

(A1)

$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) G(\vec{r}, t) = -4\pi \delta^3(\vec{r}) \delta(t)$$

write $G(\vec{r}, t) = \int \frac{d^3k}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{i\vec{k}\cdot\vec{r} - i\omega t} \tilde{G}(\vec{k}, \omega)$

as $\delta^3(\vec{r}) \delta(t) = \int \frac{d^3k}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{i\vec{k}\cdot\vec{r} - i\omega t}$

$$\Rightarrow \left(-\vec{k}^2 + \frac{\omega^2}{c^2} \right) \tilde{G} = -4\pi$$

$$\Rightarrow \tilde{G} = \frac{-4\pi}{\frac{\omega^2}{c^2} - \vec{k}^2} \quad \text{photon propagator}$$

$$\Rightarrow G(\vec{r}, t) = \int \frac{d^3k}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{i\vec{k}\cdot\vec{r} - i\omega t} \frac{-4\pi}{\frac{\omega^2}{c^2} - \vec{k}^2}$$

Does this expression make sense? Essential sing.

at $\frac{\omega}{c} = |\vec{k}|$. There are several ways to regulate

it:

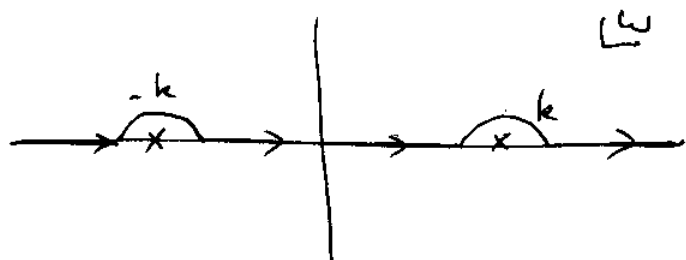
A. Retarded (causal) Green function

demand that $G(\vec{r}, t) = 0$ for $t < 0$

$$G(\vec{r}, t) = \int \frac{d^3k}{(2\pi)^3} \int \frac{d\omega}{2\pi} e^{-i\omega t + i\vec{k}\cdot\vec{r}} \frac{-4\pi}{\left(\frac{\omega}{c} - k\right)\left(\frac{\omega}{c} + k\right)} \quad \text{with } k = |\vec{k}|$$

Need $G = 0$ for $t < 0$: if $t < 0$ have to close the ω -contour into the upper half-plane. (A2)

\Rightarrow need to have poles in the lower half-plane:



$$\Rightarrow G(\vec{r}, t) = -4\pi \int \frac{d^3k}{(2\pi)^3} \int \frac{d\omega}{2\pi} e^{-i\omega t + i\vec{k} \cdot \vec{r}} \frac{1}{\left(\frac{\omega}{c} - k + i\epsilon\right)\left(\frac{\omega}{c} + k + i\epsilon\right)}$$

$$\Rightarrow G_{ret}(\vec{r}, t) = -4\pi \int \frac{d^3k}{(2\pi)^3} \frac{d\omega}{2\pi} e^{-i\omega t + i\vec{k} \cdot \vec{r}} \frac{1}{\frac{\omega^2}{c^2} - k^2 + i\epsilon\omega}$$

advanced \sim change signs of $i\epsilon$'s.

Do the Fourier transform

$$\begin{aligned} G_{ret}(\vec{r}, t) &= -4\pi \Theta(t) (-2\pi i) \frac{c^2}{2\pi} \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot \vec{r}} \left[\frac{e^{-ikct}}{2k\epsilon} - \frac{e^{ikct}}{2k\epsilon} \right] \\ &= 2\pi i \Theta(t) c \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot \vec{r}} \frac{1}{k} [e^{-ikct} - e^{ikct}] \\ &= \frac{|\vec{r}|}{(2\pi)^3} = 2\pi i c \Theta(t) \int_0^\infty dk \cdot k^2 \int_{-1}^1 d\cos\theta \int_0^{2\pi} d\varphi e^{ikr\cos\theta} \frac{1}{k} \\ [e^{-ikct} - e^{ikct}] &= \frac{ic}{2\pi} \Theta(t) \int_0^\infty dk \cdot k \cdot \frac{1}{k^2 r} [e^{ikr} - e^{-ikr}] \end{aligned}$$

$$\cdot \left[e^{-ikct} - e^{ikct} \right] = \frac{c}{2\pi r} \theta(t) \int_0^\infty dk \left[e^{ik(r-ct)} + e^{-ik(r-ct)} - e^{-ik(r+ct)} - e^{ik(r+ct)} \right] \cdot e^{-\delta k}$$

(δ is some regulator at $k \rightarrow +\infty$)

↑ this is the same formula as obtained in class, now we'll evaluate it differently

$$= \frac{c}{2\pi r} \theta(t) \cdot \left\{ \frac{-1}{i(r-ct)-\delta} + \frac{-1}{-i(r-ct)-\delta} - \frac{-1}{-i(r+ct)-\delta} - \frac{-1}{+i(r+ct)-\delta} \right\}$$

$$= \frac{ci}{2\pi r} \theta(t) \left\{ \frac{1}{r-ct+i\delta} - \frac{1}{r-ct-i\delta} + \frac{1}{r+ct-i\delta} - \frac{1}{r+ct+i\delta} \right\}$$

$$= \frac{ci}{2\pi r} \theta(t) \left\{ -2\pi i \delta(r-ct) + 2\pi i \delta(r+ct) \right\}$$

" as $r > 0, t > 0$

using $\frac{1}{x-i\delta} - \frac{1}{x+i\delta} = 2\pi i \delta(x)$
 ↑ Dirac delta-fun.
 ↑ regulator
 ! do not confuse!

$$= \frac{c}{r} \theta(t) \delta(r-ct) = \frac{1}{r} \delta\left(t - \frac{r}{c}\right) \text{ as desired!}$$

$$\Rightarrow \boxed{G_{ret}(\vec{r}, t) = \frac{1}{r} \delta\left(t - \frac{r}{c}\right)}$$

B. Advanced Green function (can be evaluated in a similar way)

$$G_{adv}(\vec{r}, t) = -4\pi \int \frac{d^3k}{(2\pi)^3} \frac{d\omega}{2\pi} e^{-i\omega t + i\vec{k} \cdot \vec{r}} \frac{1}{\frac{\omega^2}{c^2} - k^2 - i\omega\epsilon}$$