

For an EM wave: $|\vec{k}| = k_0 = \frac{\omega}{c}$, $|\vec{k}'| = k'_0 = \frac{\omega'}{c}$

\Rightarrow if the angle between \vec{k} and \vec{v} is θ in K

and θ' in K' $\Rightarrow \omega' = \gamma(\omega - \beta \omega \cos \theta) \Rightarrow$

$\Rightarrow \omega' = \gamma \omega (1 - \beta \cos \theta)$ Doppler shift

$$\begin{cases} \frac{\omega'}{c} \cos \theta' = \gamma \left(\frac{\omega}{c} \cos \theta - \beta \cdot \frac{\omega}{c} \right) \\ \frac{\omega'}{c} \sin \theta' = \frac{\omega}{c} \sin \theta \end{cases}$$

$\Rightarrow \tan \theta' = \frac{\sin \theta}{\gamma(\cos \theta - \beta)}$

(cf. with light aberration.)

Four - vectors.

We have seen one example: $x^0 = ct, x^1 = x, x^2 = y, x^3 = z$

$$\begin{pmatrix} x_0' \\ x_1' \\ x_2' \\ x_3' \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_0 \\ x_1 \\ x_2 \\ x_3 \end{pmatrix}$$

Definition A 4-vector A^M is a set of 4 quantities (A^0, A^1, A^2, A^3) , which under Lorentz transformation transform as

$$\begin{pmatrix} A^{0'} \\ A^{1'} \\ A^{2'} \\ A^{3'} \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} A^0 \\ A^1 \\ A^2 \\ A^3 \end{pmatrix}.$$

$\Rightarrow A^M, M=0,1,2,3$ is a contravariant vector if it transforms according to:

$$A^{M'} = \frac{\partial x^{M'}}{\partial x^N} A^N \quad (\text{equivalent to above})$$

$\Rightarrow B_M, M=0,1,2,3$ is a covariant vector

if
$$B_{M'} = \frac{\partial x^N}{\partial x^{M'}} B_N$$

Example: $\frac{\partial \varphi}{\partial x^M}$ is a covariant vector as

$$\frac{\partial \varphi}{\partial x^{M'}} = \frac{\partial x^N}{\partial x^{M'}} \frac{\partial \varphi}{\partial x^N}.$$

One can define tensors by

$$A^{M'} B^{N'} = \frac{\partial x^{M'}}{\partial x^\alpha} \frac{\partial x^{N'}}{\partial x^\beta} A^\alpha B^\beta \Rightarrow \text{rank two contravariant}$$

tensor would be $C^{M'N'} = \frac{\partial x^{M'}}{\partial x^\alpha} \frac{\partial x^{N'}}{\partial x^\beta} C^{\alpha\beta}$, etc.

Definition Scalar (inner) product of 2 vectors (16)
is defined by $A_\mu \cdot B^\mu$ (summation assumed)

Let's prove that it's Lorentz-invariant:

$$A'_\mu \cdot B'^\mu = \frac{\partial x^\alpha}{\partial x'^\mu} A_\alpha \frac{\partial x'^\mu}{\partial x^\beta} B^\beta = \frac{\partial x^\alpha}{\partial x^\beta} A_\alpha B^\beta = \delta^\alpha_\beta \cdot$$

$$A_\alpha B^\beta = A_\alpha B^\alpha \quad \text{Q.E.D.}$$

The interval is a scalar: (it's Lorentz-invariant)

$$ds^2 = (dx^0)^2 - (dx^1)^2 - (dx^2)^2 - (dx^3)^2$$

Define the metric tensor by

$$ds^2 \equiv g_{\mu\nu} dx^\mu dx^\nu$$

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = g^{\mu\nu} \left(\begin{array}{l} \text{Minkowski} \\ \text{space} \end{array} \right)$$

Note that $dx_\mu dx^\mu$ is also a Lorentz-scalar.

Identifying $dx_\mu = g_{\mu\nu} dx^\nu$ we see that

$g_{\mu\nu}$ raises and lowers the indices:

$$A_\mu = g_{\mu\nu} A^\nu, \quad A^\mu = g^{\mu\nu} A_\nu$$

(e.g. $x_\mu = g_{\mu\nu} x^\nu, \dots$)

$$g^{\mu\nu} = g^{\mu\alpha} \cdot g_{\alpha\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} = \delta^{\mu\nu}$$

indeed, as $A_\mu B^\mu = \delta^{\mu\nu} A_\mu B^\nu = g^{\mu\nu} A_\mu B_\nu$.

Define an abbreviated notation: $\partial_\mu \equiv \frac{\partial}{\partial x^\mu}$

$$\partial^\mu \equiv \frac{\partial}{\partial x_\mu}$$

$\Rightarrow \partial_\mu \varphi$ is a covariant vector

$\partial^\mu \varphi$ is a contravariant vector (check!)

$\partial_\mu A^\mu$ is Lorentz - invariant

Laplace operator $\frac{\partial^2}{c^2 \partial t^2} - \vec{\nabla}^2 = \partial_\mu \partial^\mu$ is

also Lorentz - invariant.

4 - velocity

Let's define a 4-vector for velocity:

$$dx^\mu = (dx^0, dx^1, dx^2, dx^3) \Rightarrow v^\mu \stackrel{?}{=} \frac{dx^\mu}{dt} \quad ?$$