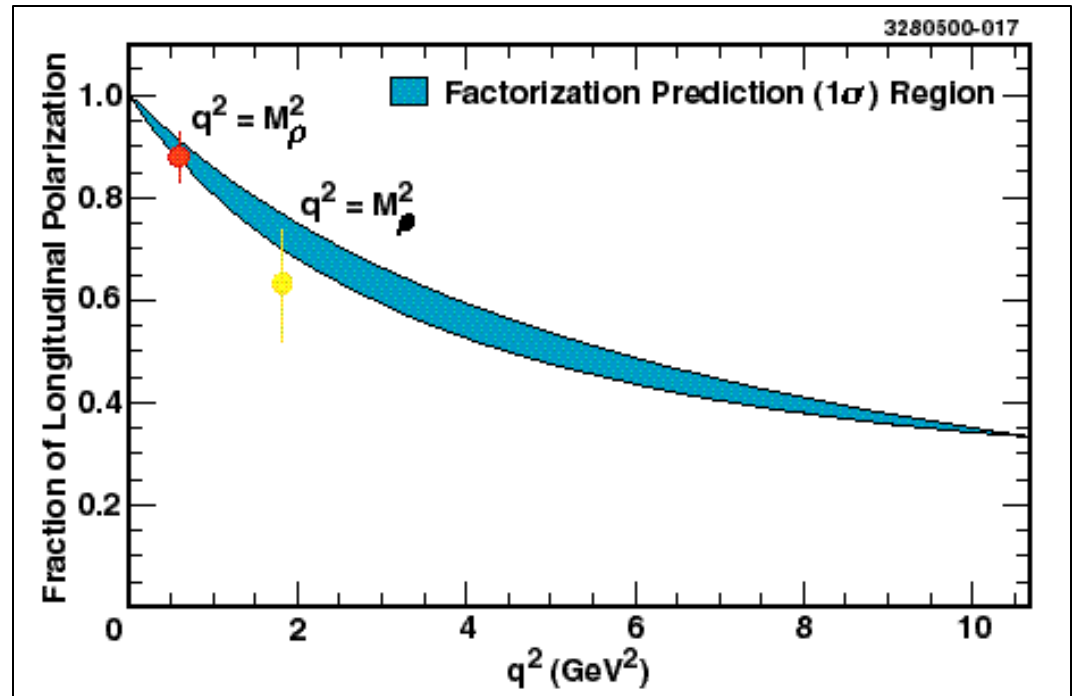
- Consistent with Heavy Quark Symmetry prediction (ratio = 1)
- With $B(r\bar{c} \otimes wp^{-}) = 39\%$, $G(B \otimes D^{(*)} r\bar{c}) \sim G(B \otimes D^{(*)} r^{-})$

Testing Factorization

Polarization:

$$\frac{\Gamma_L/\Gamma (B \rightarrow D^{*+} h^-)}{\Gamma_L/\Gamma (B \rightarrow D^{*+} l^- \mathbf{n})|_{q^2=m_h^2}}$$

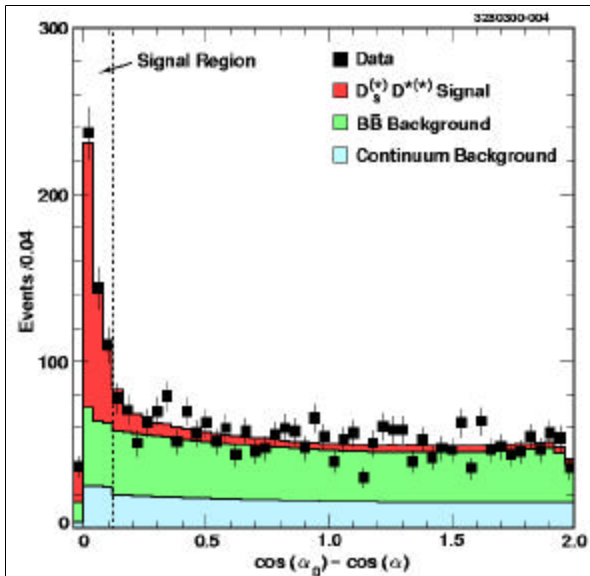
Branching Fractions:



$$\Gamma(B \rightarrow D^{*+} h^-) / d\Gamma/dq^2 (B \rightarrow D^{*+} l^- \mathbf{n})|_{q^2=m_h^2} = 6\pi^2 c_1^2 f_h^2 |V_{ud}|^2$$

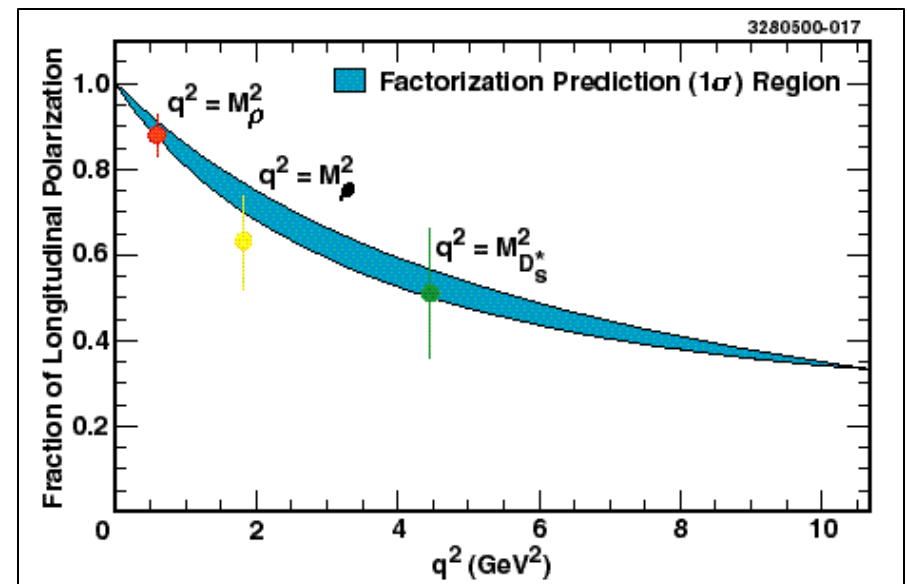
Using $B(\rho'^- \rightarrow \omega \pi^-) = 39\% \Rightarrow f_{\rho'} = 167 \pm 23 \text{ MeV}$

Extending q^2 : $B \rightarrow D^* D_s^*$



Final State	$B(\%)$
$D^* + D_s^-$	$1.10 \pm 0.18 \pm 0.10 \pm 0.28$
$D^* + D_s^0$	$1.82 \pm 0.37 \pm 0.24 \pm 0.46$
$D^{*0} + D_s^{*0}$	$2.73 \pm 0.78 \pm 0.48 \pm 0.68$

$D^{*+} +$	$G_L/G (\%)$
r^-	87.8 ± 5.3
r^0	63 ± 9
D_s^-	$50.6 \pm 13.9 \pm 3.6$



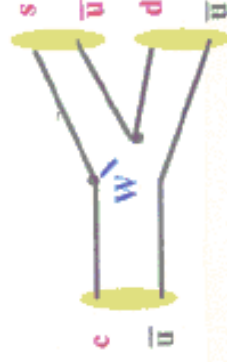
What About Charm?

$m_c \ll m_b \Rightarrow$ less energy release
stronger FSI

Factorization is in trouble!

	$R_{\text{Experiment}}$	R_{Theory}
$D^0 \rightarrow K^-\pi^+$	1.19 ± 0.14	1.18 ± 0.03
$D^0 \rightarrow K^{*-}\pi^+$	3.09 ± 0.82	1.18 ± 0.03

Color Suppression



$B(D^0 \rightarrow K^0 \rho^0) / B(D^0 \rightarrow K^- \rho^+)$
 $B(D^0 \rightarrow K^0 \pi^0) / B(D^0 \rightarrow K^- \pi^+)$
 $B(D^0 \rightarrow \bar{K}^{*0} \pi^0) / B(D^0 \rightarrow K^{*-} \pi^+)$
 $B(D^0 \rightarrow \pi^0 \pi^0) / B(D^0 \rightarrow \pi^- \pi^+)$
 $B(D_s^+ \rightarrow \bar{K}^{*0} K^+) / B(D_s \rightarrow \phi \pi^+)$
 $B(D_s^+ \rightarrow \bar{K}^0 K^+) / B(D_s \rightarrow \phi \pi^+)$

0.08 ± 0.04
 0.57 ± 0.13
 0.47 ± 0.23
 0.77 ± 0.25
 0.95 ± 0.10
 1.01 ± 0.16

0.08

0.77

not effective in Charm decay

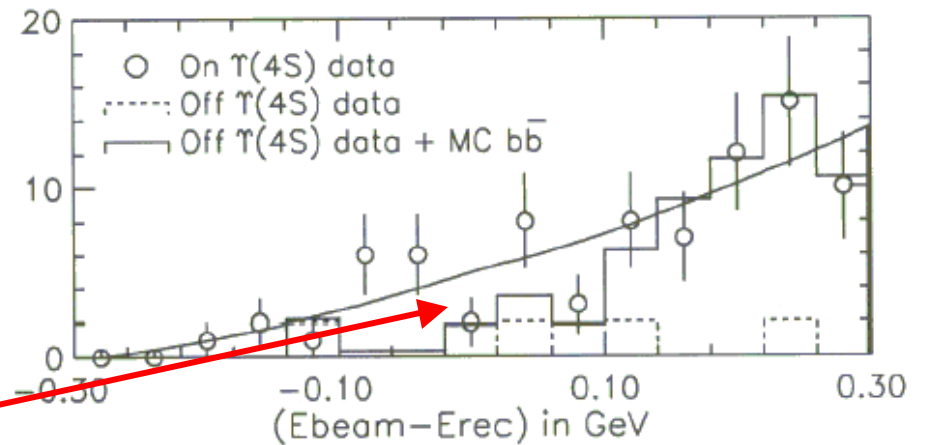
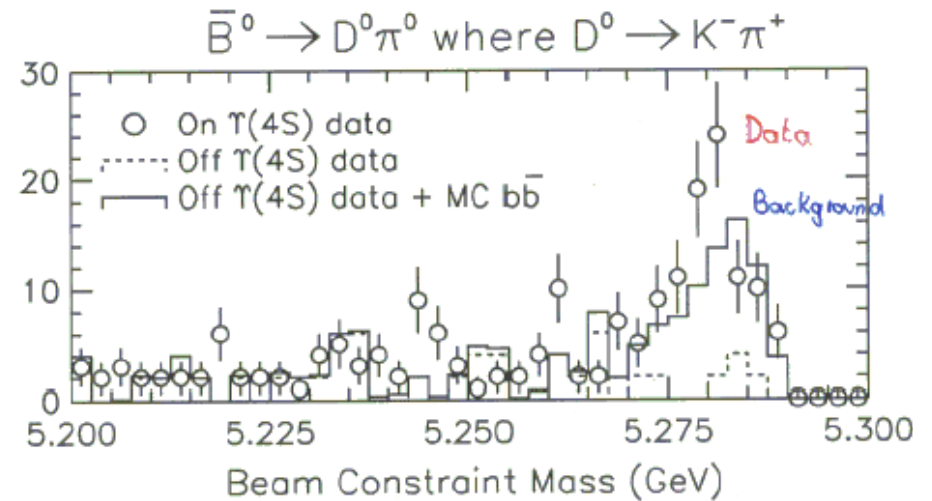
$\tau_{D^+} \sim 2.5 \tau_{D^0}$

Destructive Interference

Color Suppression in B Decays ...

CLEO:

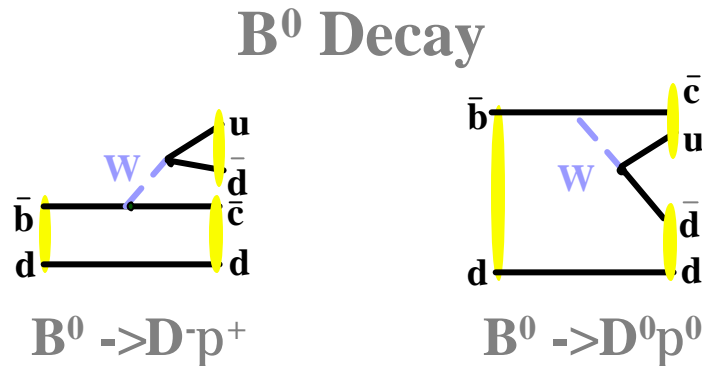
Decay Mode	U. L. (90% CL)	BSW (%)
$\bar{B}^0 \rightarrow D^0 \pi^0$	< 0.012	0.012
$\bar{B}^0 \rightarrow D^0 \rho^0$	< 0.040	0.008
$\bar{B}^0 \rightarrow D^0 \eta$	< 0.068	0.006
$\bar{B}^0 \rightarrow D^0 \eta'$	< 0.086	0.002
$\bar{B}^0 \rightarrow D^0 \omega$	< 0.056	0.008
$\bar{B}^0 \rightarrow D^{*0} \pi^0$	< 0.044	0.012
$\bar{B}^0 \rightarrow D^{*0} \rho^0$	< 0.060	0.013
$\bar{B}^0 \rightarrow D^{*0} \eta$	< 0.069	0.007
$\bar{B}^0 \rightarrow D^{*0} \eta'$	< 0.27	0.002
$\bar{B}^0 \rightarrow D^{*0} \omega$	< 0.080	0.013



No Signal!

...works

Charged B Decay and Interference

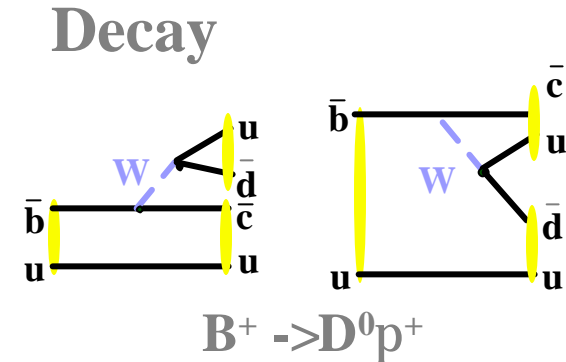


Different final states

BSW:

a_1

a_2



Interference

$a_1 + Z a_2$

- a_1, a_2 are phenomenological constants
- process dependent
- **great success in charm decay:**

destructive interference

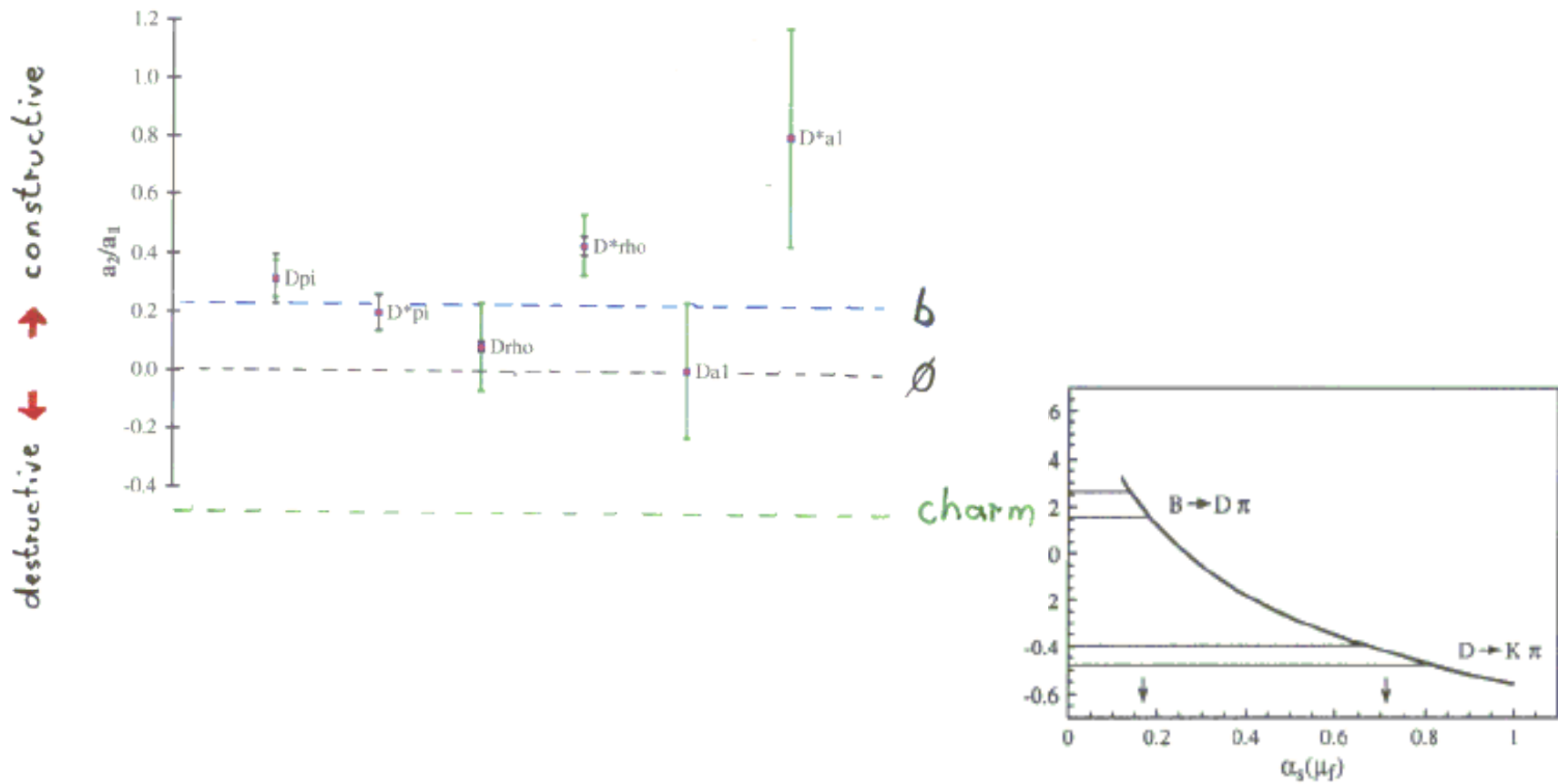
smaller G_{Hadronic} for D^+

$\tau_{D^+} \gg \tau_{D^0}$

Mode	B^0 ($\times 10^{-3}$)	B^+ ($\times 10^{-3}$)
$D\pi$	3.0 \pm 0.4	5.3 \pm 0.5
$D\rho$	7.9 \pm 1.4	13.4 \pm 1.8
Da_1	6.0 \pm 3.3	
$D\rho'$	2.8 \pm 0.6	4.1 \pm 0.8
$D^*\pi$	2.8 \pm 0.2	4.6 \pm 0.4
$D^*\rho$	6.8 \pm 3.4	15.5 \pm 3.1
D^*a_1	13 \pm 2.7	19 \pm 5
$D^*\rho'$	2.9 \pm 0.5	4.1 \pm 0.8

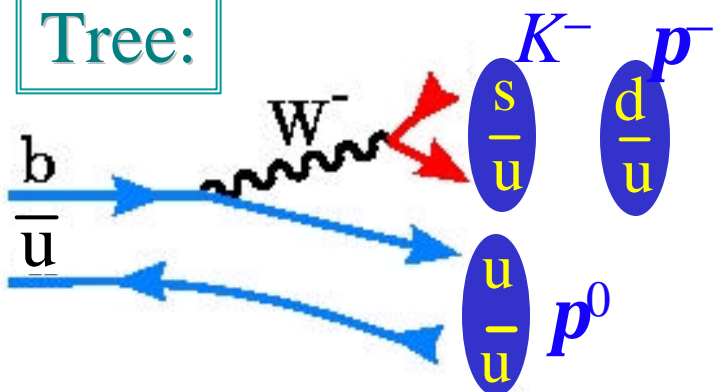
More on charged B Decays

$B \rightarrow D^{(*)}(n\pi)^- : \text{ FITS FOR } a_2/a_1$

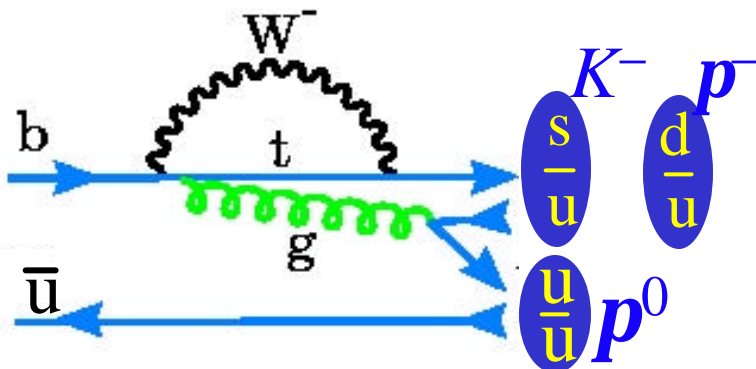


Rare B Decay

Tree:



Penguin:

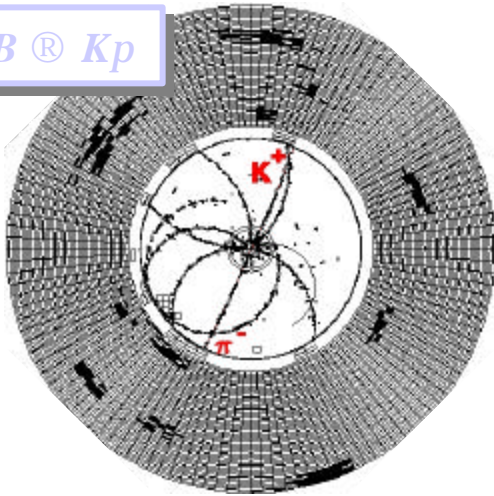


- Tree decays $b^{\text{R}} u$ vs. $b^{\text{R}} c$ suppressed by $|V_{ub}|^2/|V_{cb}|^2 \sim 0.01$
- Additional $|V_{us}|^2/|V_{ud}|^2 \sim 0.04$ for K^-
- Expect tree dominantly $b^{\text{R}} u\bar{u}d$.

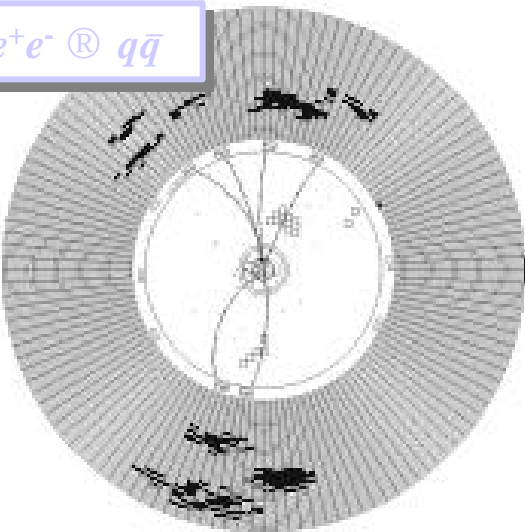
- Decays $b^{\text{R}} s, d$ GIM-suppressed
- Loop diagram $\mu(m_t/m_W)^2$.
- $|V_{td}|^2/|V_{ts}|^2 \sim 0.01$
- Expect penguins dominantly $b^{\text{R}} u\bar{u}s$.

$B \rightarrow K^+ p^- / p^+ p^-$ Topology

$B \rightarrow Kp$



$e^+e^- \rightarrow q\bar{q}$

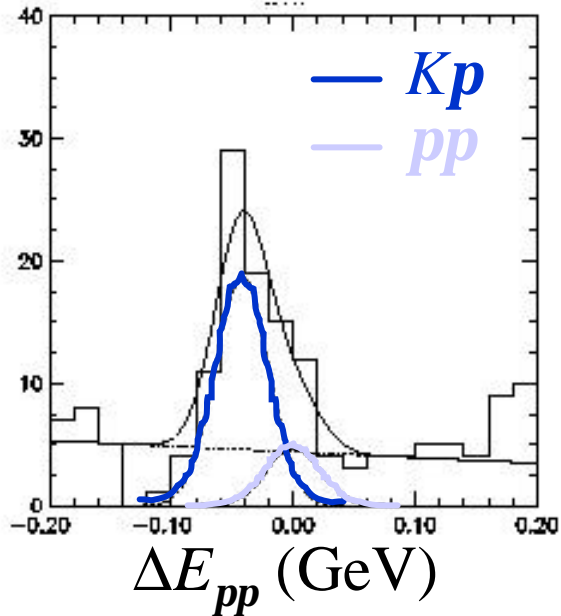


- $P_{\text{daughter}} \sim 2.55 - 2.85 \text{ GeV}/c$
(higher than for $b \rightarrow c$ decays)
- Major background from $e^+e^- \rightarrow q\bar{q}$ “continuum”
- Continuum events are “jetty” in topology
- $P_B \sim 300 \text{ MeV}/c \Rightarrow B\bar{B}$ events “spherical”
- Continuum suppression from ML fit to several kinematic and topological variables (more efficient).
- Continuum suppression factor of $\sim 10^6$, efficiency for Kp/pp of $\sim 40\%$

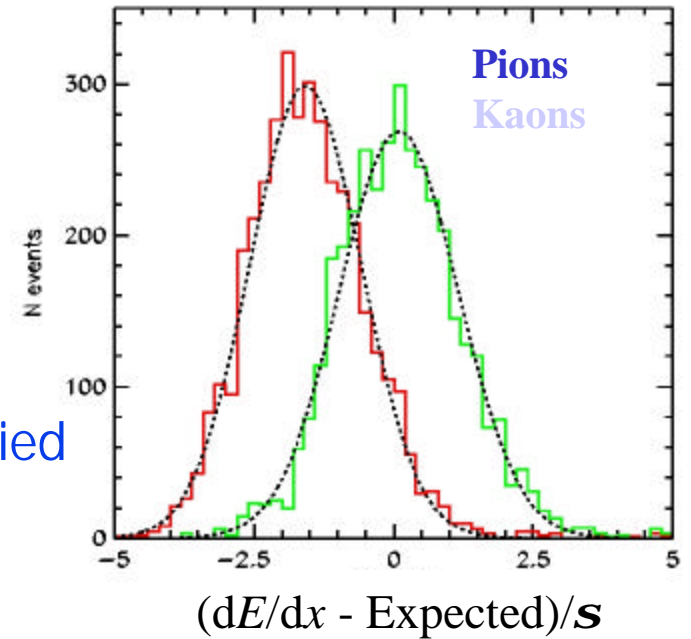
K/p Separation

- *Kp* vs. *pp* from dE/dx in drift chamber
 - Resol. confirmed with $D^{*-} \rightarrow D^0 \pi^-$, $D^0 \rightarrow K^- \pi^+$
- Also separation from kinematics:

$$\Delta E_{pp} = E_p + E_p - E_{beam}$$

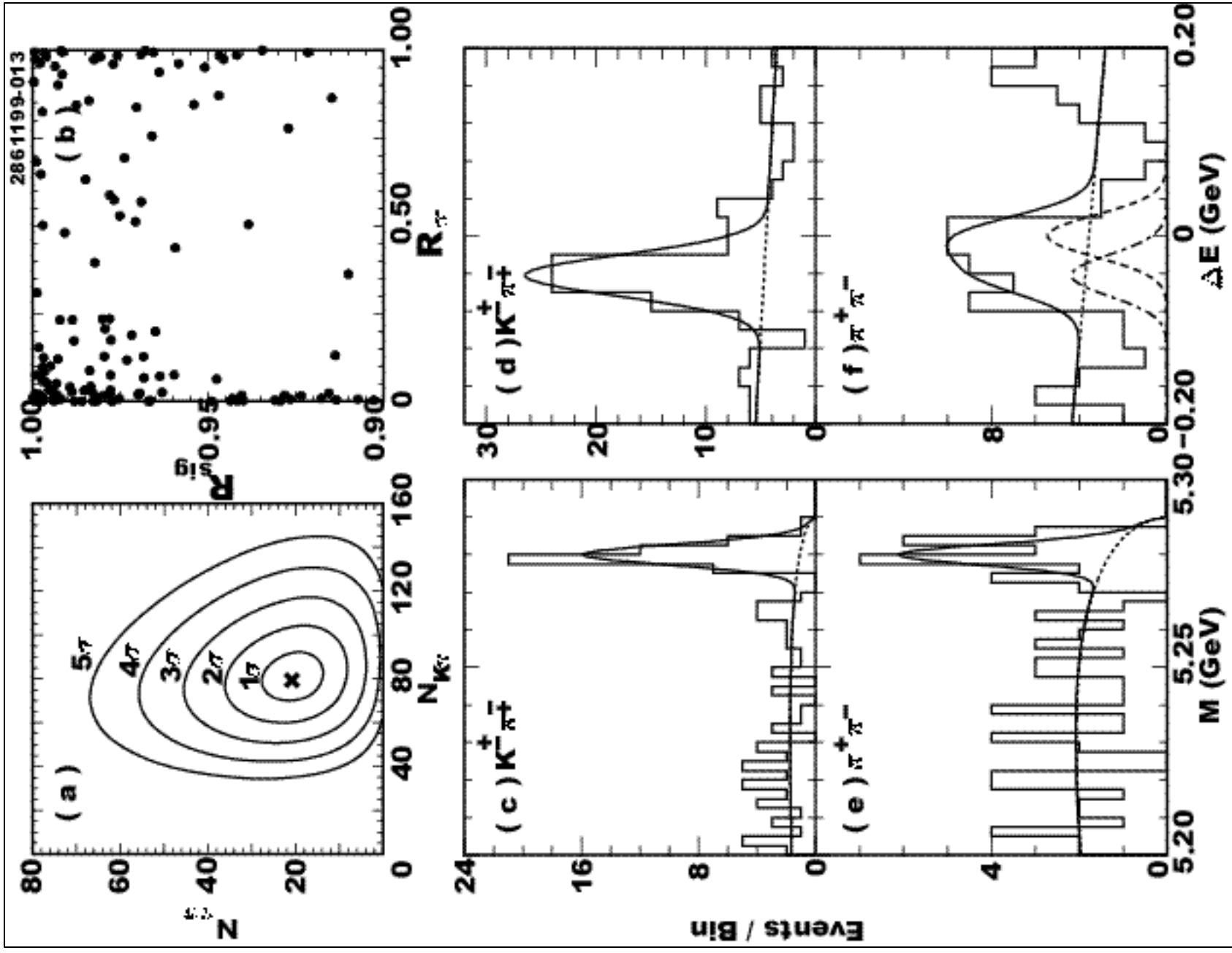


• ΔE resolution studied with $D^0 \rightarrow K^- \pi^+ (\pi^0)$ mass resolutions



	S_{DE}	DE	dE/dx
CLEO II	25 MeV	1.7s	1.7s
CLEO II.V	20 MeV	2.1s	2.0s

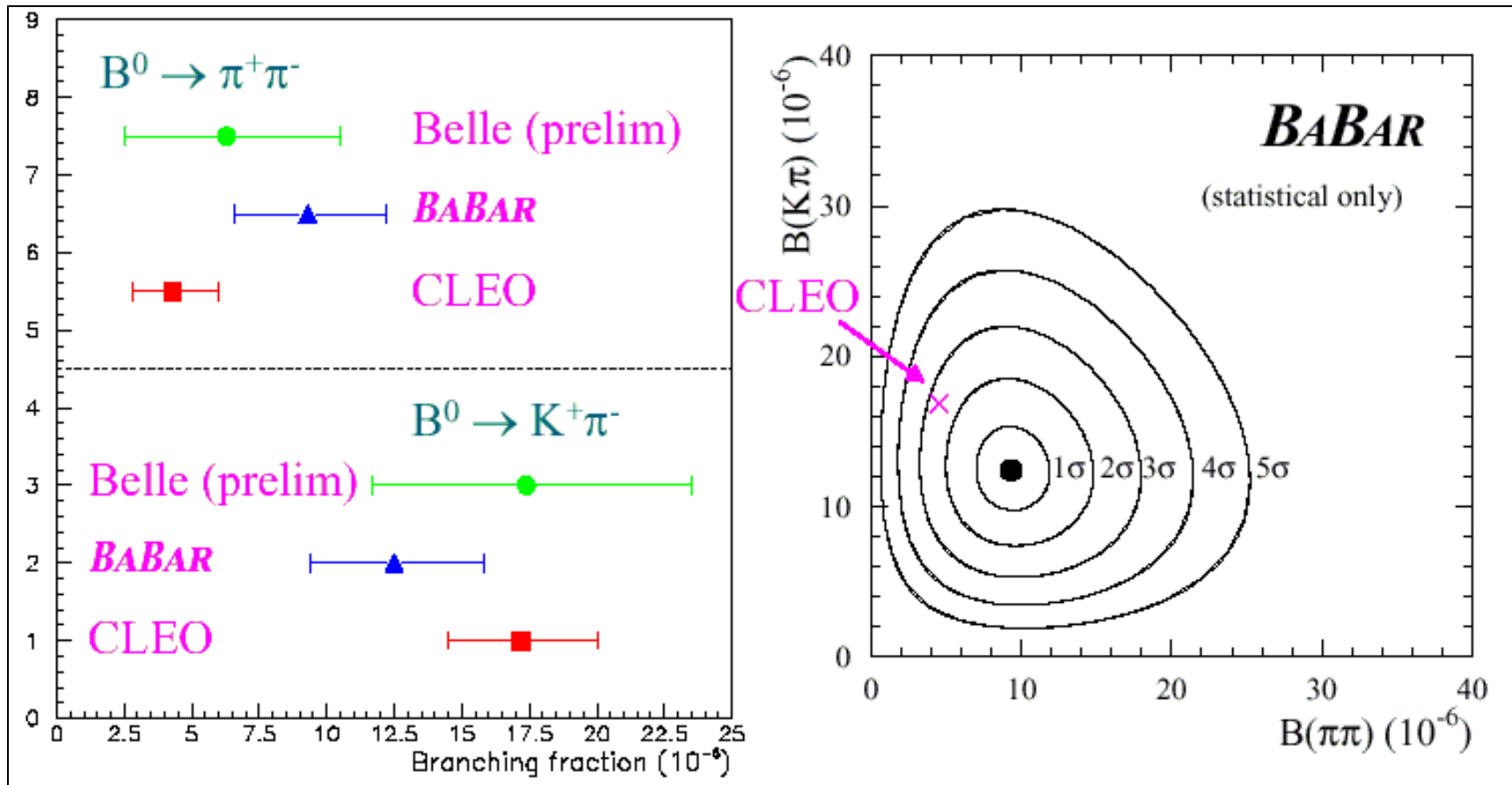
Fit Results for $B \rightarrow Kp, B \rightarrow pp$



$B^{\oplus} K p, B^{\oplus} p p$ Summary

	Signal (events)	# S	e (%)	BR ($\times 10^{-6}$)
$p^+ p^-$	$20.0^{+7.6}_{-6.5}$	4.2	48	$4.3^{+1.6}_{-1.4} \pm 0.5$
$p^+ p^0$	$21.3^{+9.7}_{-6.5}$	3.2	39	< 12.7
$p^0 p^0$	$6.2^{+4.8}_{-3.7}$	2.0	29	< 5.7
$K^+ p^-$	$80.2^{+11.8}_{-11.0}$	11.7	48	$17.2^{+2.5}_{-2.4} \pm 1.2$
$K^+ p^0$	$42.1^{+10.9}_{-9.9}$	6.1	38	$11.6^{+3.0}_{-2.7} \pm 1.4$
$K^0 p^+$	$25.4^{+6.4}_{-5.6}$	7.6	14	$18.2^{+4.6}_{-4.0} \pm 1.6$
$K^0 p^0$	$16.1^{+5.9}_{-5.1}$	4.9	11	$14.6^{+5.9}_{-5.1} \pm 2.4$
$K^+ K^-$	$0.7^{+3.4}_{-0.0}$	0.0	48	< 1.9
$K^+ K^0$	$1.4^{+2.4}_{-1.3}$	1.1	14	< 5.1

Comparison of $B \rightarrow pp$ Results

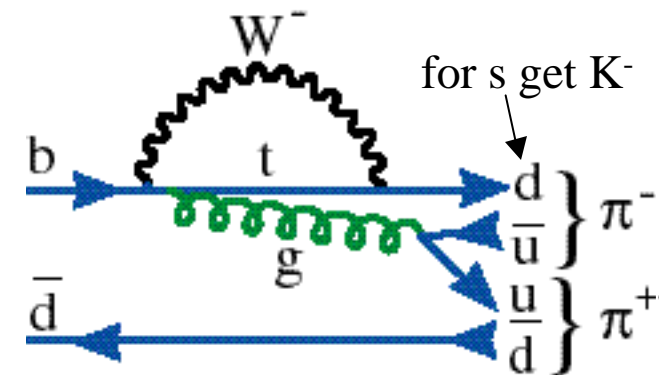
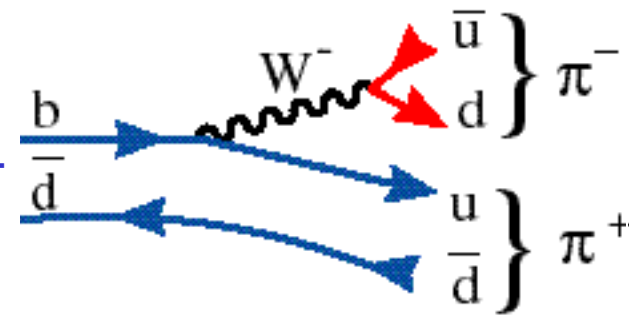


Problems with measuring a using $B^0 \rightarrow p^+ p^-$

- Using $B^0 \rightarrow p^+ p^-$ would be nice, but large Penguin term, CLEO:

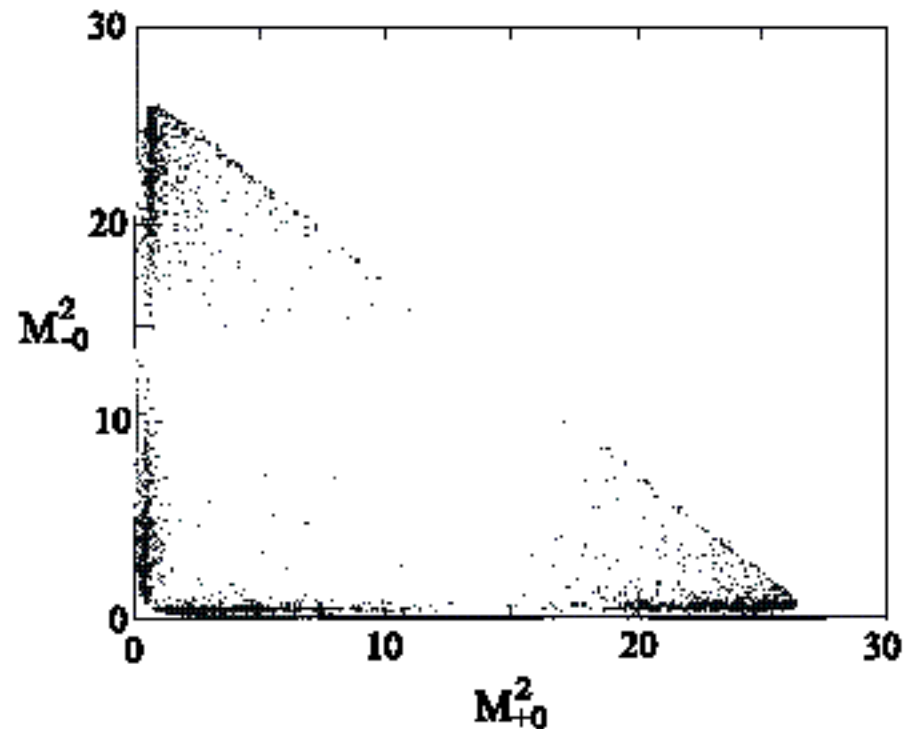
$$B(B^0 \rightarrow p^+ p^-) \sim 0.8 \times 10^{-5}, \text{ while } B(B^0 \rightarrow K^+ p^-) = (1.4 \pm 0.3 \pm 0.2) \times 10^{-5}$$

- The effect of the Penguin must be measured in order to determine a . Can be done using Isospin, but requires a rate measurements of $p^- p^0$ and $p^0 p^0$ (Gronau & London). However, this is daunting.



Measuring a using $B^0 \rightarrow \rho^0 p^+ p^- p^0$

- A Dalitz Plot analysis gives both $\sin(2a)$ and $\cos(2a)$ (Snyder & Quinn)
- CLEO has measured the branching ratios $B(B^0 \rightarrow \rho^0 p^+ p^-) = (1.5 \pm 0.5 \pm 0.2) \times 10^{-5}$
- $0.5B(B^0 \rightarrow \rho^- p^+ + \rho^+ p^-) = (3.5 \pm 1.0 \pm 0.5) \times 10^{-5}$



Snyder & Quinn showed that 1000-2000 tagged events are sufficient

Comparison of rp modes

Final State:	r^-p^+	r^+p^-	r^-p^0	r^0p^-	r^0p^0
CLEO(10^{-5}):	$<3.5 \pm 1.2>$		<7.7	$1.5 \pm 0.5 \pm 0.2$	<1.8
Ciuchini et al.:	1.0-7.5	0.2-1.9	0.3-2.6	0.5-1.1	0.00-0.02
Ali et al.:	2.1-3.4	0.6-0.9	1.1-1.6	0.1-0.7	0.00-0.02
$B(B^+ \rightarrow r^0 K^+) = < 2.2 \times 10^{-5} @ 90\% \text{ c.l. (CLEO)}$					

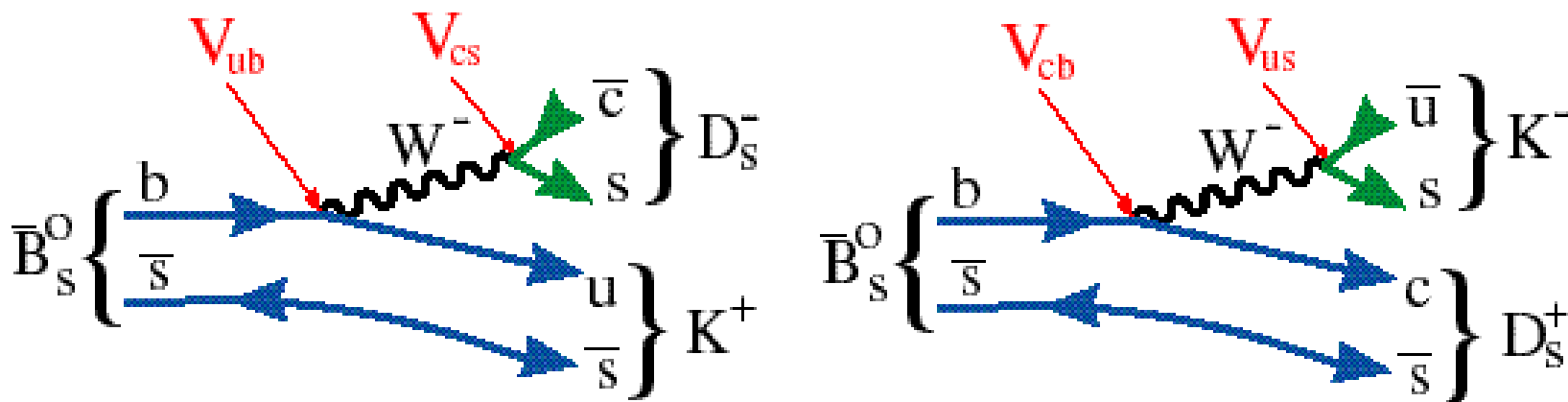
{To measure the three neutral rp modes, requires triggering on hadronic decays with high efficiency, RICH particle ID and high quality photon detection}

Ways of measuring g

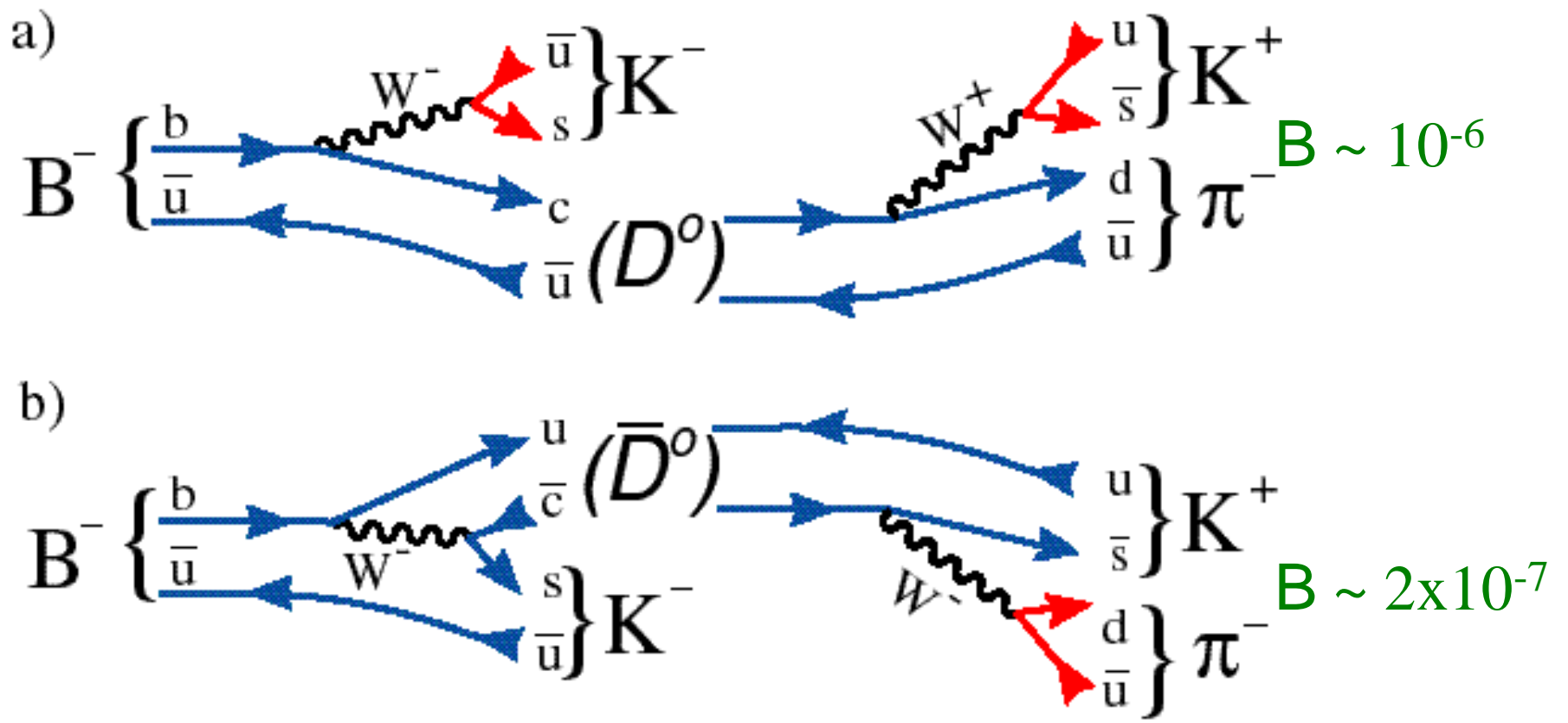
- **May be easier to measure than a**
- **There are 4 ways of determining g**
 - Time dependent flavor tagged analysis of $B_s \rightarrow D_s K^-$
 - Measure rate difference between $B^- \rightarrow D^0 K^-$ and $B^+ \rightarrow D^0 K^+$
 - Rate measurements in $K^0 \pi^\pm$ and $K^\pm \pi^\mu$ (Fleisher-Mannel) or rates in $K^0 \pi^\pm$ & asymmetry in $K^\pm \pi^0$ (Neubert-Rosner). Has theoretical uncertainties but can be useful.
 - Use U spin symmetry $d \leftrightarrow s$: measure time dependent asymmetries in both $B^0 \rightarrow \pi^+ \pi^-$ & $B_s \rightarrow K^+ K^-$ (Fleischer).
 - Ambiguities here as well but they are different in each method, and using several methods can resolve them.

$B_s \text{ (R) } D_s K^m$ Decay processes \pm

Diagrams for the two decay modes, $\mathcal{B} \sim 10^{-4}$ for each



$B^- \rightarrow [K^+ p^-] K^-$ Decay processes



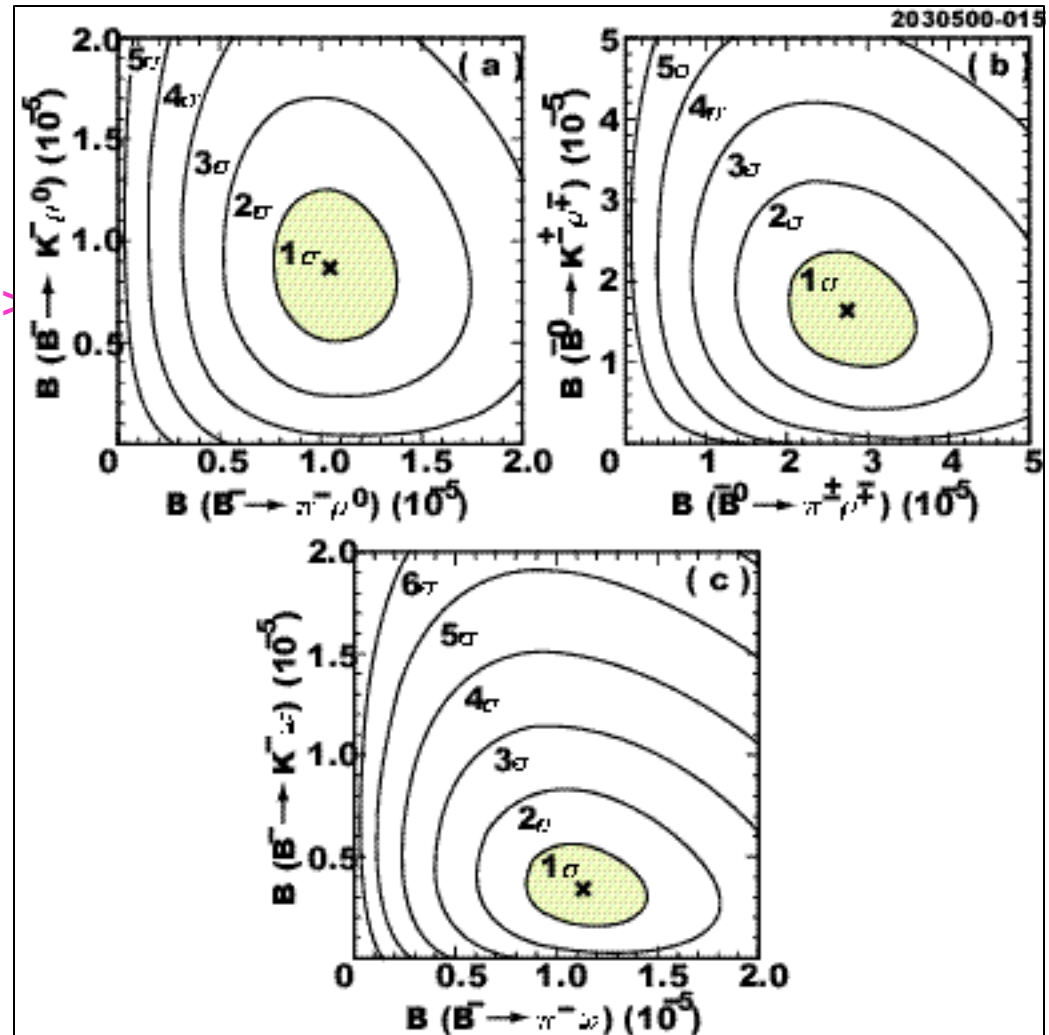
Exclusive $b \rightarrow u$ Transitions

- Many hadronic $b \rightarrow u$ transitions observed
- New study includes 14 channels (hep-ex/0006008)
- In general - good agreement with theory
- Full Dalitz analyses could determine a and g

Mode	Yield	Signif.	B (10^{-6})	Theory
$B^- \rightarrow \pi^- \rho^0$	$29.8^{+9.3}_{-9.6}$	5.4σ	$10.4^{+3.9}_{-3.6} \pm 2.1$	0.4 - 13.0
$B^- \rightarrow K^- \rho^0$	$22.4^{+10.7}_{-9.1}$	3.7σ	< 17	0.0-6.1
$B^- \rightarrow \pi^0 \rho^-$	$23.7^{+8.4}_{-7.4}$	5.1σ	< 43	3.0 - 27.0
$B^- \rightarrow \pi^0 K^{*-}$	$2.6^{+1.2}_{-2.6}$	1.0σ	< 31	0.5 - 24.0
$B^- \rightarrow \pi^- \omega$	$28.5^{+3.9}_{-7.3}$	6.2σ	$11.3^{+3.3}_{-2.9} \pm 1.4$	0.6 - 24
$B^- \rightarrow K^- \omega$	$7.9^{+6.0}_{-1.7}$	2.1σ	< 7.9	0.2 - 14.0
$B^- \rightarrow \pi^- K^{*0}$	$13.4^{+6.2}_{-5.2}$	3.6σ	< 16	3.4 - 13.0
$B^- \rightarrow K^- K^{*0}$	$0.0^{+2.3}_{-0.0}$	0.0σ	< 5.3	0.2 - 1.0
$B^0 \rightarrow \pi^+ \rho^-$	$31.0^{+0.1}_{-3}$	5.6σ	$27.5^{+8.4}_{-7.4} \pm 4.2$	12 - 93
$B^0 \rightarrow K^+ \rho^-$	$16.4^{+7.8}_{-6.6}$	3.5σ	< 32.5	0.0- 12.0
$B^0 \rightarrow \pi^0 \rho^0$	$5.4^{+6.5}_{-1.8}$	1.2σ	< 5.5	0.0 - 2.5
$B^0 \rightarrow \pi^0 \omega$	$1.5^{+3.5}_{-1.5}$	0.6σ	< 5.5	0.0 - 12.0
$B^0 \rightarrow K^0 \omega$	$7.0^{+3.8}_{-2.9}$	3.9σ	< 21	0.0 - 17.0
$B^0 \rightarrow \pi^0 K^{*0}$	$0.0^{+3.0}_{-0.0}$	0.0σ	< 3.6	0.7 - 6.1

Likelihood Contours

- $B(B^0 \rightarrow r^+ p) / BR(B^0 \rightarrow r^0 p)$ fitted simultaneously
- $BR(B^0 \rightarrow r^+ p) / BR(B^0 \rightarrow r^0 p)$ smaller than expected (x)



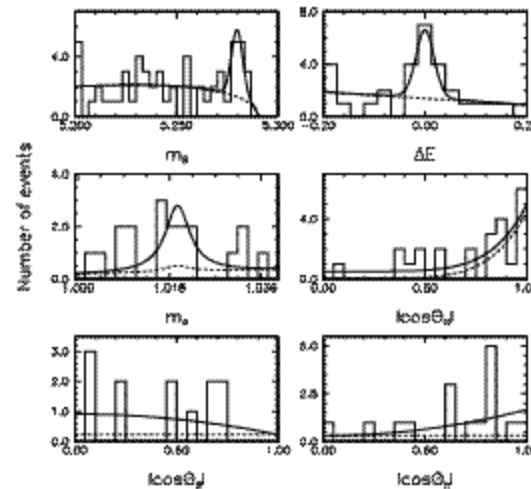
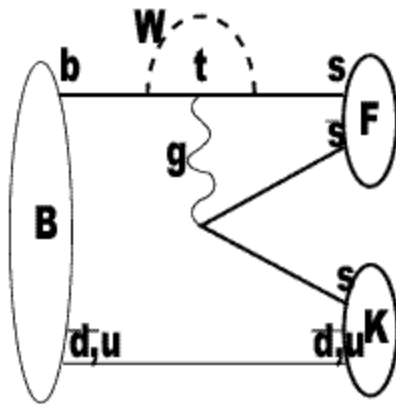
Modes with h and $h\zeta$

Mode	Signif.	B (10^{-6})
$B^+ \rightarrow \eta' K^+$	10.8σ	$89^{+10}_{-9} \pm 7$
$B^0 \rightarrow \eta' K^0$	11.7σ	$89^{+18}_{-16} \pm 9$
$B^+ \rightarrow \eta' \pi^+$	0.0σ	< 12
$B^0 \rightarrow \eta' \pi^0$	0.0σ	< 5.7
$B^+ \rightarrow \eta' K^{*+}$	1.8σ	< 35
$B^0 \rightarrow \eta' K^{*0}$	1.8σ	< 24
$B^+ \rightarrow \eta' \rho^+$	2.4σ	< 33
$B^0 \rightarrow \eta' \rho^0$	0.0σ	< 12
$B^+ \rightarrow \eta K^+$	0.8σ	< 6.9
$B^0 \rightarrow \eta K^0$	0.0σ	< 9.3
$B^+ \rightarrow \eta \pi^+$	0.6σ	< 5.7
$B^0 \rightarrow \eta \pi^0$	0.0σ	< 2.9
$B^+ \rightarrow \eta K^{*+}$	4.8σ	$26.4^{+9.6}_{-8.3} \pm 3.3$
$B^0 \rightarrow \eta K^{*0}$	5.1σ	$13.8^{+5.5}_{-4.6} \pm 1.6$
$B^+ \rightarrow \eta \rho^+$	1.3σ	< 15
$B^0 \rightarrow \eta \rho^0$	1.3σ	< 10

m_B
 m_B

PRL 80, 3710 (1998) + PRL 85, 520 (2000)

Pure Penguins: $B \rightarrow fK$



$$BR(B^- \rightarrow \phi K^-) = (6.4^{+2.6+0.5}_{-2.1-2.0}) \times 10^{-6}$$

$$BR(B^0 \rightarrow \phi K^0) = (5.9^{+1.0+1.1}_{-2.9-0.9}) \times 10^{-6} \rightarrow < 1.2 \times 10^{-5}$$

$$\text{Combined result: } BR(B \rightarrow \phi K) = (6.2^{+2.0+0.7}_{-1.8-1.7}) \times 10^{-6}$$

- Pure gluonic penguin, simple final state, sensitive to $\sin 2\beta$
- Theoretical uncertainties are large:
 - Deshpande+He: inclusive $B \rightarrow \phi X_s \sim (0.6 - 2.0) \times 10^{-4}$
 - ϕK fraction of $\phi X_s \sim 10\%$

arg(V_{ub}) (= g) from Decay Rates

- Fleischer-Mannel (*Phys. Rev. D* 57, 2752(1998))

$$R \equiv \frac{\Gamma(B^0 \rightarrow K^- \pi^+)}{\Gamma(B^+ \rightarrow K^0 \pi^+)} \geq \sin^2 g \quad \text{CLEO: } R = 1.01 \pm 0.26$$

- Neubert-Rosner (*Phys. Lett. B* 441, 403 (1998))

$$R_* \equiv \frac{\Gamma(B^+ \rightarrow K^0 \pi^+)}{2\Gamma(B^+ \rightarrow K^+ \pi^0)} \quad (1-R_*)/e_{3/2} \leq |d_{EW} - \cos g|$$

$$0.58 \pm 0.74 \leq | (0.64 \pm 0.15) - \cos g |$$

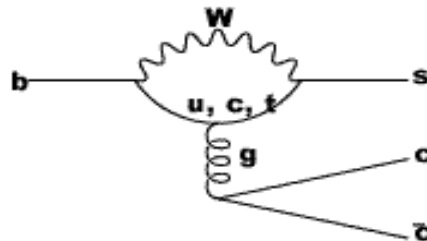
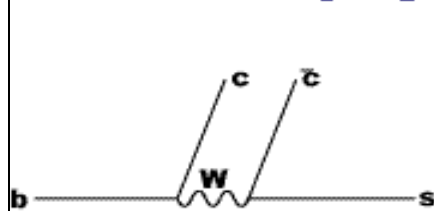
- Also model-dependent fit to many CLEO branching ratios of pp , Kp , rp , wp (Wuerthwein *et al.* hepex/9910014):

$$84 < g < 154 \quad (90\% \text{ C.L.})$$

Search for Direct CP Violation

$$A_{CP} \equiv \frac{\mathcal{B}(B^- \rightarrow \psi^0 K^-) - \mathcal{B}(B^+ \rightarrow \psi^0 K^+)}{\mathcal{B}(B^- \rightarrow \psi^0 K^-) + \mathcal{B}(B^+ \rightarrow \psi^0 K^+)} = \frac{b - \bar{b}}{b + \bar{b}}$$

- We can measure $A_{CP}(\psi K^0)$ with 4% precision
- In Standard Model $A_{CP}(\psi K^0) \ll 4\%$ (even if penguin amplitude is large)
- $A_{CP} = \frac{-2A_1 A_2 \sin(\delta_1 - \delta_2) \sin(\phi_1 - \phi_2)_{\text{weak}}}{A_1^2 + A_2^2 + 2A_1 A_2 \cos(\delta_1 - \delta_2) \cos(\phi_1 - \phi_2)}$



- $A(\psi K^0) = V_{cb}V_{cs}^*(F_{\psi K^0} + P_c - P_{\bar{c}}) + V_{ub}V_{us}^*(P_c - P_{\bar{c}})$
- $\arg(V_{cb}V_{cs}^*/V_{ub}V_{us}^*) \simeq \lambda^2 \eta + \pi$
($\lambda = 0.22, \eta \leq 1$)

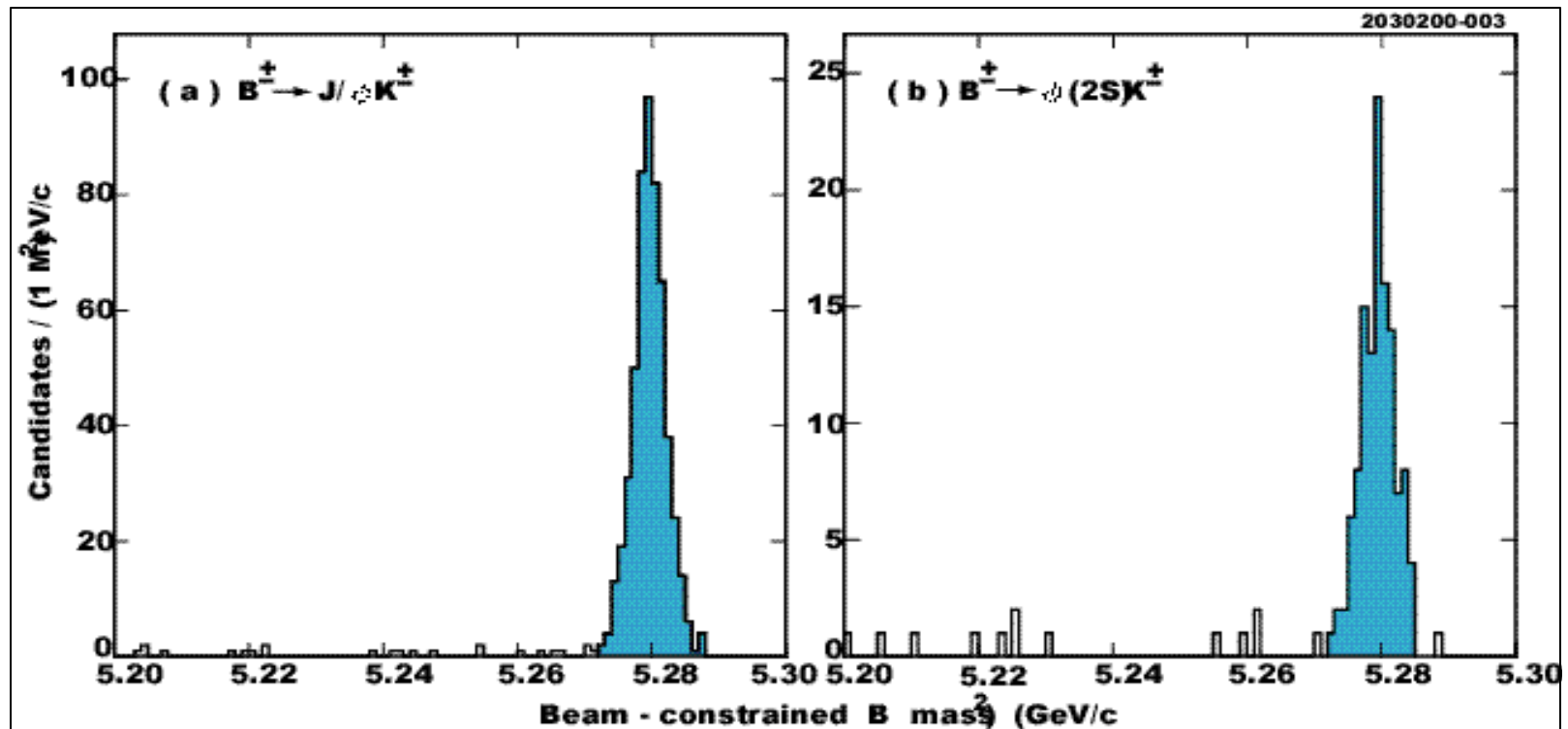
One of the models beyond the Standard

Two-Higgs doublet model with special status for top quark

(Gross, Soun, and Wu, hep-ph/9911419)

- H^- -mediated diagram competes with SM tree W^- -mediated diagram and comes with its own CP-odd phase
- $A_{CP}(\psi K^0)$ could be $\mathcal{O}(10\%) \implies$ we can measure it right now

Search for Direct CP Violation



Mode	$N(B^+)$	$N(B^-)$	$N(B^0)$	$\frac{N(B^+) - N(B^-)}{N(B^+) + N(B^-)}$	A_{CP}
$B^+ \rightarrow J/\psi K^+$	554	271	263	$(+1.5 \pm 4.3)\%$	$(+1.8 \pm 4.3 \pm 0.4)\%$
$B^+ \rightarrow \phi(2S) K^+$	120	61	59	$(+1.7 \pm 9.1)\%$	$(+2.0 \pm 9.1 \pm 1.0)\%$

CP Asymmetries

- Measure B and \bar{B} reactions described by two amplitudes:
 - $\Gamma(B \rightarrow f) = | a_1 e^{i(f_1 + d_1)} + a_2 e^{i(f_2 + d_2)} |^2$
 - $\overline{\Gamma(B \rightarrow f)} = | a_1 e^{i(-f_1 + d_1)} + a_2 e^{i(-f_2 + d_2)} |^2$
- CP asymmetry from strong and weak phase differences
 - $\Delta \equiv (\Gamma - \overline{\Gamma}) / (\Gamma + \overline{\Gamma}) \propto \sin(f_1 - f_2) \sin(d_1 - d_2)$
- Depends upon comparable magnitudes as well
- CLEO can measure decays that are sensitive to $g = \arg(V_{ub}^*)$
 - $B^+ \rightarrow K^+ p^0, B^+ \rightarrow K^0 p^+, B^+ \rightarrow K^+ h'$

A_{CP} Expectations

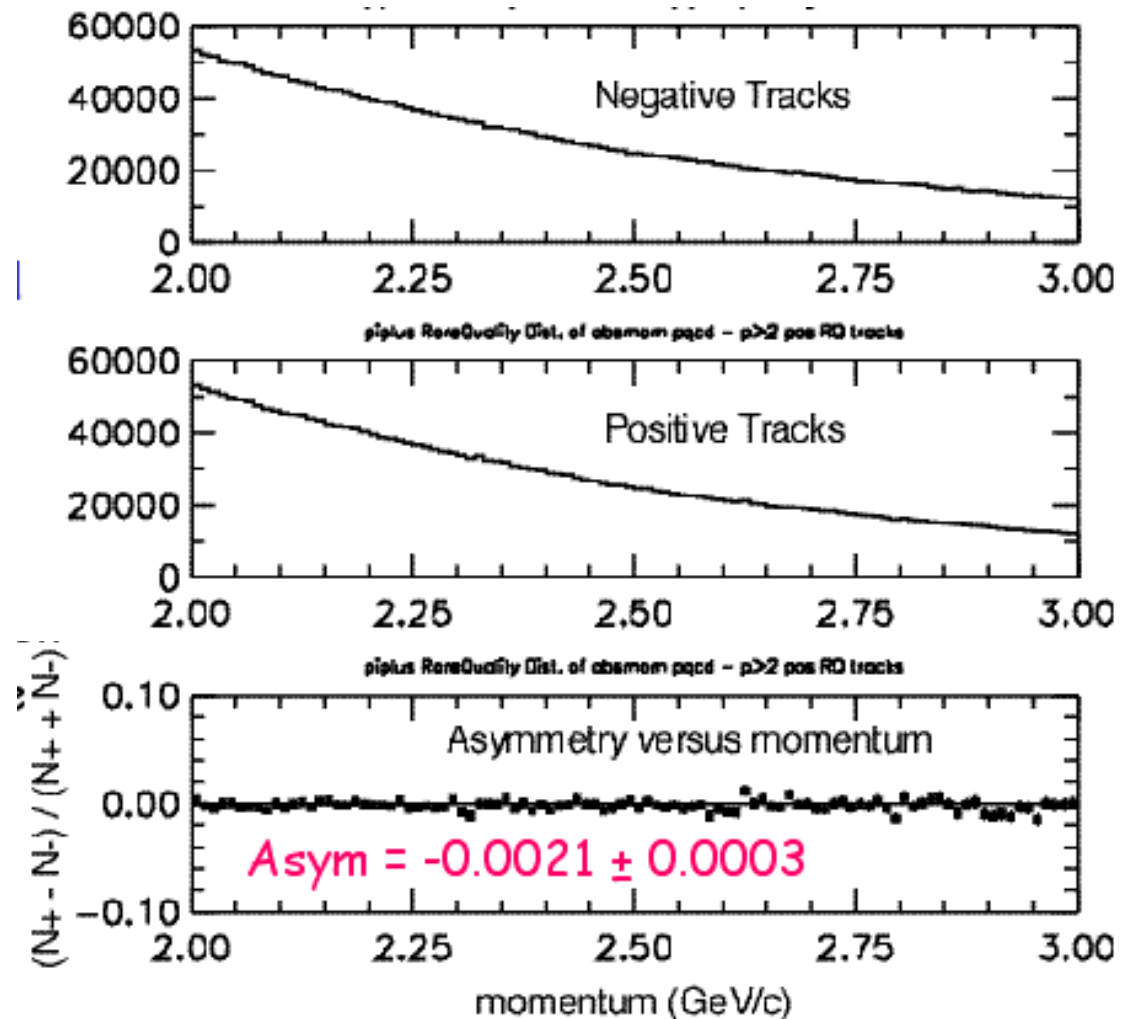
- Factorization model calculations (no FSI interactions)
Ali, Kramer, Liu, hep-ph/9805403

» K^+p^+	0.04 - 0.11	» K^+p^0	0.03 - 0.09
» K^0p^+	0.01	» K^+h'	0.02 - 0.06
» wp^+	-0.12 - +0.02		

- Final state interactions may boost $A_{CP} \sim 20 - 40\%$.
 - He et al, *Phys. Rev. Lett.* **81**, 5738 (1998)
 - Neubert, *JHEP* 9902, 014 (1999)
 - Deshpande et al., *Phys. Rev. Lett.* **82**, 2240 (1999)
- New physics could boost $A_{CP} \sim 40 - 60\%$.
 - He et al., hep-ph/980982

Experimental Bias(es)

- B flavor tagged by high momentum track
- Must demonstrate reconstruction not charge dependent.
- Charge difference in K^-N and K^+N cross sections
- Track reconstruction difference confirmed in Monte Carlo ~ 0.002



CP Asymmetry Results

80.2 \pm 11.8 events
- 11.0

42.1 \pm 10.9 events
- 9.9

25.2 \pm 6.4 events
- 5.6

101 \pm 13 events
- 12

28.5 \pm 8.2 events
- 7.3

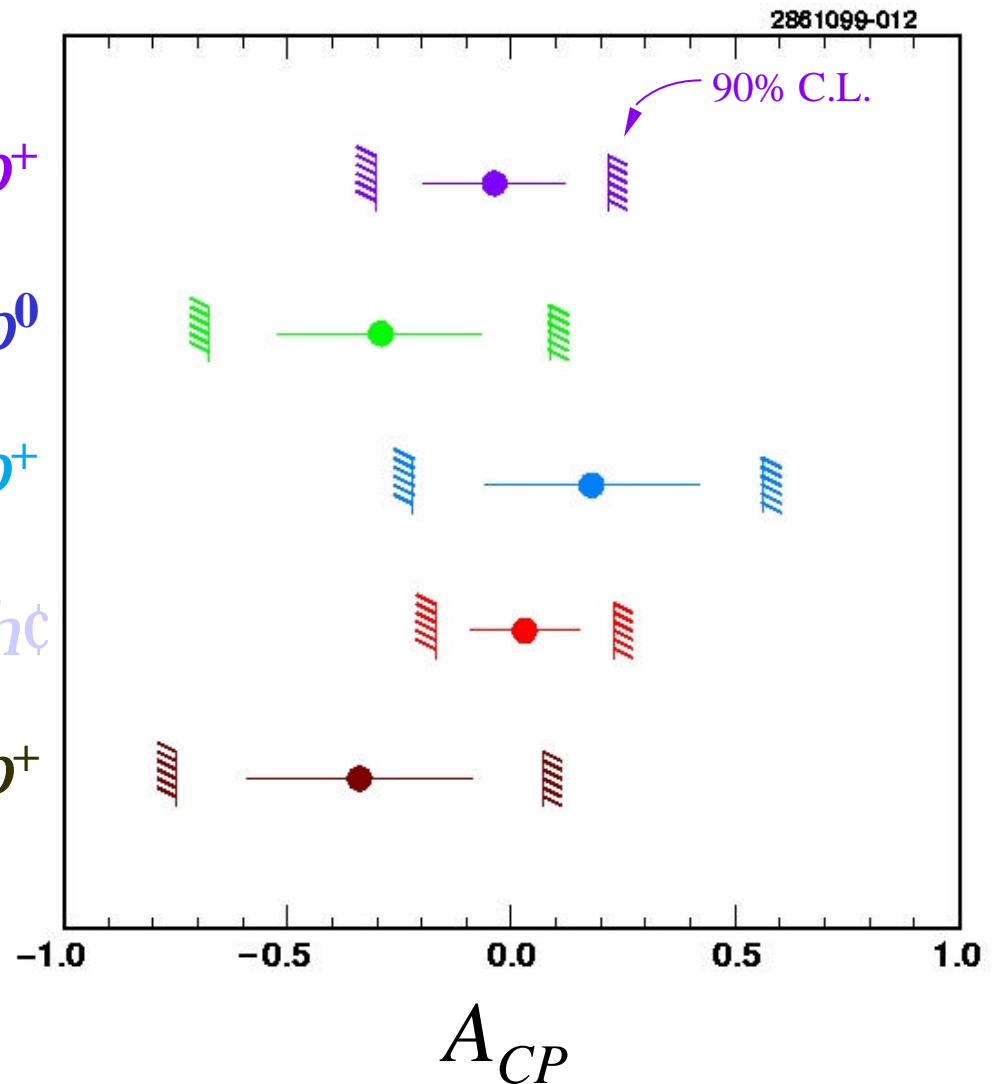
$K^- p^+$

$K^+ p^0$

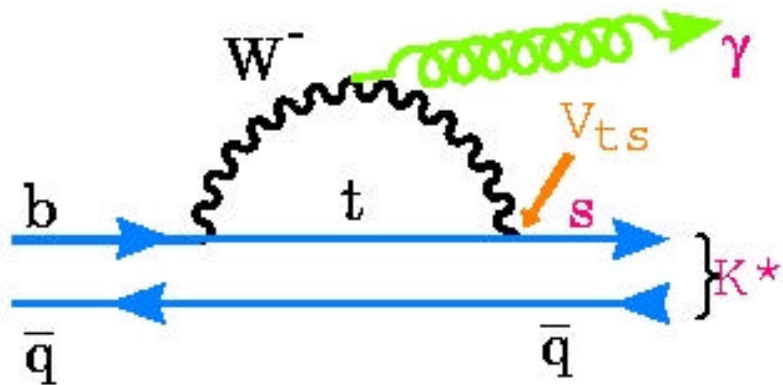
$K^0 p^+$

$K^+ h_c$

wp^+

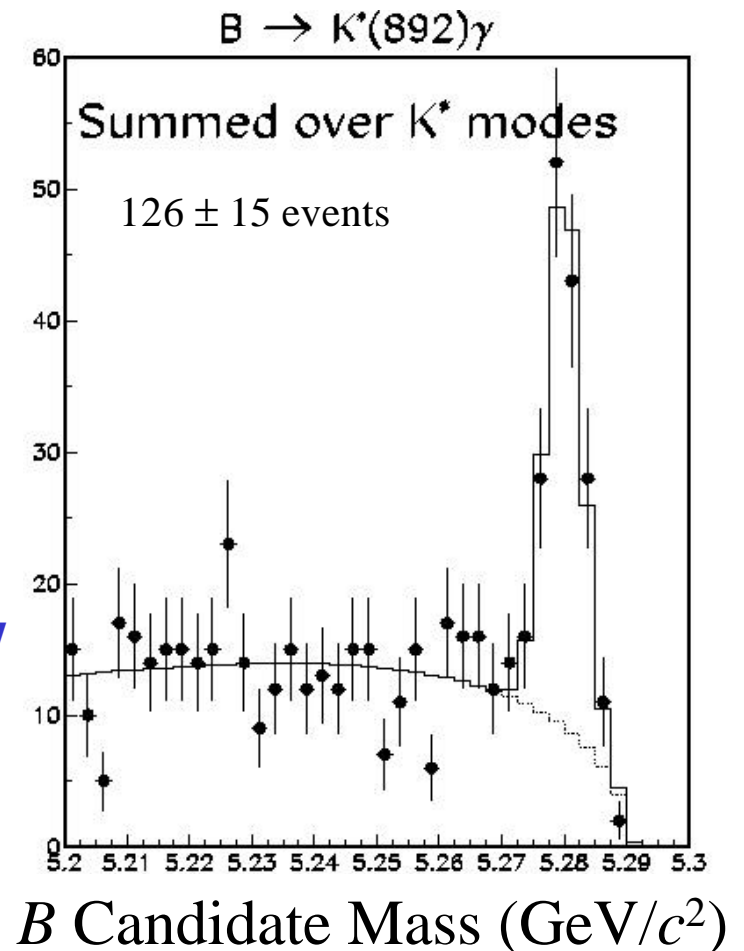


~~CP~~ from New Physics?



- Penguin amplitude $\propto |V_{ts}|$
- Other amplitudes, ~~CP~~, small in SM
- Some Higgs models introduce ~~CP~~, possibly even if $b \rightarrow sg$ rate unaffected.

- Wolfenstein & Wu, *Phys. Rev. Lett.* **73**, 2809 (1998)
- Asatrian & Ioannissian, *Phys. Rev.* **D54**, 5642
- Kagan & Neubert, *Phys. Rev.* **D58**, 094012



$b \rightarrow sg$ Results

- Updated branching ratio results:

$$\text{BR}(B^0 \rightarrow K^{*0} g) = (4.5 \pm 0.7 \pm 0.3) \times 10^{-5}$$

$$\text{BR}(B^+ \rightarrow K^{*+} g) = (3.8 \pm 0.9 \pm 0.3) \times 10^{-5}$$

$$\text{BR}(B \rightarrow K_2^{*}(1430) g) = (1.6 \pm 0.55 \pm 0.13) \times 10^{-5}$$

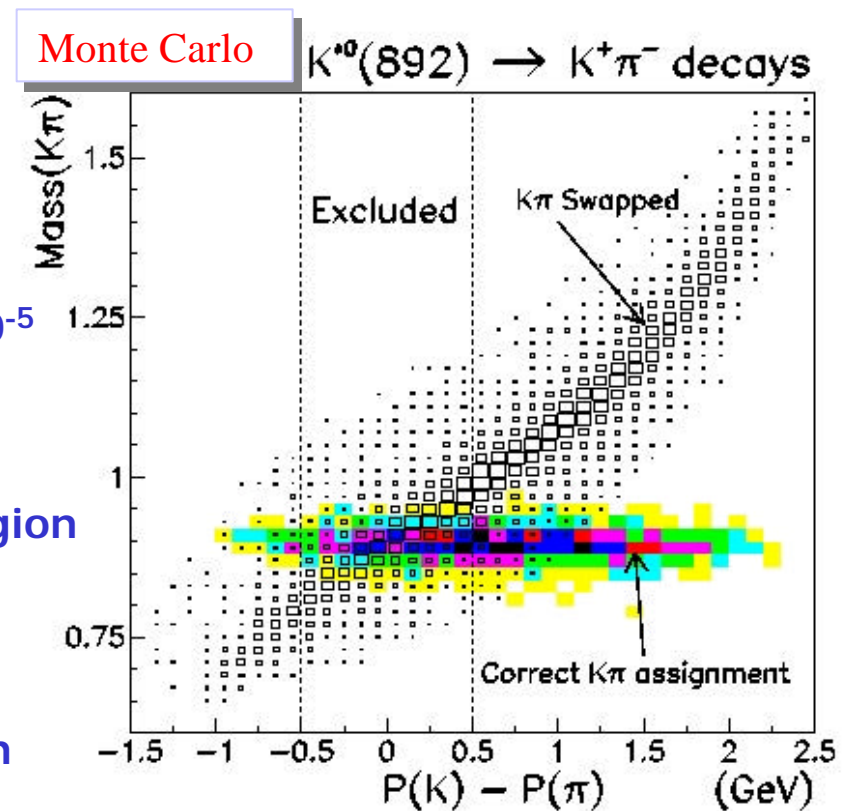
- CP asymmetry from special kinematic region for best K/p identification

$$A_{CP} = 0.08 \pm 0.13 \pm 0.03$$

- Asymmetry for inclusive $b \rightarrow sg$ (based on 9.7M $B\bar{B}$ pairs only):

$$A_{CP} = -0.063 \pm 0.090 \pm 0.009 \text{ or}$$

$$-0.21 < A_{CP} < 0.09 \text{ (90\% C.L.)}$$



- Upper limit on $b \rightarrow d$ exclusive penguins:

$$\text{BR}(B \rightarrow (r, w) g) < \sim 10^{-5}$$

Search for $b \rightarrow dg$

- Expect that $B \rightarrow rg$ also described by penguin amplitude - dominant top?

$$\frac{G(B \rightarrow rg)}{G(B \rightarrow K^*g)} = \frac{|V_{td}|^2}{|V_{ts}|^2} \times \alpha \quad \alpha \sim 0.6 - 0.8$$

- Updated branching ratio limits:

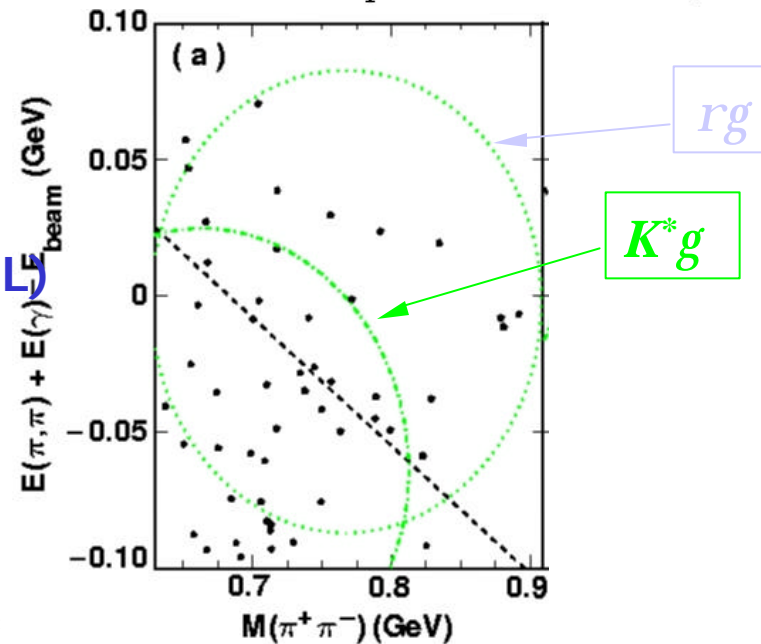
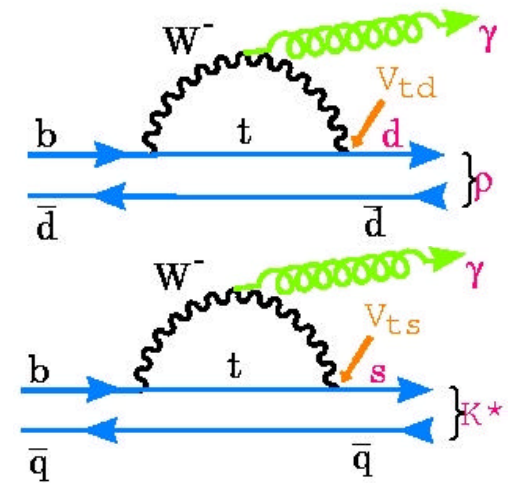
$$BR(B^0 \rightarrow r^0g) < 1.7 \cdot 10^{-5}$$

$$BR(B^+ \rightarrow r^+g) < 1.3 \cdot 10^{-5}$$

$$BR(B^0 \rightarrow wg) < 9.2 \cdot 10^{-6}$$

- $BR(B \rightarrow rg) / BR(B \rightarrow K^*g) < 0.32$ (90%CL)

- $|V_{td}/V_{ts}| < 0.72$ (90%CL)



Short Term Future

- **B Factories (BABAR, Belle, (CLEO))**
 - $\sin(2\beta)$ with $\sim 5\text{-}10\%$ precision
 - Rare decays with BR $\sim 10^{-6}$
 - Better understanding of V_{cb} and V_{ub}
- **Hadron colliders**
 - HERA-b has the potential to also measure $\sin(2\beta)$
 - CDF already has already taken the first steps toward measuring $\sin(2\beta)$. Next run will start \sim Aug. 2000.
 - CDF could also measure x_s .

Why do b & c decay physics at hadron colliders?

- Large samples of b quarks are available, with the Fermilab Main Injector, the collider will produce $\sim 4 \times 10^{11}$ b hadrons per 10^7 sec at $L = 2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$.
- e^+e^- machines operating at the Y(4S) at L of 3×10^{33} produce 6×10^7 B's per 10^7 s.
- B_s & L_b and other b-flavored hadrons are accessible for study.
- Charm rates are $\sim 10x$ larger than the b rate

Main detector challenges

- **Problems:**

- $\sigma_b/\sigma_{\text{tot}} \sim 1/500$ at Fermilab, $1/100$ at LHC
- Background from b's can overwhelm "rare" processes
- Large data rate just from b's - 1 kHz into detector
- Large rates cause Radiation damage to EM calorimeter; photon multiplicities may obscure signals

- **Solutions for BTeV:**

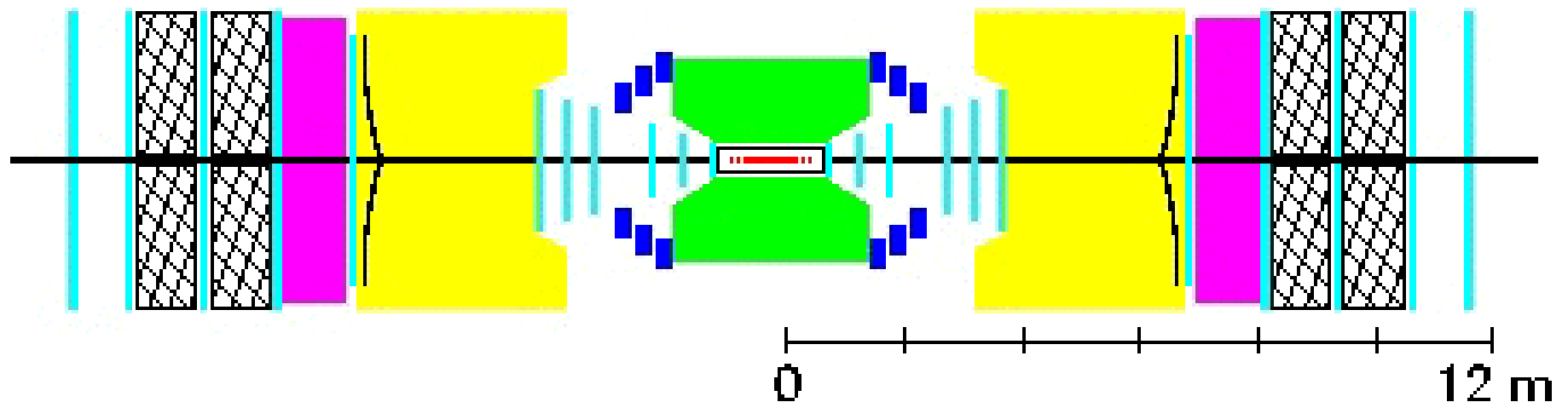
- Use **detached vertices** for trigger and background rejection
- Have excellent charged **particle identification** & lepton id
- Dead-timeless trigger and DAQ system capable of writing kHz of events to tape
- Use PbWO_4 crystal calorimeter

Fundamental Detector Principles

- Necessary to trigger efficiently on purely hadronic final states - detached vertex trigger
- Necessary to reconstruct final states with excellent decay time resolution, good efficiency and mass resolution
- Necessary to detect final states with g or p^0 efficiently with good energy resolution
- Necessary to be able to identify $p/K/p$

The BTeV Detector

- Pixel Detector
Inside the beam pipe
- SM3 Dipole Magnet
- Magnet Coils
- Beam Pipe Vacuum
- Wire Chambers
Aperature $\tan\theta=0.3$
- RICH
- EM Cal-PbWO₄ crystals
- Muon detector

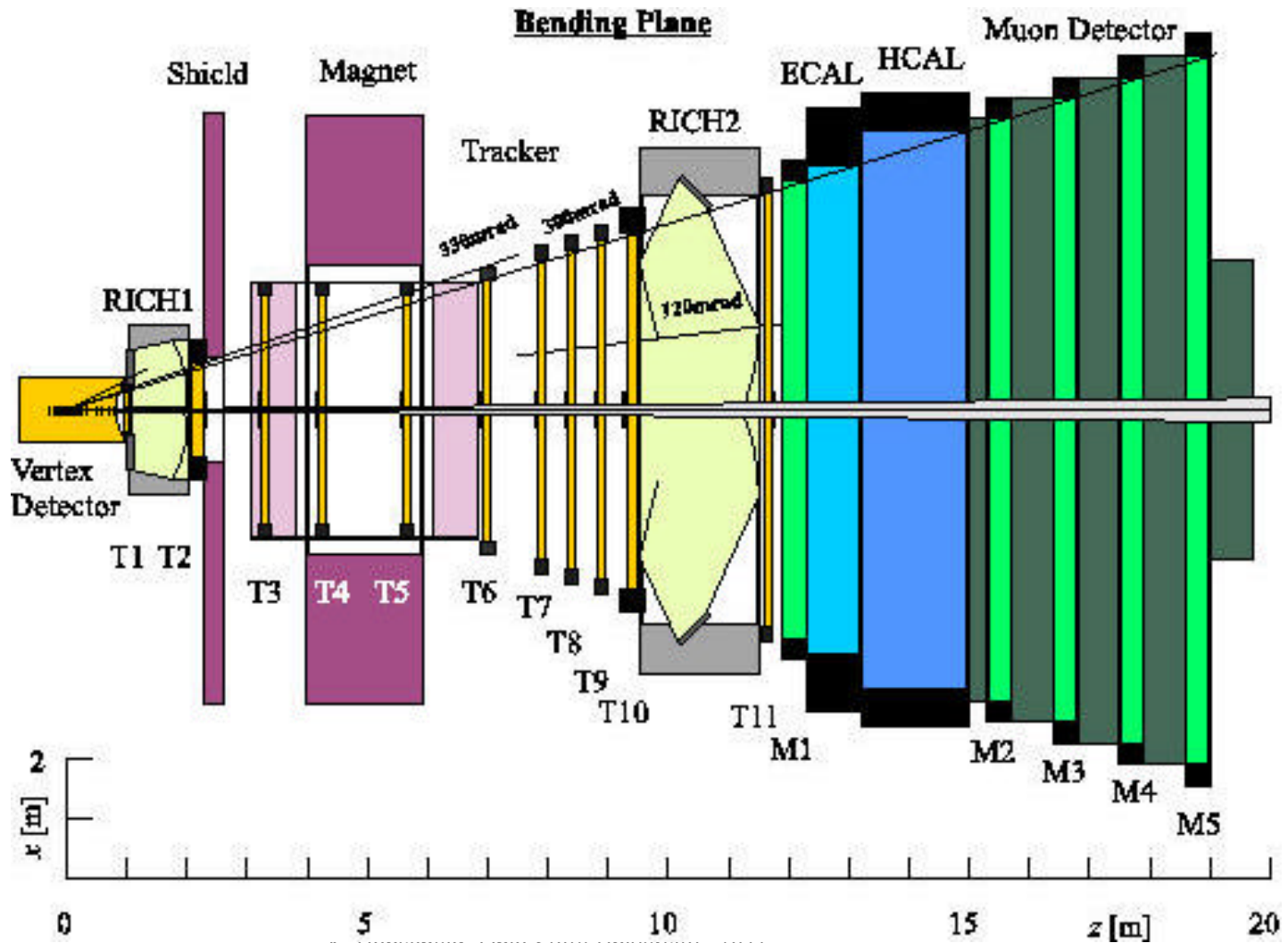


The C0 Interaction Region

- Construction finished
- BTeV is designed to be compatible



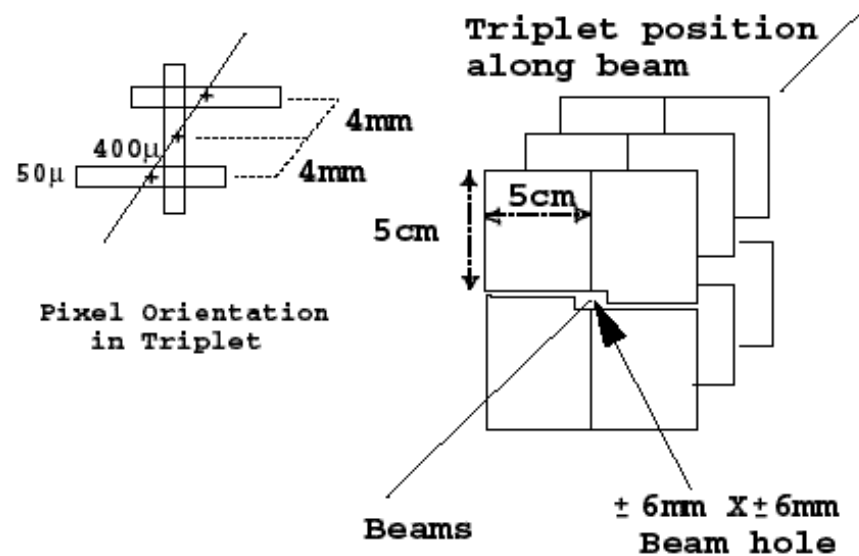
The LHCb Detector



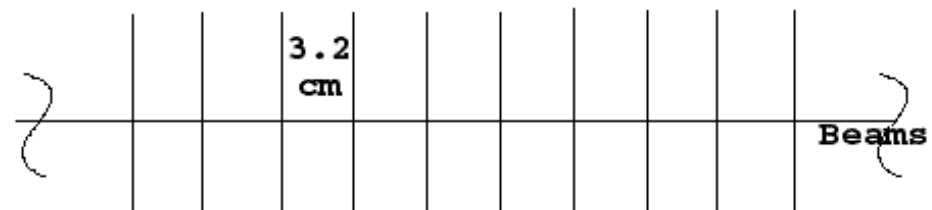
The BTeV Pixel Detector

- **Pixels necessary to eliminate ambiguity problems with high track density; Essential to our detached vertex trigger**
- **Crucial for accurate decay length measurement**
- **Radiation hard**
- **Low noise**

The BTeV Baseline Pixel Detector

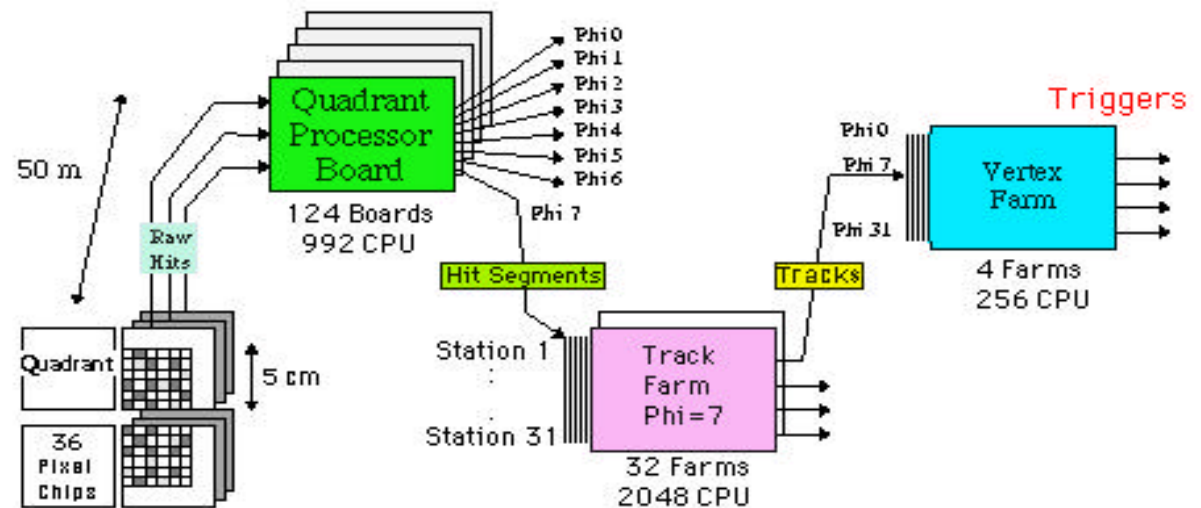


Elevation View
10 of 31 Triplet Stations



Pixel Trigger Description

- Triplets used to get space point & mini-vector, called a 'station hit'
- Station hits are organized
- into f-slices
- Tracks are found in these f-slices
 - full pattern recognition is performed
 - Minimum track p cu
- Event level processor



Detached Vertex Trigger

- Level I Trigger uses information from the Pixel Detector to find the primary vertex and then look for tracks that are detached from it
- The simulation does the pattern recognition. It uses hits from MCFast including multiple scattering, bremsstrahlung, pair conversions, hadronic interactions and decays in flight
- Detailed studies of efficiency and rejection for up to an average of three interactions/crossing

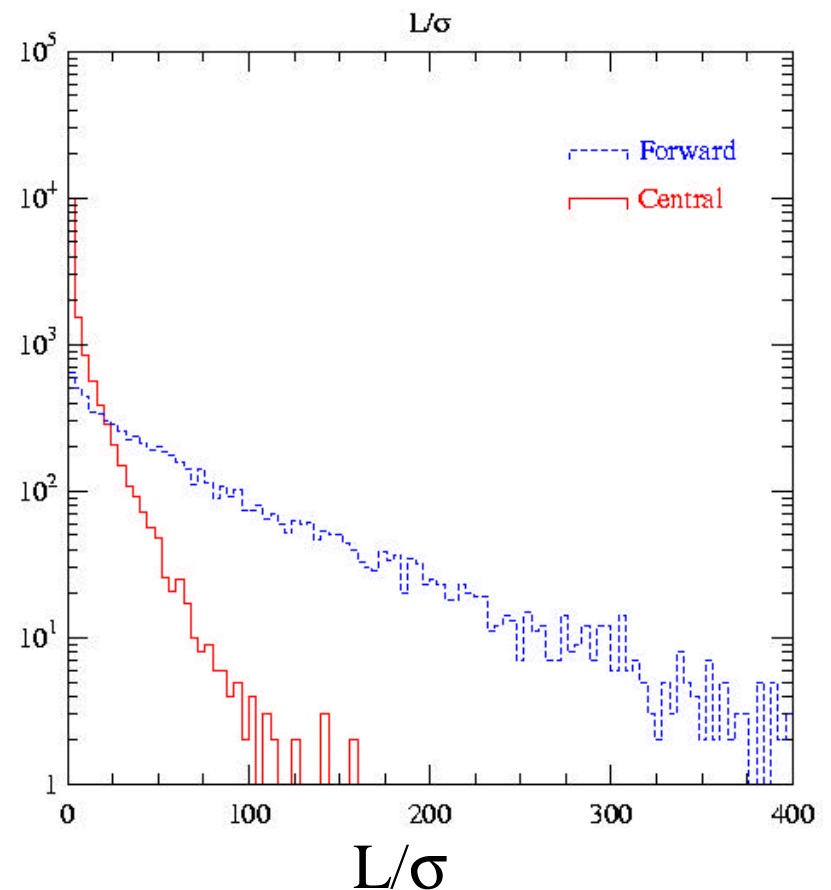
BTeV Trigger Performance

- For a requirement of at least 2 tracks detached by more than 4s, BTeV triggers on only 1% of the beam crossings and achieve the following efficiencies for these states:

State	efficiency(%)	State	efficiency(%)
$B \rightarrow \pi^+\pi^-$	55	$B^0 \rightarrow K^+\pi^-$	54
$B_s \rightarrow D_s K$	70	$B^0 \rightarrow J/\psi K_s$	50
$B^- \rightarrow D^0 K^-$	60	$B_s \rightarrow J/\psi K^*$	69
$B^- \rightarrow K_s \pi^-$	40	$B^0 \rightarrow K^* \gamma$	40

$B^0 \otimes p^+ p^-$: L/s distribution

- L/s = Decay length/error is very important in rejecting background both at trigger level and in analysis
- Much better in Forward (BTeV) geometry than **Central geometry** because b's are moving faster



A sample calculation: $B^0 \text{ @ } p^+p^-$

	<u>BTeV</u>	<u>LHCb</u>
Cross-section	100 μb	500 μb
Luminosity	2×10^{32}	2×10^{32}
# of B^0 /Year (10^7 s)	1.4×10^{11}	7×10^{11}
$B(B^0 \rightarrow \pi^+\pi^-)$	0.75×10^{-5}	0.75×10^{-5}
Reconstruction efficiency	0.06	0.032
Triggering efficiency (after all other cuts)	0.50	0.17
# ($\pi^+\pi^-$)	34,000	28,560
ϵD^2 for flavor tags (K^\pm, λ^\pm , same + opposite side jet tags)	0.1	0.1
# of tagged $\pi^+\pi^-$	3,400	2,900
Signal/Background	0.6	1
Error in $\pi^+\pi^-$ asymmetry (including bkgrd)	± 0.023	± 0.019

Comparisons of BTeV With e^+e^- B factories

- Number of flavor tagged $B^0 \rightarrow p^+ p^-$ ($B=0.75 \times 10^{-5}$)

	L ($\text{cm}^{-2}\text{s}^{-1}$)	σ	$\#B^0/10^7\text{s}$	ϵ	ϵD^2	#tagged
e^+e^-	3×10^{33}	1nb	3.0×10^7	0.4	0.4	46
BTeV	2×10^{32}	100 μb	1.4×10^{11}	0.03	0.1	3400

- Number of $B^- \rightarrow D^0 K^-$

	L ($\text{cm}^{-2}\text{s}^{-1}$)	σ	$\#B^0/10^7\text{s}$	ϵ	#
e^+e^-	3×10^{33}	1nb	3.0×10^7	0.5	2
BTeV	2×10^{32}	100 μb	1.4×10^{11}	0.015	320

- B_s , B_c and L_b not done at $Y(4S)$ e^+e^- machines
- Number of tagged, reconstructed B^0 decays to rp is a factor of at least 10 higher for BTeV.

Comparisons of BTeV with LHCb

- **Advantages of LHCb**

- σ_b 5x larger at LHC, while σ_t is only 1.6x larger
- The mean number of interactions per beam crossing is 3x lower at LHC, when the FNAL bunch spacing is 132 ns
- LHCb HAS BEEN APPROVED!

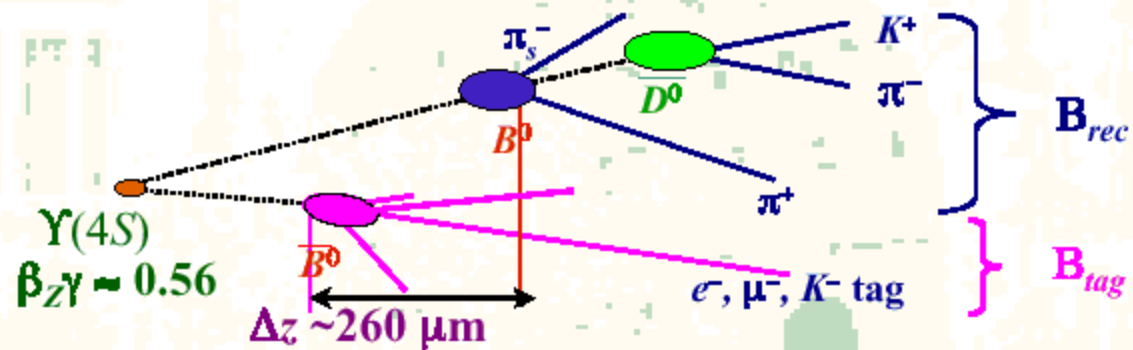
- **Advantages of BTeV (machine specific)**

- The 25 ns bunch spacing at LHC makes 1st level detached vertex triggering more difficult.
- The 7x larger LHC beam energy causes problems: much larger range of track momenta that need to be analyzed and large increase in track multiplicity, which causes triggering and tracking problems
- The long interaction region at FNAL, $\sigma=30$ cm compared with 5 cm at LHC, somewhat compensates for the larger number of interactions per crossing, since the interactions are well separated

Comparisons with LHCb II

- **Advantages of BTeV (detector specific)**
 - BTeV is a two-arm spectrometer (gives 2x advantage)
 - BTeV has vertex detector in magnetic field which allows rejection of high multiple scattering (low p) tracks in the trigger
 - BTeV is designed around a pixel vertex detector which has much less occupancy, and allows for a detached vertex trigger in the first trigger level.
 - Important for accumulation of large samples of rare hadronic decays and charm physics.
 - Allows BTeV to run with multiple interactions per crossing, L in excess of $2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
 - BTeV will have a much better EM calorimeter

Babar and Belle time measurement:

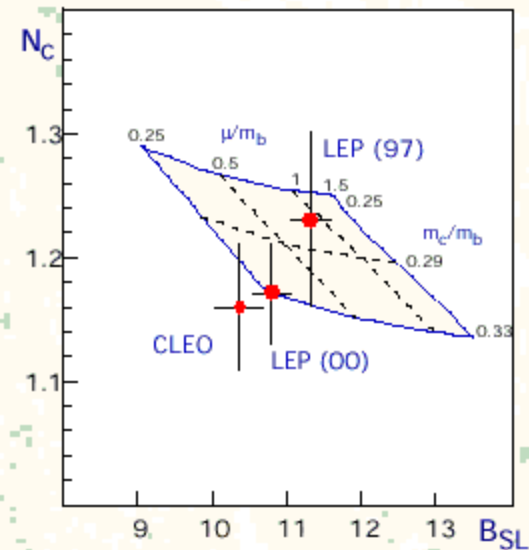


The time resolution is about $\tau_B/2$.

The primary vertex is not used (z_V is poorly known),
 except in dilepton analyses.
 One B decay defines $t = 0$.

Charm Counting

It is possible to calculate both the inclusive semileptonic BR and the number of charmed particles per decay in terms of fundamental quantities. In the past, the measurements have disagreed, making interpretation difficult. New **ALEPH** & **DELPHI** results change the picture. There is now experimental consistency.



$$N_c = 1.171 \pm 0.040$$

$$B_{SL} = 10.79 \pm 0.25$$

(Barker & Blyth @ ICHEP)

Using new lifetimes

$$\text{CLEO: } B_{SL} = (10.49 \pm 0.17 \pm 0.43)\%$$

$$N_c = 1.10 \pm 0.05$$

Note: LEP Vcb WG says: $B_{SL} = 10.56 \pm 0.11 \pm 0.18$