

Relativistic Particles in Electromagnetic Fields (23)

The Lagrangian and the action

Let us construct the Lagrangian for the interaction of electromagnetic field with a point particle of charge e and mass m .

The Lagrangian for the particle is

$$L_{\text{free}} = -mc^2 \sqrt{1 - \frac{v^2}{c^2}}$$

Given the action for a system

$$S = \int_{t_1}^{t_2} dt L(q_i(t), \dot{q}_i(t), t)$$

the least action principle gives us equations of motion ^(EOM) (Euler-Lagrange equations)

$$\boxed{\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0}$$

$q_i \sim$ generalized coordinates, $\dot{q}_i = \frac{dq_i}{dt} \sim$ velocities

Apply this to L_{free} : $q_i(t) = x^i(t)$,

$$\dot{q}_i(t) = \frac{d}{dt} x^i(t) = \dot{x}^i(t) = v^i(t) \Rightarrow L_{\text{free}} = L_{\text{free}}(\dot{x}^i)$$

$$\Rightarrow \frac{\partial L_{\text{free}}}{\partial x^i} = 0, \quad \frac{\partial L}{\partial \dot{x}^i} = -mc^2 \frac{-2v^i/c^2}{2\sqrt{1 - \frac{v^2}{c^2}}} = m \gamma v^i \quad \begin{array}{l} \text{momentum} \\ \text{conservation} \\ \downarrow \end{array}$$

$$\Rightarrow \text{EOM are } \frac{d}{dt} (m \gamma \vec{v}) = 0 \Rightarrow \frac{d\vec{p}}{dt} = 0 \Rightarrow \vec{p} = \text{const}$$

Electromagnetic field: imagine a 4-vector field $A_\mu(x)$, where x denotes a point in 4-dim space-time x^μ .

$$A^\mu(x) = (A^0(x), \vec{A}(x))$$

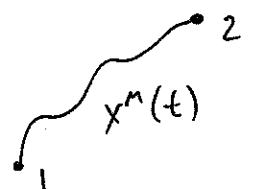
(Def.) $A^0(x) \equiv \Phi(x)$ is the scalar potential

(Def.) $\vec{A}(x) = (A^1(x), A^2(x), A^3(x))$ is the vector potential.

We need to construct a term in the action (Lagrangian) which combines both the field $A^\mu(x)$ and the motion of the particle, described by $x^\mu(t) \sim$ a trajectory.

Physics is Lorentz-invariant \Rightarrow so is the action. To get a L-inv. object multiply the field A^μ by the infinitesimal element of the particle's world line (dx^μ):

$$S_{\text{int}} \sim \int_1^2 A^\mu dx^\mu$$



The integral goes along the particle's trajectory. (Took S to be linear in $A^\mu \sim$ experimental result. (Sds $A_\mu A^\mu$ allowed).

It turns out that the proportionality coefficient is $-\frac{e}{c}$, such that

$$\mathcal{S}_{int} = -\frac{e}{c} \int_1^2 dx^\mu A^\mu$$

c = speed of light, e = particle's charge
(naturally neutral particles do not feel EM field \Rightarrow if $e=0$ need $\mathcal{S}_{int}=0$)

$$\mathcal{S}_{int} = -\frac{e}{c} \int_{t_1}^{t_2} dt \frac{dx^\mu}{dt} A^\mu = -\frac{e}{c} \int_{t_1}^{t_2} dt \frac{1}{\gamma} u_\mu A^\mu$$

as $dt = \gamma d\tau$

where $u^\mu = \frac{dx^\mu}{d\tau}$ is the 4-velocity.

Since $\mathcal{S} = \int dt L \Rightarrow$ $L_{int} = -\frac{e}{c} \frac{1}{\gamma} u_\mu A^\mu$

The total Lagrangian for particle in the EM field is

$$L = -mc^2 \sqrt{1 - \frac{v^2}{c^2}} - \frac{e}{c\gamma} u_\mu A^\mu \Rightarrow \text{as } u^\mu = \gamma(c, \vec{v})$$

$$L = -mc^2 \sqrt{1 - \frac{v^2}{c^2}} - e\Phi + \frac{e}{c} \vec{v} \cdot \vec{A}$$

familiar from electrostatics term

\Rightarrow could be used to fix the prefactor in \mathcal{S}_{int}

Equations of motion: $\frac{d}{dt} \left(\frac{\partial L}{\partial v^i} \right) - \frac{\partial L}{\partial x^i} = 0$

$\Rightarrow \frac{\partial L}{\partial v^i} = m \gamma v^i + \frac{e}{c} A^i = p^i \sim \text{canonical momentum}$

$$\vec{P} = \vec{p} + \frac{e}{c} \vec{A}$$

$$\frac{\partial L}{\partial x^i} = -e \frac{\partial \Phi}{\partial x^i} + \frac{e}{c} v^j \frac{\partial A^j}{\partial x^i}$$

$\Rightarrow \frac{d}{dt} \left(\underbrace{m \gamma v^i + \frac{e}{c} A^i}_{p^i} \right) + e \frac{\partial \Phi}{\partial x^i} - \frac{e}{c} v^j \frac{\partial A^j}{\partial x^i} = 0$

$$\frac{dp^i}{dt} = -e \frac{\partial \Phi}{\partial x^i} - \frac{e}{c} \left[\frac{dA^i}{dt} - v^j \frac{\partial A^j}{\partial x^i} \right]$$

As $\frac{dA^i(x)}{dt} = \frac{\partial A^i(x)}{\partial t} + \frac{\partial x^j}{\partial t} \frac{\partial A^i(x)}{\partial x^j} = \frac{\partial A^i}{\partial t} + v^j \frac{\partial A^i}{\partial x^j}$

We get

$$\frac{dp^i}{dt} = -\frac{e}{c} \left[c \left(\overset{\Phi}{\frac{\partial A^0}{\partial x^i} + \frac{1}{c} \frac{\partial A^i}{\partial t}} \right) + v^j \frac{\partial A^i}{\partial x^j} - v^j \frac{\partial A^j}{\partial x^i} \right]$$

Note that $\frac{\partial}{\partial x^i} = -\frac{\partial}{\partial x_i}$ and $x^0 = ct \Rightarrow$

$$\frac{dp^i}{dt} = -\frac{e}{c} \left[c \left(-\frac{\partial A^0}{\partial x_i} + \frac{\partial A^i}{\partial x^0} \right) - v^j \left(-\frac{\partial A^j}{\partial x_i} + \frac{\partial A^i}{\partial x_j} \right) \right]$$

Euler-Lagrange Equations

(23)

$$S = \int_{t_1}^{t_2} dt L(q, \dot{q})$$

\Rightarrow least action principle requires that $\delta S = 0$

under $q \rightarrow q + \delta q$. We get

$$0 = \delta S = \int_{t_1}^{t_2} dt \left[\frac{\delta L}{\delta q} \delta q + \frac{\delta L}{\delta \dot{q}} \delta \dot{q} \right] = \frac{\delta L}{\delta \dot{q}} \delta q \Big|_{t_1}^{t_2} +$$

$\underbrace{\hspace{10em}}_{\frac{d}{dt} \delta q \Rightarrow \text{parts}}$

$$+ \int_{t_1}^{t_2} dt \left[\frac{\delta L}{\delta q} - \frac{d}{dt} \left(\frac{\delta L}{\delta \dot{q}} \right) \right] \delta q$$

\Rightarrow require that $\delta q(t_1) = \delta q(t_2) = 0 \Rightarrow$ drop the first term

\Rightarrow since the above result is true for any $\delta q(t)$

$$\Rightarrow \boxed{\frac{\delta L}{\delta q} - \frac{d}{dt} \left(\frac{\delta L}{\delta \dot{q}} \right) = 0}$$

Euler-Lagrange equation
aka
equation of motion (EOM)

