

We know $L \Rightarrow$ can find canonical momentum

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$$\vec{P} \text{ by } P_i = \frac{\partial L}{\partial v_i} = \underbrace{\gamma m v_i}_{\text{momentum of a free particle}} + \frac{e}{c} A_i = \left(\vec{p} + \frac{e}{c} \vec{A} \right)_i$$

$$\Rightarrow \vec{P} = \vec{p} + \frac{e}{c} \vec{A} \quad \left(\text{in QM } \vec{P} \rightarrow -i\hbar \vec{\nabla} \right)$$

quantization

Now we can find the Hamiltonian of the system:

$$H = \vec{P} \cdot \vec{v} - L \Rightarrow \left. \begin{array}{l} \text{with } \vec{v} \text{ expressed} \\ \text{in terms of } \vec{P} \end{array} \right\} \begin{array}{l} \vec{P} = \frac{m\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}} + \frac{e}{c} \vec{A} \\ \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2 = \frac{m^2 v^2}{1-\frac{v^2}{c^2}} \end{array}$$

$$\Rightarrow v^2 \left(m^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2 \frac{1}{c^2} \right) = \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2$$

$$\Rightarrow \vec{v} = \frac{c \left(\vec{P} - \frac{e}{c} \vec{A} \right)}{\sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}}$$

$$\Rightarrow H = \frac{\vec{P} \cdot \left(c \vec{P} - e \vec{A} \right)}{\sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}} + e \Phi - \frac{e}{c} \frac{\vec{A} \cdot \left(c \vec{P} - e \vec{A} \right)}{\sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}}$$

$$+ m c^2 \sqrt{1 - \frac{\left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}} = \frac{c \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}{\sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}} + e \Phi +$$

$$+ \frac{m^2 c^3}{\sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2}} = c \sqrt{m^2 c^2 + \left(\vec{P} - \frac{e}{c} \vec{A} \right)^2} + e \Phi$$

$$\Rightarrow H = \sqrt{m^2 c^4 + (c\vec{P} - e\vec{A})^2} + e\Phi$$

Motion of a point charge in external \vec{E}, \vec{B} fields:

We have Lorentz force $\frac{d\vec{P}}{dt} = q(\vec{E} + \frac{1}{c}\vec{v} \times \vec{B})$
and energy change $\frac{dE}{dt} = q\vec{v} \cdot \vec{E}$.

A. Constant Electric Field.

$$\frac{d\vec{P}}{dt} = q\vec{E} \Rightarrow \vec{P} = q\vec{E}t + \text{const} \Rightarrow \text{if}$$

the particle starts from rest $\Rightarrow \vec{P}|_{t=0} = 0$

$$\Rightarrow \vec{P} = q\vec{E}t \Rightarrow \text{is } \vec{E} = E\hat{x} \Rightarrow$$

$$\Rightarrow p_x = qEt, \quad p_y = p_z = 0$$

$$\Rightarrow \frac{mV}{\sqrt{1 - \frac{v^2}{c^2}}} = qEt \Rightarrow m^2 \left(\frac{dx}{dt}\right)^2 = q^2 E^2 t^2 \left(1 - \frac{1}{c^2} \left(\frac{dx}{dt}\right)^2\right)$$

$$\Rightarrow \frac{dx}{dt} = \frac{qEt}{\sqrt{m^2 + \frac{q^2}{c^2} E^2 t^2}}$$

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$$\Rightarrow x(t) = \int_0^t dt' \frac{q E t'}{\sqrt{m^2 + \frac{q^2}{c^2} E^2 t'^2}} = \frac{q E}{m} \frac{c^2 m^2}{q^2 E^2} \left(\sqrt{1 + \frac{q^2 E^2}{m^2 c^2} t^2} - 1 \right)$$

assume $x(0) = 0$

$$\Rightarrow x(t) = \frac{m c^2}{q E} \left(\sqrt{1 + \frac{q^2 E^2}{m^2 c^2} t^2} - 1 \right)$$

moves with
speed of light!

\Rightarrow as $t \rightarrow \infty \Rightarrow x(t) \approx c t$ ~ linear in t !

\Rightarrow if c is large \Rightarrow expand in powers of $\frac{1}{c} \Rightarrow$

$\Rightarrow x(t) \approx \frac{1}{2} \frac{q E}{m} t^2 = \frac{1}{2} a t^2$ ~ well-known classical NR result!

B. Constant Uniform Magnetic Field.

$$\frac{d\vec{p}}{dt} = \frac{q}{c} \vec{v} \times \vec{B}, \quad \frac{dE}{dt} = 0 \Rightarrow E = \text{const.}$$

$$\Rightarrow \text{write } \vec{p} = m \gamma \vec{v} = m \gamma c^2 \cdot \frac{\vec{v}}{c^2} = E \cdot \frac{\vec{v}}{c^2}$$

$$\Rightarrow \frac{E}{c^2} \frac{d\vec{v}}{dt} = \frac{q}{c} \vec{v} \times \vec{B} \Rightarrow \text{define } \vec{\omega}_B = \frac{q \vec{B} c}{E} = \frac{q \vec{B}}{\gamma m c}$$

(precession frequency)

$$\Rightarrow \frac{d\vec{v}}{dt} = \vec{v} \times \vec{\omega}_B \Rightarrow \text{if } \vec{B} = B \hat{z} \Rightarrow \vec{\omega}_B = \omega_B \hat{z}$$

$$\Rightarrow \text{get } \dot{V}_x = \omega_B V_y, \quad \dot{V}_y = -\omega_B V_x, \quad \dot{V}_z = 0$$

$$\Rightarrow \ddot{V}_x = \omega_B \dot{V}_y = -\omega_B^2 V_x \Rightarrow V_x = V_{0\perp} \cdot e^{\pm i\omega_B t}$$

$$\Rightarrow V_y = \frac{1}{\omega_B} \dot{V}_x = \pm i V_{0\perp} e^{\pm i\omega_B t}$$

\Rightarrow taking real parts write $V_x = V_{0\perp} \cos(\omega_B t + \alpha)$

$$\Rightarrow V_y = -V_{0\perp} \sin(\omega_B t + \alpha) \Rightarrow \sqrt{V_x^2 + V_y^2} = V_{0\perp}$$

\sim transverse (w.r. to \vec{B}) velocity

$$\Rightarrow \text{as } V_x = \dot{x} = V_{0\perp} \cos(\omega_B t + \alpha) \Rightarrow$$

$$\Rightarrow \begin{cases} x(t) = x_0 + r \sin(\omega_B t + \alpha) \\ y(t) = y_0 + r \cos(\omega_B t + \alpha) \\ z(t) = z_0 + V_{0z} t \end{cases}$$

$$r = \frac{V_{0\perp}}{\omega_B} = \frac{V_{0\perp} E}{q B c} = \frac{c p_{0\perp}}{q B}$$

Motion of a positive charge is shown here \rightarrow

