

Retarded time is defined by  $t_{ret} = t - \frac{R(t_{ret})}{c}$

$\Rightarrow t = t_{ret} + \frac{R(t_{ret})}{c} \Rightarrow$  as  $\frac{dt}{dt_{ret}} = [1 - \hat{n} \cdot \vec{\beta}]_{ret} \Rightarrow$

$$\vec{A}(\omega) = \left( \frac{e^2}{8\pi^2 c} \right)^{1/2} \int_{-\infty}^{\infty} dt' e^{i\omega(t' + \frac{R(t')}{c})} \frac{\hat{n} \times [(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{\beta} \cdot \hat{n})^2}$$

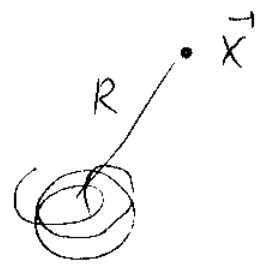
leading to

$$\frac{dI}{d\omega d\Omega} = \frac{e^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt' e^{i\omega(t' + \frac{R(t')}{c})} \frac{\hat{n} \times [(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{\beta} \cdot \hat{n})^2} \right|^2$$

General result for radiation  $\omega$ -spectrum & angular dependence.

$\Rightarrow$  assume that we are far from the charge, which is moving in a localized trajectory:

$R \gg |\vec{r}'| \Rightarrow R = |\vec{x} - \vec{r}(t)| \approx x - \hat{n} \cdot \vec{r}$



$\Rightarrow \vec{A}(\omega) = \left( \frac{e^2}{8\pi^2 c} \right)^{1/2} \int_{-\infty}^{\infty} dt' e^{i\omega(t' + \frac{x}{c} - \frac{\hat{n} \cdot \vec{r}(t')}{c})}$

$\cdot \frac{\hat{n} \times [(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{\beta} \cdot \hat{n})^2} \Rightarrow \frac{dI}{d\omega d\Omega} = \frac{e^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt' e^{i\omega(t' - \frac{\hat{n} \cdot \vec{r}(t')}{c})} \right|^2$

$$\cdot \frac{\hat{n} \times [(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{\beta} \cdot \hat{n})^2} \Big|^2$$

Now as 
$$\frac{\hat{n} \times [(\hat{n} - \vec{\beta}) \times \vec{\beta}]}{(1 - \vec{\beta} \cdot \hat{n})^2} = \frac{d}{dt} \left[ \frac{\hat{n} \times (\hat{n} \times \vec{\beta})}{(1 - \vec{\beta} \cdot \hat{n})} \right]$$

⇒ inserting and integrating by parts yields

$$\frac{dI}{d\omega d\Omega} = \frac{e^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt' e^{i\omega(t' - \frac{\hat{n} \cdot \vec{r}(t')}{c})} i\omega (1 - \hat{n} \cdot \vec{\beta}) \right.$$

$$\left. \frac{\hat{n} \times (\hat{n} \times \vec{\beta})}{1 - \hat{n} \cdot \vec{\beta}} \right|^2 = \frac{e^2 \omega^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt' e^{i\omega(t' - \frac{\hat{n} \cdot \vec{r}(t')}{c})} \hat{n} \times (\hat{n} \times \vec{\beta}) \right|^2$$

⇒ 
$$\frac{dI}{d\omega d\Omega} = \frac{e^2 \omega^2}{4\pi^2 c} \left| \int_{-\infty}^{\infty} dt' e^{i\omega(t' - \frac{\hat{n} \cdot \vec{r}(t')}{c})} \hat{n} \times (\hat{n} \times \vec{\beta}) \right|^2$$

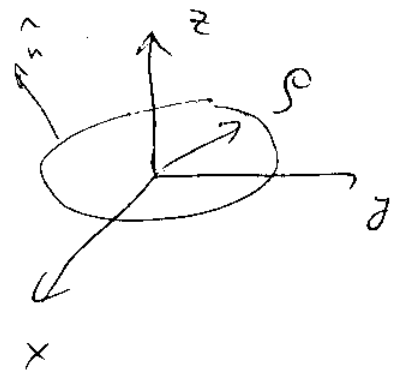
### Synchrotron Radiation

Imagine a particle in circular motion in

xy-plane: 
$$\vec{\beta} = \beta \left( \hat{x} \cos\left(\frac{vt}{\rho}\right) + \hat{y} \sin\left(\frac{vt}{\rho}\right) \right)$$

$\rho$  is the radius of the orbit,

$v$  is particle's velocity



⇒ 
$$\hat{n} \times (\hat{n} \times \vec{\beta}) = \hat{n} (\hat{n} \cdot \vec{\beta}) - \vec{\beta}$$

if 
$$\hat{n} = \sin\theta \cos\varphi \hat{x} + \sin\theta \sin\varphi \hat{y} + \cos\theta \hat{z} \Rightarrow$$

$$\Rightarrow \hat{n} \cdot \vec{\beta} = \beta \sin \theta \left( \cos \varphi \cos \frac{vt}{\rho} + \sin \varphi \sin \frac{vt}{\rho} \right) = \quad (77)$$

$$= \beta \sin \theta \cos \left( \varphi - \frac{vt}{\rho} \right)$$

$$\Rightarrow \hat{n} \times (\hat{n} \times \vec{\beta}) = \hat{x} \left[ \beta \sin^2 \theta \cos \varphi \cos \left( \varphi - \frac{vt}{\rho} \right) - \beta \cos \frac{vt}{\rho} \right]$$

$$+ \hat{y} \left[ \beta \sin^2 \theta \sin \varphi \cos \left( \varphi - \frac{vt}{\rho} \right) - \beta \sin \frac{vt}{\rho} \right] + \hat{z} \cos \theta \cdot \beta \sin \theta \cos \left( \varphi - \frac{vt}{\rho} \right)$$

~~$$+ \beta \hat{y} \sin \left( \frac{vt}{\rho} \right) + \beta (\hat{n} \times \hat{y}) \cos \left( \frac{vt}{\rho} \right) \sin \theta$$~~

$$= \sin \frac{vt}{\rho} \left[ \beta \hat{x} \sin^2 \theta \cos \varphi \sin \varphi + \beta \hat{y} (\sin^2 \theta \sin^2 \varphi - 1) \right]$$

$$+ \hat{z} \beta \sin \theta \cos \theta \cdot \sin \varphi + \cos \frac{vt}{\rho} \left[ \hat{x} (\beta \sin^2 \theta \cos^2 \varphi - \beta) \right]$$

$$+ \hat{y} \left[ \beta \sin^2 \theta \sin \varphi \cos \varphi + \hat{z} \beta \sin \theta \cos \varphi \cos \theta \right]$$

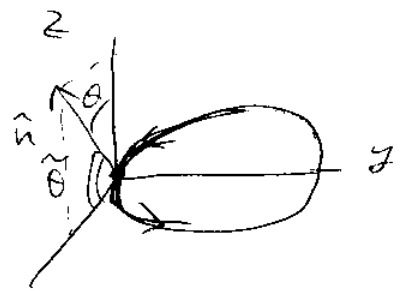
Imagine that at  $t=0$   $\varphi=0$ :

$$\Rightarrow \hat{n} \times (\hat{n} \times \vec{\beta}) \approx -\sin \left( \frac{vt}{\rho} \right) \beta \hat{y} +$$

$$+ \cos \left( \frac{vt}{\rho} \right) \left[ -\hat{x} \beta \cos^2 \theta + \hat{z} \beta \sin \theta \cos \theta \right] \times$$

$$= -\sin \left( \frac{vt}{\rho} \right) \beta \underbrace{\hat{y}}_{\hat{\epsilon}_{\parallel}} + \cos \left( \frac{vt}{\rho} \right) \cdot \cos \theta \beta \underbrace{(\hat{n} \times \hat{y})}_{\hat{\epsilon}_{\perp}}$$

$$= \beta \left[ -\sin \left( \frac{vt}{\rho} \right) \hat{\epsilon}_{\parallel} + \cos \left( \frac{vt}{\rho} \right) \cos \theta \hat{\epsilon}_{\perp} \right]$$



$$\omega \left( t' - \frac{\hat{n} \cdot \vec{r}}{c} \right) = \omega \left[ t' - \frac{\rho}{c} \sin \theta \sin \left( \frac{vt}{\rho} \right) \right] \approx$$

$$\approx \left| \begin{matrix} t \approx 0 \\ \theta \approx \frac{\pi}{2} - \tilde{\theta} \end{matrix} \right. = \omega \left[ t - \frac{\rho}{c} \cos \tilde{\theta} \sin \frac{vt}{\rho} \right] \approx$$

$$\approx \omega \left[ t - \frac{\rho}{c} \left( \frac{vt}{\rho} - \frac{1}{6} \left( \frac{vt}{\rho} \right)^3 \right) \left( 1 - \frac{\tilde{\theta}^2}{2} \right) \right] \approx \omega \left[ t (1 - \beta^2) + \right.$$

$$\left. + \beta \frac{t \tilde{\theta}^2}{2} + \frac{v^3 t^3}{6 c \rho^2} \right] \approx \left| \begin{matrix} \beta \rightarrow 1 \\ \Rightarrow \beta = \sqrt{\frac{1}{\gamma^2} + 1} \approx 1 - \frac{1}{2\gamma^2} \end{matrix} \right.$$

$$\approx \omega \left[ t \frac{1}{2\gamma^2} + \frac{t \tilde{\theta}^2}{2} + \frac{c^2 t^3}{6 \rho^2} \right] \approx \frac{\omega}{2} \left[ t \left( \frac{1}{\gamma^2} + \tilde{\theta}^2 \right) + \right.$$

$$\left. + \frac{c^2 t^3}{3 \rho^2} \right]. \quad \text{We get}$$

$$\frac{dI}{d\omega d\Omega} = \frac{e^2 \omega^2}{4\pi^2 c} \left| -\hat{\epsilon}_{\parallel} A_{\parallel}(\omega) + \hat{\epsilon}_{\perp} A_{\perp}(\omega) \right|^2$$

with

$$A_{\parallel}(\omega) = \frac{c}{\rho} \int_{-\infty}^{\infty} dt' \cdot t' \cdot e^{i \frac{\omega}{2} \left[ t' \left( \frac{1}{\gamma^2} + \tilde{\theta}^2 \right) + \frac{c^2 t'^3}{3 \rho^2} \right]}$$

$$A_{\perp}(\omega) \approx \tilde{\theta} \int_{-\infty}^{\infty} dt' e^{i \frac{\omega}{2} \left[ t' \left( \frac{1}{\gamma^2} + \tilde{\theta}^2 \right) + \frac{c^2 t'^3}{3 \rho^2} \right]}$$

A few more steps lead to