

Last time

Separation of Variables in Cylindrical Coordinates

(cont'd)

A. z-independent geometry

$$\Phi(\rho, \varphi) = R(\rho) \Psi(\varphi)$$

$$\Rightarrow \begin{cases} R(\rho) = a\rho^\nu + b\rho^{-\nu} \\ \Psi(\varphi) = A \cos(\nu\varphi) + B \sin(\nu\varphi) \end{cases}, \nu \neq 0$$

$$\begin{cases} R(\rho) = a_0 + b_0 \ln \rho \\ \Psi(\varphi) = A_0 + B_0 \varphi \end{cases}, \nu = 0$$

for $\varphi \in [0, 2\pi]$ $\Rightarrow \nu = \text{integer}$

\Rightarrow

$$\Phi(\rho, \varphi) = a_0 + b_0 \ln \rho + \sum_{n=1}^{\infty} \left[a_n \rho^n \sin(n\varphi + \alpha_n) + b_n \rho^{-n} \sin(n\varphi + \beta_n) \right]$$

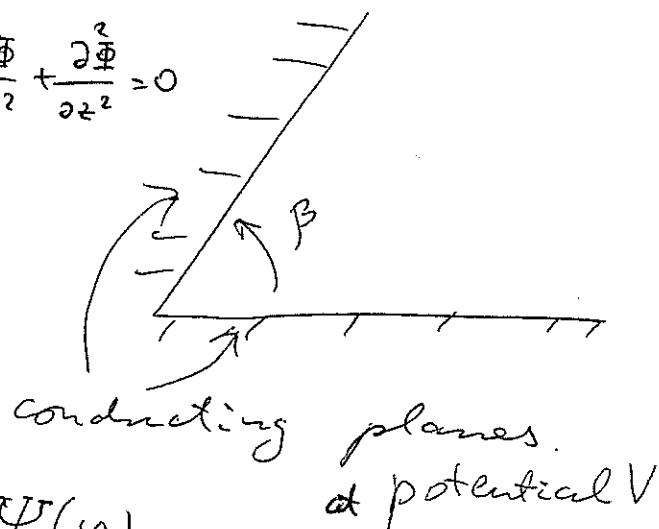
general solution of Laplace equation
in cylindrical coordinates in the z-indep.
geometry

A. z-independent geometry.

$$\nabla^2 \Phi = 0 \Rightarrow \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \varphi^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Phi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \varphi^2} = 0$$

if $\Phi = \Phi(\rho, \varphi)$



Try separation of

variables: $\Phi(\rho, \varphi) = R(\rho) \Psi(\varphi)$

\Rightarrow multiplying Laplace equation by $\frac{\rho^2}{\Phi}$ we get

$$\frac{\rho}{R} \frac{d}{d\rho} \left(\rho \frac{dR}{d\rho} \right) + \frac{1}{\Psi} \Psi'' = 0$$

$$\underbrace{\frac{\rho}{R} \frac{d}{d\rho} \left(\rho \frac{dR}{d\rho} \right)}_{= \nu^2} + \underbrace{\frac{1}{\Psi} \Psi''}_{= -\nu^2} = 0$$

$$\Rightarrow \frac{\rho}{R} \frac{d}{d\rho} \left(\rho \frac{dR}{d\rho} \right) = \nu^2, \quad \frac{\Psi''}{\Psi} = -\nu^2$$

$$\Rightarrow \rho^2 R'' + \rho R' - \nu^2 R = 0$$

$$R(\rho) = \rho^\lambda \Rightarrow \lambda(\lambda-1) + \lambda - \nu^2 = 0 \Rightarrow \lambda = \pm \nu$$

$$\Rightarrow \int R(\rho) = a \rho^\nu + b \rho^{-\nu}$$

$$\int \Psi(\varphi) = A \cos(\nu \varphi) + B \sin(\nu \varphi)$$

for $\nu = 0$:
$$\begin{cases} R(\rho) = a_0 + b_0 \ln \rho \\ \Psi(\varphi) = A_0 + B_0 \varphi \end{cases}$$

(151)

General solutions of problems with cylindrical symmetry & z -independence!
 for $\varphi \in [0, 2\pi]$

$$\Phi(\rho, \varphi) = a_0 + b_0 \ln \rho + \sum_{n=1}^{\infty} [a_n \rho^n \sin(n\varphi + \alpha_n) + b_n \rho^{-n} \sin(n\varphi + \beta_n)]$$

For our problem, we want $\Phi(\varphi=0) = \Phi(\varphi=\beta) = V$

for all $\rho \Rightarrow b_0 = 0, B_0 = 0, A_0 = 0$
 $\Rightarrow a_0 A_0 = V$

no singularity at $\rho = 0 \Rightarrow b = 0$

Also, at $\varphi = \beta$: $\sin(\nu\beta) = 0 \Rightarrow \nu = \frac{\pi m}{\beta}, m \in \text{int.}$

$$\Rightarrow \Phi(\rho, \varphi) = V + \sum_{m=1}^{\infty} a_m \rho^{\frac{m\pi}{\beta}} \sin\left(\frac{m\pi}{\beta} \varphi\right)$$

a_m 's to be fixed by b.c.'s at ∞

\Rightarrow need the field near $\rho \approx 0 \Rightarrow m > 0 \Rightarrow$

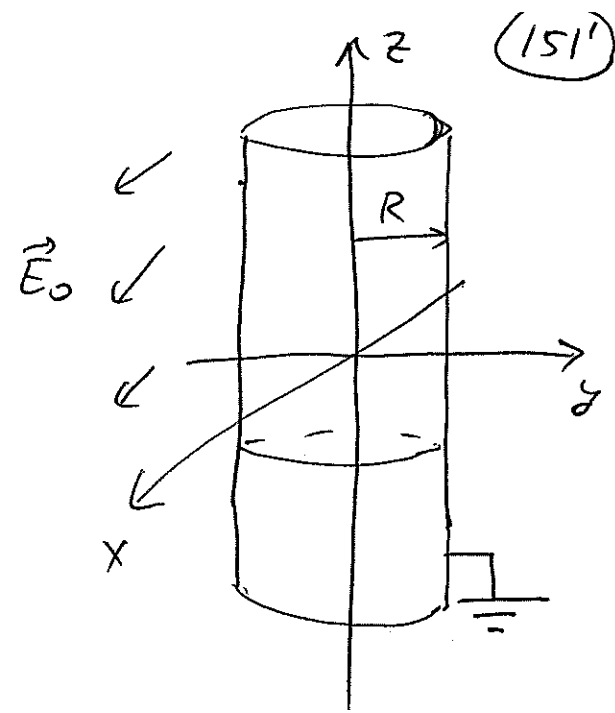
only $m=1$ can be kept \Rightarrow

$$\Phi(\rho, \varphi) \approx V + a_1 \rho^{\frac{\pi}{\beta}} \sin\left(\frac{\pi}{\beta} \varphi\right) \text{ for small-}\rho$$

$$\Rightarrow E_\rho = -\frac{\partial \Phi}{\partial \rho} = -\frac{\pi a_1}{\beta} \rho^{\left(\frac{\pi}{\beta}-1\right)} \sin\left(\frac{\pi}{\beta} \varphi\right)$$

Example | $\vec{E}_0 = E_0 \hat{x}$

Infinite grounded conducting cylinder of radius R in a uniform external electric field \vec{E}_0 . Find Φ .



General solution:

$$\Phi(\rho, \varphi) = a_0 + b_0 \ln \rho + \sum_{n=1}^{\infty} \left[a_n \rho^n \sin(n\varphi + \alpha_n) + b_n \rho^{-n} \sin(n\varphi + \beta_n) \right]$$

(i) As $\rho \rightarrow \infty$ \Rightarrow ^{should} get $\Phi_0 = -E_0 x = -E_0 \rho \cos \varphi$

\Rightarrow require that $\Phi(\rho, \varphi) \Big|_{\rho \rightarrow \infty} = -E_0 \rho \cos \varphi$

$\Rightarrow b_0 = 0, a_0 = 0, a_n = 0$ for $n \neq 1,$

$a_1 = -E_0, \quad \alpha_1 = \frac{\pi}{2}$ (as $\sin(\varphi + \frac{\pi}{2}) = \cos \varphi$)

$\Rightarrow \Phi(\rho, \varphi) = -E_0 \rho \cos \varphi + \sum_{n=1}^{\infty} b_n \rho^{-n} \sin(n\varphi + \beta_n)$

(ii) Cylinder is grounded $\Rightarrow \Phi(\rho=R, \varphi) = 0$

$\Rightarrow -E_0 R \cos \varphi + \sum_{n=1}^{\infty} b_n R^{-n} \sin(n\varphi + \beta_n) = 0$

$$\Rightarrow b_n = 0 \text{ for } n \neq 1$$

$$-E_0 R \cos \varphi + \frac{b_1}{R} \sin(\varphi + \beta_1) = 0$$

$$\Rightarrow \beta_1 = \frac{\pi}{2}, \quad b_1 = E_0 R^2$$

\Rightarrow we obtain

$$\Phi(\rho, \varphi) = -E_0 \rho \cos \varphi + \frac{E_0 R^2}{\rho} \cos \varphi$$

$$\Rightarrow \Phi(\rho, \varphi) = -E_0 \cos \varphi \left[\rho - \frac{R^2}{\rho} \right]$$

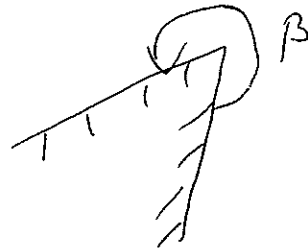
$$E_\varphi = -\frac{1}{\rho} \frac{\partial \Phi}{\partial \varphi} = -\frac{\bar{n} a_1}{\beta} \rho^{\left(\frac{\bar{n}}{\beta}-1\right)} \cos\left(\frac{\bar{n}}{\beta} \varphi\right).$$

$\sigma(\rho)$ at $\varphi=0$ is $\sigma = \epsilon_0 E_\varphi(\rho, 0) \Rightarrow$

$$\sigma = -\frac{\epsilon_0 \bar{n} a_1}{\beta} \rho^{\frac{\bar{n}}{\beta}-1} \rightarrow \infty \text{ as } \rho \rightarrow 0$$

$\text{if } \beta > \bar{n}$

strong fields lead to
electrical discharge \Rightarrow



\Rightarrow leads to lightning "hitting" lightning rod!

B. z-dependent case.

Laplace equation becomes $\frac{\partial^2 \Phi}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial \Phi}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \varphi^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$

$$\Phi(\rho, \varphi, z) = R(\rho) Q(\varphi) Z(z)$$

$$\Rightarrow R'' Q Z + \frac{1}{\rho} R' Q Z + \frac{1}{\rho^2} R Q'' Z + R Q Z'' = 0$$

divide by $R Q Z$ to obtain

$$\left\{ \begin{array}{l} \frac{d^2 Z}{dz^2} - k^2 Z = 0 \\ \frac{d^2 Q}{d\varphi^2} + \nu^2 Q = 0 \end{array} \right.$$

$$\left(\frac{d^2 R}{d\rho^2} + \frac{1}{\rho} \frac{dR}{d\rho} + \left(k^2 - \frac{\nu^2}{\rho^2} \right) R = 0. \right.$$

First two are easy: $Z = e^{\pm k z}$, $Q = e^{\pm i\nu \varphi}$

Last one is a bit trickier: first rescale: $x = k\rho$

$$\Rightarrow \frac{d^2 R}{dx^2} + \frac{1}{x} \frac{dR}{dx} + \left(1 - \frac{\nu^2}{x^2} \right) R = 0$$

Look for solution in the form of a series:

$$R(x) = x^\alpha \sum_{j=0}^{\infty} a_j x^j, \quad a_0 \neq 0$$

$$\Rightarrow \sum_{j=0}^{\infty} \left[a_j \cdot (j+\alpha)(j+\alpha-1) x^{j-2+\alpha} + a_j \cdot (j+\alpha) \cdot \right.$$

$$\left. \cdot x^{\alpha+j-2} - \nu^2 a_j x^{\alpha+j-2} \right] + \sum_{j=0}^{\infty} a_j x^{j+\alpha} = 0$$

$j=0$: $a_0 (\alpha^2 - \nu^2) \cdot x^{\alpha-2} = 0 \Rightarrow$ if $a_0 \neq 0 \Rightarrow \alpha = \pm \nu$. ↙ 1st non-vanishing coef.

$j=1$: $a_1 [(\alpha+1)^2 - \nu^2] x^{\alpha-1} = 0 \Rightarrow a_1 = 0.$

other j 's: $a_{j+2} [(j+\alpha+2)^2 - \nu^2] + a_j = 0$ | coef. at x^{j+2}

$$\Rightarrow a_{j+2} = -a_j \frac{1}{(j+2)^2 + 2\alpha(j+2)}$$