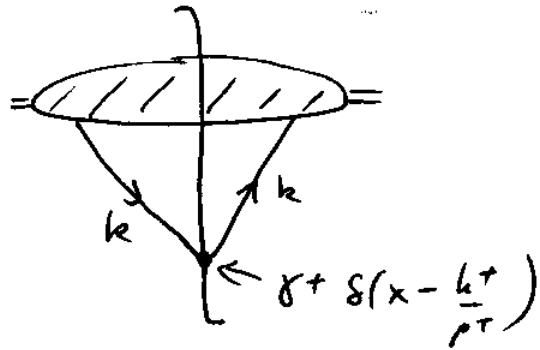


Last time:

QCD-Improved Parton Model (DGLAP Equation)

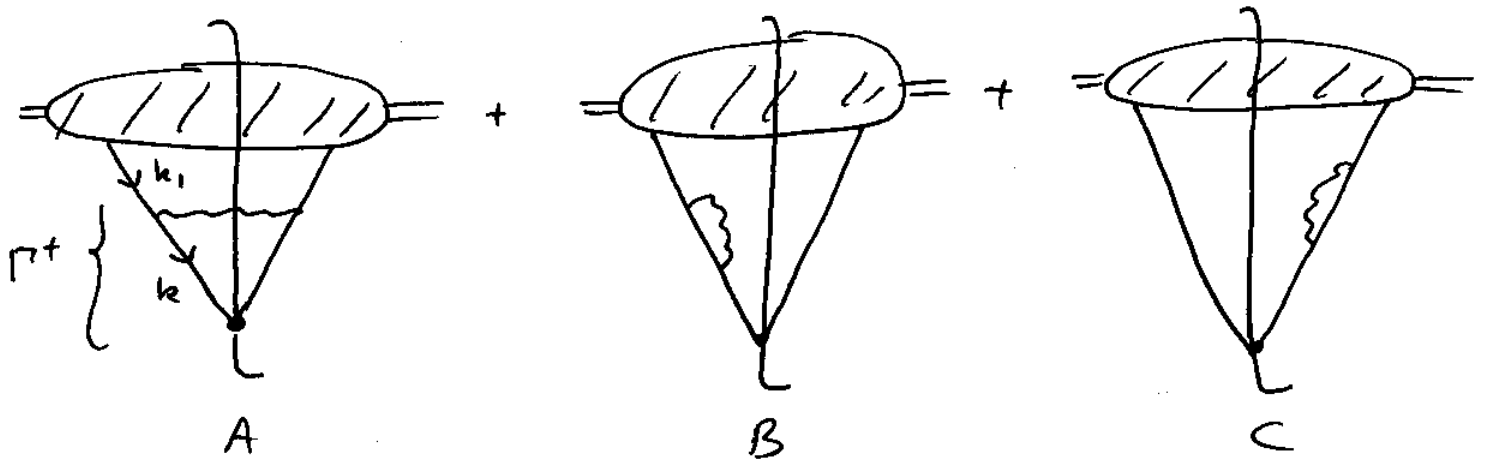
(cont'd)

$$g^f(x) = \frac{\epsilon_f}{p^+}$$



"naive" parton model

QCD $\mathcal{O}(\alpha_s)$ corrections to this picture:



We calculated diagram A assuming that

$\underline{k}^2 \gg \underline{k}_1^2, \Lambda^2 \dots$ We got

$$\Gamma^+ = \gamma^+ \frac{1}{x_1} \frac{\alpha_s C_F}{2\pi} \int \frac{d\underline{k}_\perp^2}{k_\perp^2} \frac{1+z^2}{1-z}$$

with $x_1 = \frac{k_1^+}{p^+}$, $z = \frac{k^+}{k_1^+} \Rightarrow$ as $z \leq 1 \Rightarrow x = zx_1 \leq x_1$

$$\Gamma^+ \sim \alpha_s \int_{\Lambda^2}^{Q^2} \frac{d\underline{k}_\perp^2}{k_\perp^2} \sim \alpha_s \ln \frac{Q^2}{\Lambda^2} \Rightarrow \left. \begin{array}{l} \alpha_s \ll 1 \\ \ln \frac{Q^2}{\Lambda^2} \gg 1 \end{array} \right\} \alpha_s \ln \frac{Q^2}{\Lambda^2} \sim 1$$

leading log approximation (LLA)

\Rightarrow one can resum powers of d_s enhanced by $\ln \frac{Q^2}{\Lambda^2} \Rightarrow$ resum powers of $(d_s \ln \frac{Q^2}{\Lambda^2}) \dots$

We had:

$$g_0^f(x) = \frac{\epsilon_p}{p^+} \int d^4k A_{ab}^f(p, k) (\gamma^+)_{ba} \delta(x - \frac{k^+}{p^+})$$

We now have:

$$S g^f = \frac{\epsilon_p}{p^+} \int d^4k_1 A_{ab}^f(p, k_1) (\gamma^+)_{ba} = \frac{\epsilon_p}{p^+} \int d^4k_1 A_{ab}^f(p, k_1) \cdot$$

$$\cdot (\gamma^+)_{ba} \cdot \frac{1}{x_1} \frac{d_s C_F}{2\pi} \int_{-1}^{Q^2} \frac{dk_{\perp}^2}{k_{\perp}^2} \frac{1+z^2}{1-z} \stackrel{\text{inserting } \int dx_1 \delta(x_1 - \frac{k_1^+}{p^+}) = 1}{=} \frac{\epsilon_p}{p^+} \int \frac{dx_1}{x_1} \cdot \int d^4k_1 \cdot$$

$$\cdot A_{ab}^f(p, k_1) (\gamma^+)_{ba} \delta(x_1 - \frac{k_1^+}{p^+}) \frac{d_s C_F}{2\pi} \int_{-1}^{Q^2} \frac{dk_{\perp}^2}{k_{\perp}^2} \frac{1+z^2}{1-z}$$

$$\sim g_0^f(x_1)$$

$$\Rightarrow S g^f = \int \frac{dx_1}{x_1} \int_{-1}^{Q^2} \frac{dk_{\perp}^2}{k_{\perp}^2} \frac{d C_F}{2\pi} \frac{1 + (\frac{x}{x_1})^2}{1 - \frac{x}{x_1}} \cdot g^f(x_1, Q^2)$$

as $z = \frac{x}{x_1}$

put Q^2 -dep. back.

$$\Rightarrow \frac{\partial g^f(x, Q^2)}{\partial \ln Q^2} = \frac{d C_F}{2\pi} \int \frac{dx_1}{x_1} \frac{1 + (\frac{x}{x_1})^2}{1 - \frac{x}{x_1}} g^f(x_1, Q^2)$$

\Rightarrow but this can not be the full story, as we are forgetting graphs B & C.

$$\Gamma_+ = \gamma_+ \frac{1}{x_1} \frac{\alpha_s(F)}{2\pi} \int_{\underline{k}_1^2}^{Q^2} \frac{d\underline{k}^2}{\underline{k}^2} \frac{1+z^2}{1-z^2}$$

$$\Gamma_+ \sim \alpha_s \cdot \ln(Q^2/\underline{k}_1^2) \sim \alpha_s \ln Q^2$$

$\alpha_s \ll 1$ (perturbation theory, small coupling)

$\ln Q^2 \gg 1$ (DIS with large Q^2)

$\alpha \ln Q^2 \sim 1$ our resummation parameter!

"Leading Logarithmic Approximation"

Remember: we neglected terms suppressed

by $\frac{\underline{k}_1^2}{\underline{k}^2}$, $\frac{\underline{k}_1^4}{\underline{k}^4}$, ... \Rightarrow they give

$$\int_{\underline{k}_1^2}^{Q^2} \frac{d\underline{k}^2}{\underline{k}^4} \underline{k}_1^2 \sim \left(\frac{1}{\underline{k}_1^2} - \frac{1}{Q^2} \right) \underline{k}_1^2 \sim 1 - \frac{\underline{k}_1^2}{Q^2}$$

\swarrow no log \nwarrow higher twist

Old (LO) Parton Model vertex (Mueller vertex)

was $\gamma_+ \delta(x - \frac{k_+}{p_+}) \sim$ same γ_+ matrix as Γ_+

$$\Rightarrow Q^2 \frac{\partial}{\partial Q^2} g^+(x, Q^2) = \frac{\alpha}{2\pi} \int_x^1 \frac{dx_1}{x_1} \gamma_{\gamma\gamma} \left(\frac{x}{x_1} \right) g^+(x_1, Q^2)$$

where $x = \frac{k_+}{p_+}$, $x_1 = \frac{k_{1+}}{p_+} \Rightarrow z = \frac{k_+}{k_{1+}} = \frac{x}{x_1}$

as $z < 1 \Rightarrow x_1 > x$ in the integral.

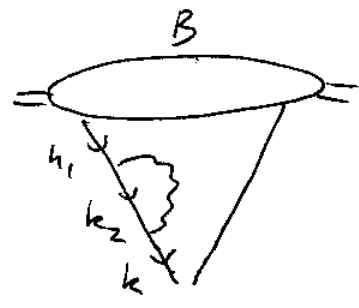
Including the virtual terms (B and C) gives

$$\chi_{gg}(z) = C_F \left(\frac{1+z^2}{1-z} \right)_+$$

where

$$\int_0^1 dz [h(z)]_+ f(z) = \int_0^1 dz h(z) [f(z) - f(1)]$$

Easy to understand:



$\propto \delta(k-k_1)$
 \Downarrow
 $x = x_1$

$\Rightarrow g_f(x_1, Q^2) = g_f(x, Q^2)$
 as $x = x_1$

"real" part,
 \downarrow diagram A

$$Q^2 \frac{\partial}{\partial Q^2} g_f(x, Q^2) = \frac{\alpha C_F}{2\pi} \int_x^1 \frac{dx_1}{x_1} \frac{1 + (x/x_1)^2}{1 - x/x_1} \cdot [g_f(x_1, Q^2) - g_f(x, Q^2)]$$

Virtual corrections, graphs B & C

bare quark state $|\psi_0\rangle = \text{---} \Rightarrow \langle \psi_0 | \psi_0 \rangle = 1$ (59)

(normalization)

dressed quark state $|\psi\rangle = \text{---} + \text{---} + \text{---}$
 $\underbrace{\hspace{1cm}}_{|\psi_0\rangle} \quad \underbrace{\hspace{1cm}}_{|\psi_1\rangle} \quad \underbrace{\hspace{1cm}}_{V|\psi_0\rangle}$

normalization:

$$\langle \psi | \psi \rangle = 1 = \langle \psi_0 | \psi_0 \rangle + \text{---} + \text{---} + \text{---}$$

$$= 1 + \langle \psi_1 | \psi_1 \rangle + 2V \langle \psi_0 | \psi_0 \rangle = 1 + \langle \psi_1 | \psi_1 \rangle + 2V$$

$$\Rightarrow V = -\frac{1}{2} \langle \psi_1 | \psi_1 \rangle$$

$$\Rightarrow \text{graphs } B, C = -\frac{1}{2} A \Rightarrow B + C = -A$$

\approx simply imposed probability conservation!

Defining flavor singlet distribution

$$\Sigma(x, Q^2) \equiv \sum_f [q_f^+(x, Q^2) + q_f^-(x, Q^2)]$$

and flavor non-singlet

$$\Delta_{ff'}(x, Q^2) \equiv q_f^+(x, Q^2) - q_{f'}^+(x, Q^2)$$

we write

$$Q^2 \frac{\partial}{\partial Q^2} \Delta_{ff'}(x, Q^2) = \frac{\alpha(Q^2)}{2\pi} \int_x^1 \frac{dx_1}{x_1} \gamma_{gg}(\frac{x}{x_1}) \cdot \Delta_{ff'}(x_1, Q^2)$$

and

$$Q^2 \frac{\partial}{\partial Q^2} \begin{pmatrix} \Sigma(x, Q^2) \\ G(x, Q^2) \end{pmatrix} = \frac{\alpha(Q^2)}{2\pi} \int_x^1 \frac{dx_1}{x_1} \begin{pmatrix} \gamma_{gg}(\frac{x}{x_1}) & \gamma_{gq}(\frac{x}{x_1}) \\ \gamma_{qg}(\frac{x}{x_1}) & \gamma_{qq}(\frac{x}{x_1}) \end{pmatrix} \cdot \begin{pmatrix} \Sigma(x_1, Q^2) \\ G(x_1, Q^2) \end{pmatrix}$$

Dokshitzer, Gribov, Lipatov, Altarelli, Parisi

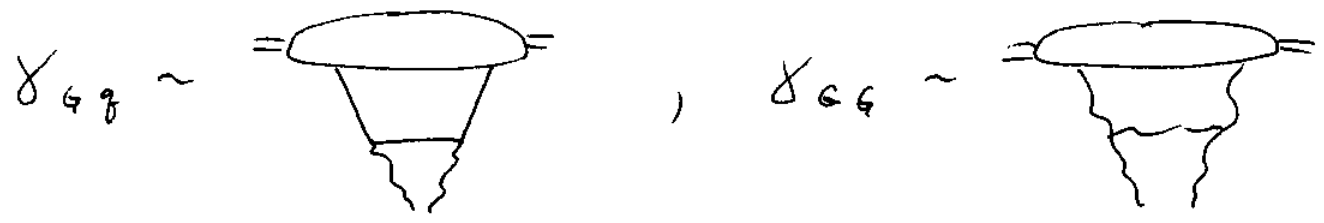
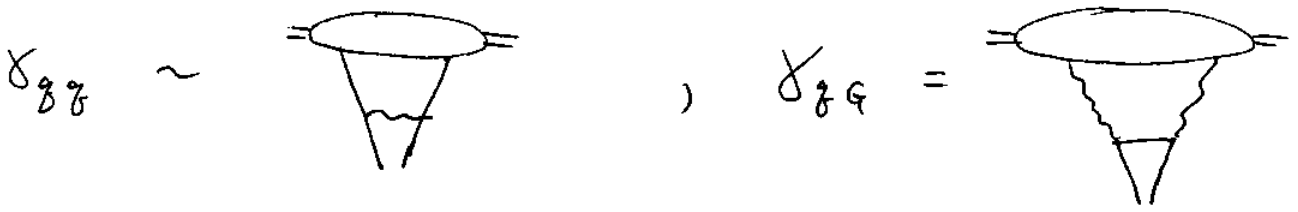
(DGLAP) Equations

GL ~ QED case ~ 172

P, A & P ~ QCD case, '77

$$G(x, Q^2) = \text{[Diagram: a loop with a triangle below it]} \sim \langle A_i A_i \rangle \text{ in } A_+ = 0 \text{ gauge}$$

gluon distribution function

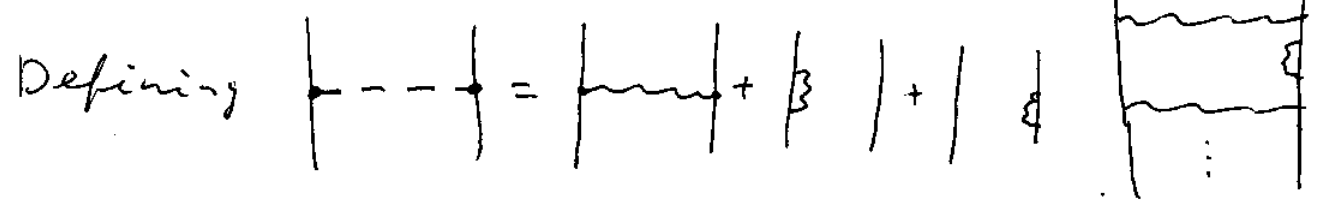


After explicit calculations one gets:

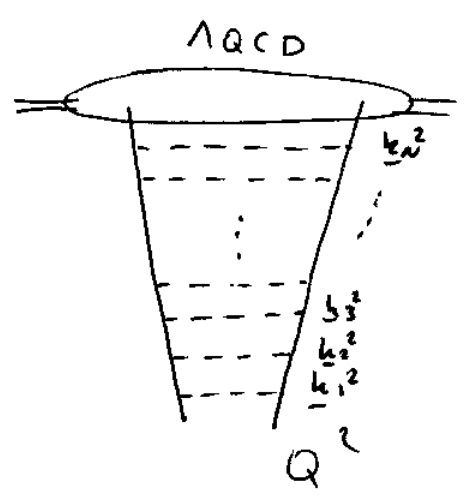
$$\left\{ \begin{aligned} \delta_{gg}(z) &= C_F \left(\frac{1+z^2}{1-z} \right)_+ \\ \delta_{gq}(z) &= C_F \frac{1+(1-z)^2}{z} \\ \delta_{qg}(z) &= N_F [z^2 + (1-z)^2] \\ \delta_{qq}(z) &= 2N_C \left[\frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z) \right] + \frac{11N_C - 2N_F}{6} \delta(z-1) \end{aligned} \right.$$

Note that $\delta_{gq}(z)$ can be obtained from $\delta_{qg}(z)$ by substituting $z \rightarrow 1-z$ and dropping virtual corrections.

Iterate the evolution for $P_+(x, Q^2)$:



We get a ladder diagram:

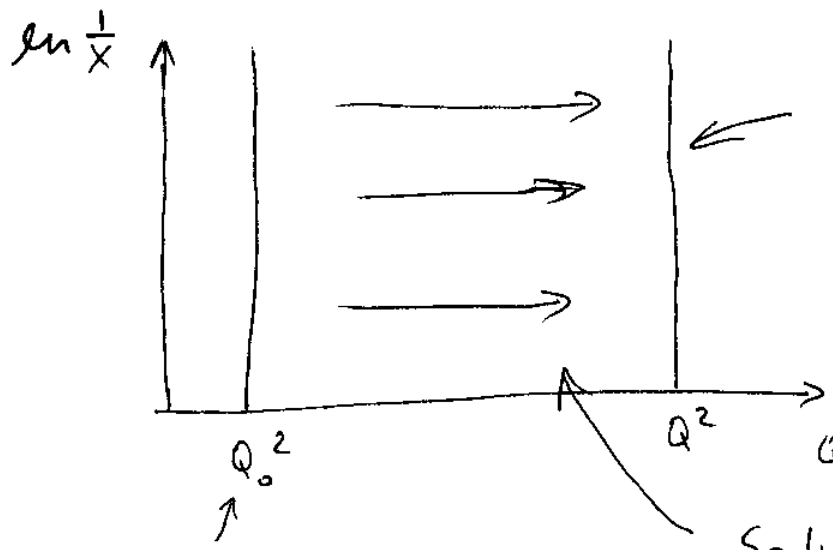


$$Q^2 \gg k_1^2 \gg k_2^2 \gg \dots \gg k_N^2 \gg \Lambda_{QCD}^2$$

DGLAP resums ladder graphs with the ladder connecting scales Q^2 and Λ_{QCD}^2 , $Q^2 \gg \Lambda_{QCD}^2$

such that $\ln Q^2 / \Lambda_{QCD}^2 \gg 1$ and $d_s \ln \frac{Q^2}{\Lambda_{QCD}^2} \sim 1$ is the resummation parameter.

How does DGLAP work?



obtain distribution at whatever Q^2 you wanted.

start with some initial condition

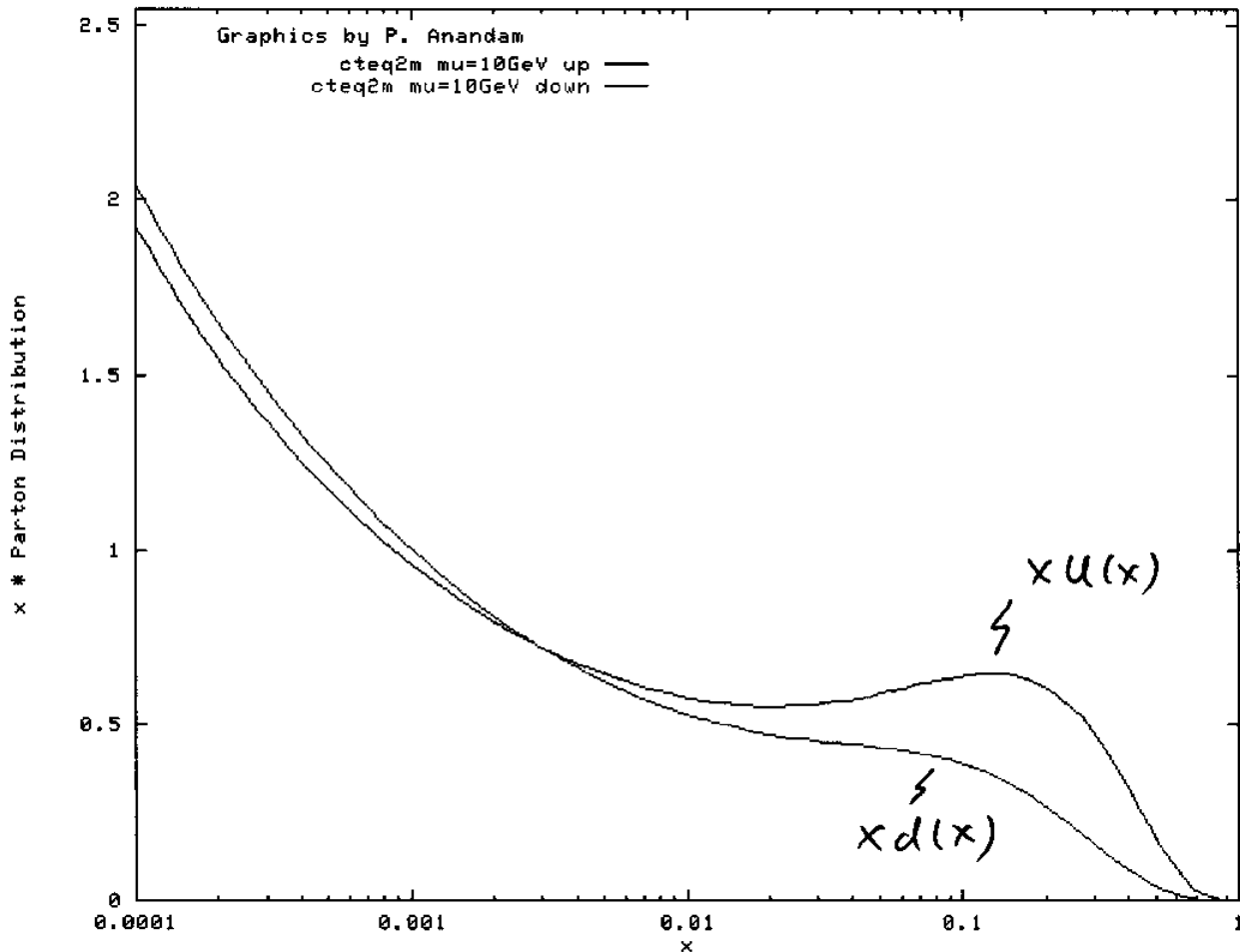
\sim function $P^A(x, Q_0^2)$

Solve the DGLAP equations ("evolve" the distribution function)

\Rightarrow people calculate PDF's (Parton Distribution Functions) & fit the data. See attachments for PDF examples.

Parton Distribution Graph

(Number of graphs plotted since 21 November 2000: 658)



$$Q = 10 \text{ GeV} \Rightarrow Q^2 = 100 \text{ GeV}^2$$

at large - x valence quarks dominate

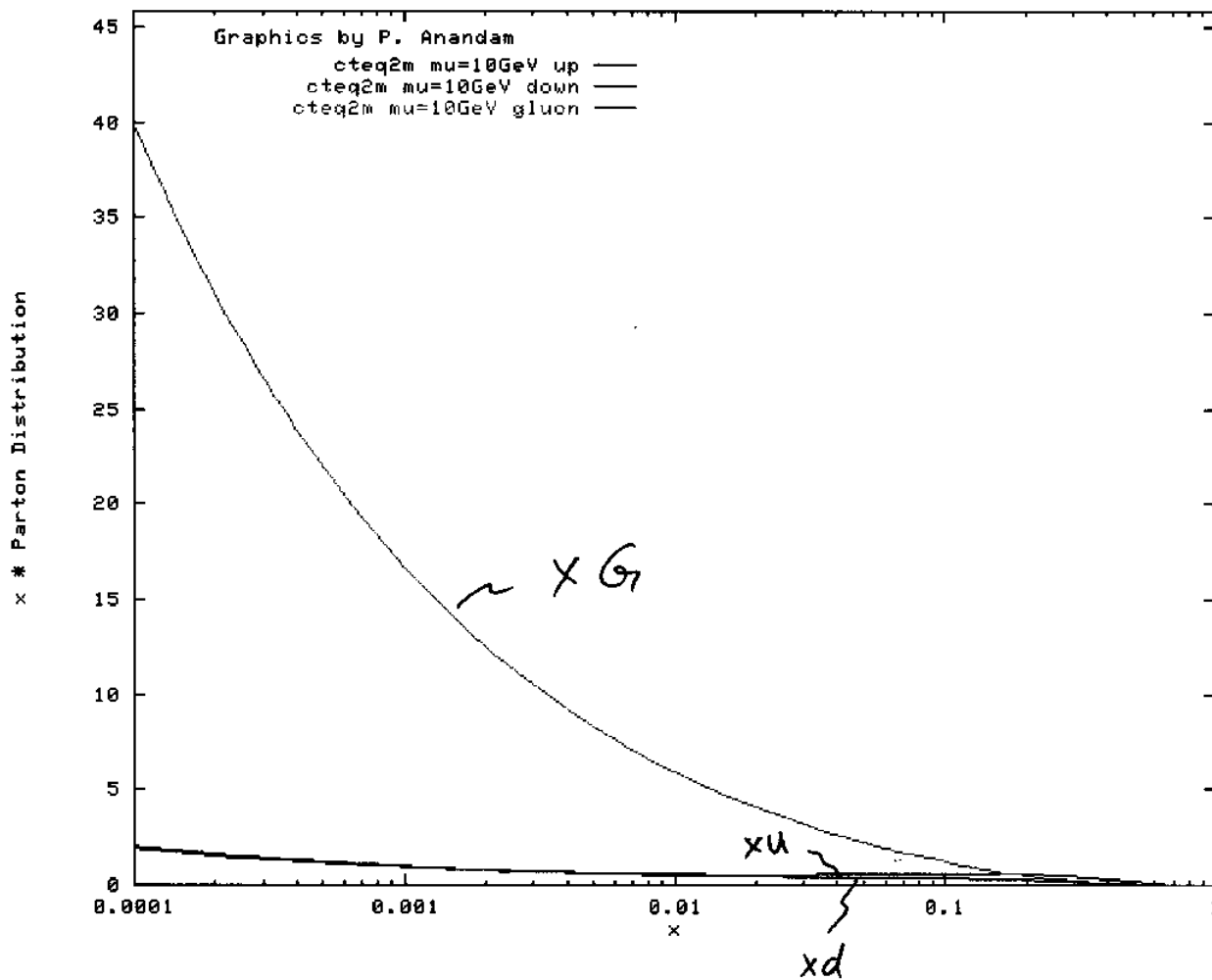
$$\Rightarrow x u_v(x) = 2 x d_v(x)$$

\Rightarrow not so at small - x

go to <http://zebu.uoregon.edu/~parton/partongraph.html>

Parton Distribution Graph

(Number of graphs plotted since 21 November 2000: 659)



the same plot with xG (gluon distribution) plotted as well

now, who's ya daddy?

\Rightarrow at small- x gluons dominate by far...