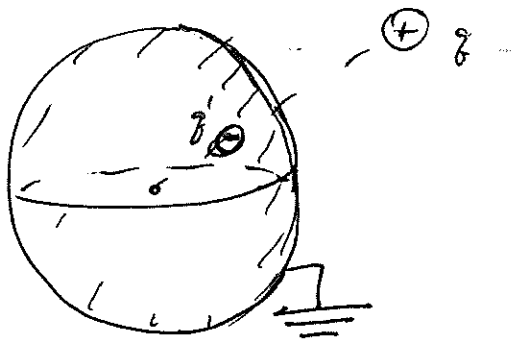


Last time | We finished constructing Dirichlet Green function for a sphere (outside a sphere)

$$G_D(\vec{x}, \vec{x}') = \frac{1}{|\vec{x} - \vec{x}'|} - \frac{R}{x'} \frac{1}{\left| \vec{x} - \frac{R^2}{x'^2} \vec{x}' \right|}$$

using the method of images:



grounded conducting sphere

Employing the solution of this problem we found the force on charge  $q$ :

$$F = \frac{-1}{4\pi\epsilon_0} \frac{q^2 R/y}{\left[ y - \frac{R^2}{y} \right]^2} \quad (\text{attractive})$$

and the induced charge density on the sphere

$$\sigma = -\epsilon_0 \left. \frac{\partial \Phi}{\partial r} \right|_{r=R} = \frac{-q}{4\pi R} \frac{y^2 - R^2}{\left[ y^2 + R^2 - 2Ry \cos\theta \right]^{3/2}}$$

where  $\theta$  is the angle between  $\vec{y}$  &  $\vec{x}$ .

$\Rightarrow$  Note: the sphere is at  $\Phi = 0 \Rightarrow$  if we now disconnect the ground, nothing would change

$\Rightarrow$  can bring in a charge  $Q - q'$  and place it on the sphere  $\Rightarrow$  would distribute itself uniformly on the surface

$$\Rightarrow \Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \left[ \frac{q}{|\vec{x} - \vec{y}|} - q \frac{R}{s} \frac{1}{|\vec{x} - \frac{R^2}{s^2} \vec{y}|} + \frac{Q + q \frac{R}{s}}{|\vec{x}|} \right]$$

Suppose the sphere has <sup>a total</sup> charge  $Q \Rightarrow$

split it into  $q'$  &  $Q - q'$  [alternatively; start with a grounded sphere and charge  $q$ ; disconnect the ground; you now have  $q'$  on the sphere and  $\Phi = 0$  then add  $Q - q'$ .

$\Rightarrow$  In conductors all charge sits on the surface  
bring  $Q$  from infinity: first bring  $q'$  then  $Q - q'$ .

$\Rightarrow q'$  along with  $q$  creates  $\Phi = 0$  on surface

$\Rightarrow Q - q'$  is uniformly distributed on the surface, giving extra

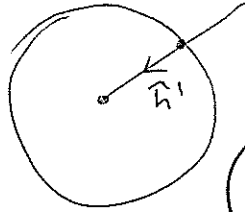
$$\Delta \Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \frac{Q - q'}{|\vec{x}|} = \frac{1}{4\pi\epsilon_0} \frac{Q + q R/y}{|\vec{x}|}$$

in addition to the potential found above, such that total potential is

$$\Phi(\vec{x}) = \left[ \frac{q}{|\vec{x} - \vec{y}|} - q \frac{R}{y} \frac{1}{|\vec{x} - \frac{R^2}{y^2} \vec{y}|} + \frac{Q + q R/y}{|\vec{x}|} \right] \frac{1}{4\pi\epsilon_0}$$

For general boundary condition  $\Phi(R, \theta, \varphi)$  on the surface of the sphere, <sup>to find</sup> the potential

we need  $\left. \frac{\partial G_D(\vec{x}, \vec{x}')}{\partial n'} \right|_{|\vec{x}'|=R} = - \left. \frac{\partial G_D}{\partial r'} \right|_{r'=R} \Rightarrow$



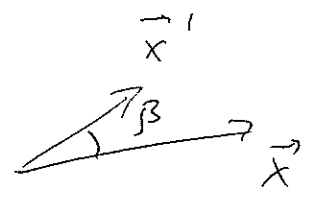
General formula: (Dirichlet Green function)

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int_V d^3x' G_D(\vec{x}, \vec{x}') \rho(\vec{x}') - \frac{1}{4\pi} \oint_S d\Omega' \Phi(\vec{x}') \frac{\partial G_D(\vec{x}, \vec{x}')}{\partial n'}$$

Write

$$G_D(\vec{x}, \vec{x}') = \frac{1}{\sqrt{x^2 + x'^2 - 2xx' \cos \beta}}$$

$$- \frac{R}{x'} \frac{1}{\sqrt{x^2 + \frac{R^4}{x'^2} - 2x \frac{R^2}{x'} \cos \beta}}$$



$$\Rightarrow - \left. \frac{\partial G_D}{\partial x'} \right|_{x'=R} = \left\{ \frac{x' - x \cos \beta}{\left( \sqrt{x^2 + x'^2 - 2xx' \cos \beta} \right)^3} - R \right.$$

$$\left. \cdot \frac{x' x^2 - x R^2 \cos \beta}{\left[ x^2 x'^2 + R^4 - 2xx'R^2 \cos \beta \right]^{3/2}} \right\} \Big|_{x'=R} =$$

$$= \frac{R - x^2/R}{\left( x^2 + R^2 - 2xR \cos \beta \right)^{3/2}}$$

outside charges

$$\Rightarrow \Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' \rho(\vec{x}') G_D(\vec{x}, \vec{x}') - \frac{1}{4\pi} \int d\Omega' \frac{R(R^2 - x^2)}{\left( x^2 + R^2 - 2xR \cos \beta' \right)^{3/2}} \Phi(R, \theta', \varphi')$$

where  $d\Omega = d\cos \theta d\varphi$

and  $\beta =$  angle between  $\vec{x}$  and  $\vec{x}'$

If  $\Phi(R, \theta', \varphi') = V$  constant &  $\rho = 0$

$$\Rightarrow \Phi(\vec{x}) = \frac{1}{4\pi} \int d\Omega V \frac{R(x^2 - R^2)}{(x^2 + R^2 - 2xR \cos\beta)^{3/2}} =$$

$$= \frac{1}{4\pi} \cdot 2\pi \cdot \int_{-1}^1 d\cos\beta \frac{R(x^2 - R^2)}{(x^2 + R^2 - 2xR \cos\beta)^{3/2}} V =$$

$$= \frac{1}{2} \frac{R(x^2 - R^2)}{xR} V \left[ \frac{1}{(x^2 + R^2 - 2xR \cos\beta)^{1/2}} \right]_{-1}^1 =$$

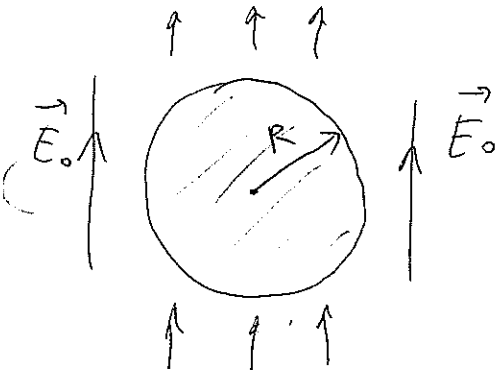
$$= \frac{1}{2} \frac{x^2 - R^2}{x} V \left( \frac{1}{|x-R|} - \frac{1}{|x+R|} \right) = \text{as } x > R =$$

$$= V \frac{R}{x} \Rightarrow \Phi(\vec{x}) = V \frac{R}{x}$$

as one'd expect from Gauss's law.

(skip)

Consider a conducting sphere in a uniform electric field  $\vec{E}_0$ : let's guess the answer for  $\phi$ !



$$\phi(\vec{r}) = \underbrace{-\vec{E}_0 \cdot \vec{r}}_{\text{potential due to field } \vec{E}_0} + \underbrace{\phi_{\text{sphere}}(\vec{r})}_{\text{potential due to the sphere.}}$$

Laplace equation  $\nabla^2 \phi = 0$  is valid everywhere outside of the sphere  $\Rightarrow$

$$\Rightarrow \nabla^2 \phi_{\text{sphere}}(\vec{r}) = 0.$$

( & vanishes at  $\infty$  )

$\phi_{\text{sphere}}(\vec{r})$  depends on  $r$  and  $\vec{E}_0$  & satisfies

$$\nabla^2 \phi_{\text{sphere}} = 0 \Rightarrow \phi_{\text{sphere}}(\vec{r}) \propto \vec{E}_0 \cdot \vec{\nabla} \frac{1}{r}$$

is a natural guess  $\Rightarrow$  as  $\vec{\nabla} \frac{1}{r} = -\frac{\vec{r}}{r^3} \Rightarrow$

$$\phi(\vec{r}) = -\vec{E}_0 \cdot \vec{r} + C \cdot \vec{E}_0 \cdot \frac{\vec{r}}{r^3}$$

$\uparrow$  constant

Require that  $\phi(\vec{r})|_{|\vec{r}|=R} = 0 \Rightarrow -1 + \frac{C}{R^3} = 0$

$$\Rightarrow C = R^3 \Rightarrow$$

$$\phi(\vec{r}) = -\vec{E}_0 \cdot \vec{r} \left( 1 - \frac{R^3}{r^3} \right)$$

(cf. Jackson's Eq. (2.14)).

Surface charge density

$$\sigma = -\epsilon_0 \left. \frac{\partial \phi}{\partial r} \right|_{r=a} = \epsilon_0 E_0 \cos \theta + 2\epsilon_0 E_0 \cos \theta = 3\epsilon_0 E_0 \cos \theta$$

$$\Rightarrow Q = \oint_{\vec{r}} da \sigma = \int_0^{2\pi} d\varphi \int_{-1}^1 d\cos \theta \ 3\epsilon_0 E_0 \cos \theta = 0$$

$\Rightarrow$  sphere could be insulated or grounded...