

Last time

Monochromatic Plane Wave ~~(2011)~~

energy density:

$$u = \frac{1}{2} (\vec{E} \cdot \vec{D} + \vec{B} \cdot \vec{H})$$

Poynting vector:

$$\vec{S} = \vec{E} \times \vec{H}$$

(energy flow)

Maxwell stress tensor:

$$\sigma^{ij} = \epsilon E^i E^j + \mu H^i H^j - \frac{1}{2} \delta^{ij} (\epsilon E^2 + \mu H^2)$$

(momentum flow)

(for LHM media)

$$\begin{cases} \vec{E} = \vec{E}_0 \cos(\omega t - \vec{k} \cdot \vec{x}) \\ \vec{B} = \vec{B}_0 \cos(\omega t - \vec{k} \cdot \vec{x}) \end{cases}$$

\sim monochromatic plane wave

$$\langle u \rangle = \frac{1}{2} \epsilon E_0^2$$

$$\langle \vec{S} \rangle = \frac{1}{2} \sqrt{\frac{\epsilon}{\mu}} E_0^2 \hat{k} = \langle u \rangle \vec{v} \sim \text{energy flow}$$

where $\vec{v} = \frac{\hat{k}}{\sqrt{\mu \epsilon}} \sim$ phase velocity

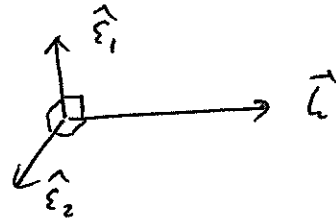
$\langle \dots \rangle \sim$ time averaged

Polarization

$$\vec{E} = \text{Re} \left\{ \vec{E}_0 e^{-i\omega t + i\vec{k} \cdot \vec{x}} \right\} \sim \text{plane wave}$$

choose a basis:

$$\vec{E}_0 = E_1 \hat{\epsilon}_1 + E_2 \hat{\epsilon}_2$$



E_1 & E_2 are complex in general (may contain phase shifts)

$\hat{\epsilon}_1, \hat{\epsilon}_2 \sim$ linear polarizations

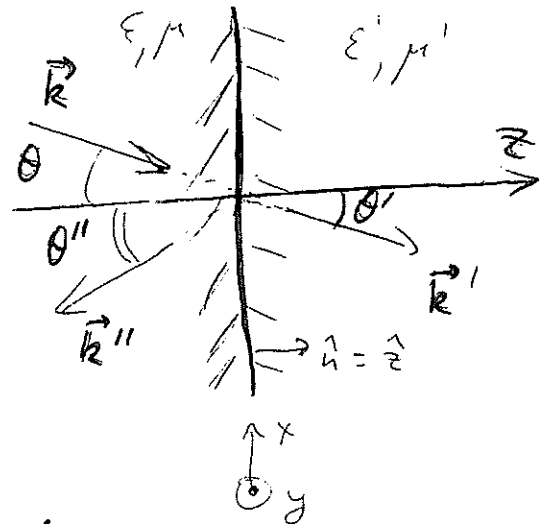
(Def.) $\hat{\epsilon}_{\pm} = \frac{1}{\sqrt{2}} (\hat{\epsilon}_1 \pm i\hat{\epsilon}_2)$ \sim circular polarizations
(a different basis)

$$\Rightarrow \vec{E}_0 = E_+ \hat{\epsilon}_+ + E_- \hat{\epsilon}_-$$

Reflection and Refraction

Incident wave:

$$\begin{cases} \vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{x} - \omega t)} \\ \vec{B} = \frac{1}{\omega} \vec{k} \times \vec{E} = \sqrt{\mu \epsilon} \frac{\vec{k} \times \vec{E}}{k} \end{cases}$$



Refracted wave:

$$\begin{cases} \vec{E}' = \vec{E}_0' e^{i(\vec{k}' \cdot \vec{x} - \omega' t)} \\ \vec{B}' = \sqrt{\mu' \epsilon'} \frac{\vec{k}' \times \vec{E}'}{k'} \end{cases}$$

(do not assume that all \vec{k} 's are in one plane)

Reflected wave:

$$\begin{cases} \vec{E}'' = \vec{E}_0'' e^{i(\vec{k}'' \cdot \vec{x} - \omega'' t)} \\ \vec{B}'' = \sqrt{\mu \epsilon} \frac{\vec{k}'' \times \vec{E}''}{k''} \end{cases}$$

Match boundary conditions. $\vec{\nabla} \cdot \vec{D} = 0 \Rightarrow D_n$ is cont.

$\vec{\nabla} \cdot \vec{B} = 0 \Rightarrow B_n$ is continuous

$\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0 \Rightarrow \vec{\nabla} \times \vec{E} - i\omega \vec{B} = 0 \Rightarrow E_t$ is continuous

as \vec{B} has no δ -function singularity at $z=0$

$\vec{\nabla} \times \vec{H} = \frac{\partial \vec{D}}{\partial t} = -i\omega \vec{D} \Rightarrow H_t$ is continuous (same reason) at $z=0$

to have any boundary conditions need

$$\omega = \omega' = \omega'' \Rightarrow k = k'' = \sqrt{\mu \epsilon} \omega, \quad k' = \sqrt{\mu' \epsilon'} \omega$$

spatial phase factors should also be equal

at $z=0$: $\vec{k} \cdot \vec{x} \Big|_{z=0} = \vec{k}' \cdot \vec{x} \Big|_{z=0} = \vec{k}'' \cdot \vec{x} \Big|_{z=0}, \quad \forall x, y$

\Rightarrow choose $\vec{h} = (h_x, 0, h_z) \Rightarrow \vec{h} \cdot \vec{x} \Big|_{z=0} = h_x \cdot x \Rightarrow$ no y -dep (245)

\Rightarrow there should be no y -dependence in $\vec{k}' \cdot \vec{x}$ and in $\vec{k}'' \cdot \vec{x}$ as well $\Rightarrow k'_y = k''_y = 0 \Rightarrow$ all lie in the same plane

$$k \cdot \sin \theta = k' \cdot \sin \theta' = k'' \cdot \sin \theta''$$

\Rightarrow as $k = k'' \Rightarrow \theta = \theta''$ \approx angle of reflection is equal to angle of incidence!

$$\text{as } k = \sqrt{\mu \epsilon} \omega \text{ and } k' = \sqrt{\mu' \epsilon'} \omega \Rightarrow$$

$$\sqrt{\mu \epsilon} \sin \theta = \sqrt{\mu' \epsilon'} \sin \theta'. \text{ Remember } n = c \sqrt{\mu \epsilon}$$

(index of refraction) $\Rightarrow n \sin \theta = n' \sin \theta'$

Snell's law!

The only thing left is to find \vec{E}_0', \vec{E}_0'' using b.c.'s:

$$D_n^{\perp} \text{ continuous} \Rightarrow \hat{n} \cdot [\epsilon (\vec{E}_0 + \vec{E}_0'') - \epsilon' \vec{E}_0'] = 0$$

$$B_n \text{ continuous} \Rightarrow \hat{n} \cdot [\vec{k} \times \vec{E}_0 + \vec{k}'' \times \vec{E}_0'' - \vec{k}' \times \vec{E}_0'] = 0$$

(and $\omega = \omega' = \omega''$)

$$E_t \text{ continuous} \Rightarrow \hat{n} \times [\vec{E}_0 + \vec{E}_0'' - \vec{E}_0'] = 0$$

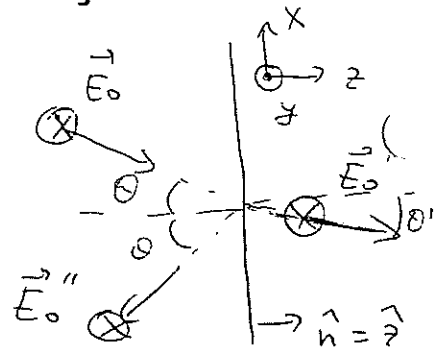
$$H_t \text{ continuous} \Rightarrow \left[\frac{1}{\mu} (\vec{k} \times \vec{E}_0 + \vec{k}'' \times \vec{E}_0'') - \frac{1}{\mu'} (\vec{k}' \times \vec{E}_0') \right] \times \hat{n} = 0$$

$\vec{E}_0', \vec{E}_0'' = 6$ unknowns, while there are 6 equations above

Consider 2 cases: (linear polarization)

(I) $\vec{E}_0 \perp$ to the plane of incidence.

$$\vec{E}_0, \vec{E}_0', \vec{E}_0'' \parallel \hat{y}$$



3rd & 4th equations. ($\hat{n} = \hat{z}$)

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$$\begin{cases} E_0 + E_0'' - E_0' = 0 \\ \frac{1}{\mu} (k E_0 \cos \theta - k'' E_0'' \cos \theta'') - \frac{1}{\mu'} k' E_0' \cos \theta' = 0 \end{cases}$$

as $k = k'' = \sqrt{\mu \epsilon} \omega$, $k' = \sqrt{\mu' \epsilon'} \omega \Rightarrow$ and $\theta = \theta''$

$$\begin{cases} E_0 + E_0'' - E_0' = 0 \\ \sqrt{\frac{\epsilon}{\mu}} (E_0 - E_0'') \cos \theta - \sqrt{\frac{\epsilon'}{\mu'}} E_0' \cos \theta' = 0 \end{cases}$$

1st eqn. $0 = 0$; 2nd eqn.: $k E_0 \sin \theta + k'' E_0'' \sin \theta - k' E_0' \sin \theta' = 0 \Rightarrow (E_0 + E_0'') \sin \theta - \sqrt{\frac{\mu' \epsilon'}{\mu \epsilon}} E_0' \sin \theta' = 0$

as $\sqrt{\mu \epsilon} \sin \theta = \sqrt{\mu' \epsilon'} \sin \theta'$ (Snell's law) $\Rightarrow E_0 + E_0'' - E_0' = 0$
 \Rightarrow duplicates the 3rd one.

Using Snell's law ($n \sin \theta = n' \sin \theta'$) to get rid of n' we write (work it out yourself):

$$\boxed{\begin{aligned} \frac{E_0'}{E_0} &= \frac{2n \cos \theta}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}} \\ \frac{E_0''}{E_0} &= \frac{n \cos \theta - \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}}{n \cos \theta + \frac{\mu}{\mu'} \sqrt{n'^2 - n^2 \sin^2 \theta}} \end{aligned}}$$

$$n = \sqrt{\frac{\mu \epsilon}{\mu_0 \epsilon_0}}$$

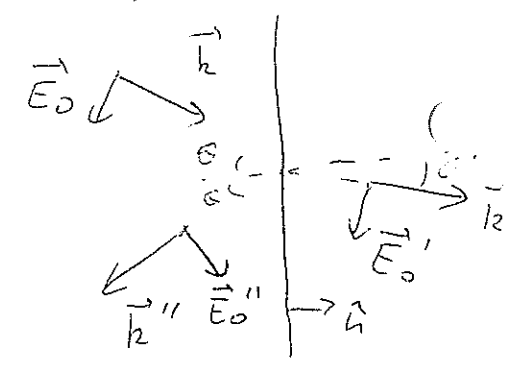
$$n' = \sqrt{\frac{\mu' \epsilon'}{\mu_0 \epsilon_0}}$$

Fresnel equations

Ⓜ $\vec{E}_0 \parallel$ plane of incidence (xz plane)

2 independent equations (3rd & 4th):

$$\begin{cases} (E_0 + E_0'') \cos \theta - E_0' \cos \theta' = 0 \\ \sqrt{\frac{\epsilon}{\mu}} (E_0 - E_0'') - \sqrt{\frac{\epsilon'}{\mu'}} E_0' = 0 \end{cases}$$



(other two can be reduced to those)

Solve:

(Using Snell's Law)

$$\frac{E_0'}{E_0} = \frac{2n n' \cos \theta}{\frac{n}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

$$\frac{E_0''}{E_0} = \frac{-\frac{n}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}{\frac{n}{\mu'} n'^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$$

Fresnel equations

Normal incidence: $\theta = 0 \Rightarrow$ both Ⓚ and Ⓜ

give

$$\frac{E_0'}{E_0} = \frac{2n}{n + \frac{n}{\mu'} n'}$$

$$\frac{E_0''}{E_0} = \frac{n - \frac{n}{\mu'} n'}{n + \frac{n}{\mu'} n'}$$

if $\mu = \mu'$
 $n' > n$
 \Rightarrow reflected wave change sign - phase revers

Polarization by reflection: put $\mu = \mu'$ for simplicity.

Ⓚ: $\frac{E_0''}{E_0} = \frac{n \cos \theta - \sqrt{n'^2 - n^2 \sin^2 \theta}}{n \cos \theta + \sqrt{n'^2 - n^2 \sin^2 \theta}}$

\leftarrow different \Rightarrow

Ⓜ: $\frac{E_0''}{E_0} = \frac{-n^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}{n^2 \cos \theta + n \sqrt{n'^2 - n^2 \sin^2 \theta}}$

$\leftarrow \Rightarrow$ reflected light is polarized!