

Last time:

Canonical Quantization of Free Vector Field (cont'd)

(I) Coulomb gauge: $\vec{\nabla} \cdot \vec{A} = 0$, $A^0 = 0$ (vacuum)

$\Rightarrow A^i(x) \sim$ free fields

$$\pi_i = \frac{\delta \mathcal{L}}{\delta \dot{A}_i} = F^{i0} = E^i \sim \text{canonical momenta}$$

$$[A_i(\vec{x}, t), \pi_j(\vec{y}, t)] = i \left(\delta_{ij} + \frac{\partial_i \partial_j}{\nabla^2} \right) \delta(\vec{x} - \vec{y})$$

$$[A_i(\vec{x}, t), A_j(\vec{y}, t)] = 0 = [\pi_i(\vec{x}, t), \pi_j(\vec{y}, t)]$$

$\square A^i = 0 \sim$ Maxwell eqn's

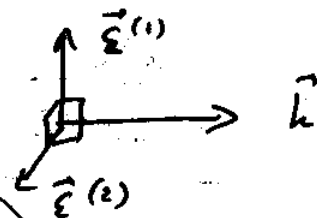
$$\vec{A}(x) = \int \frac{d^3k}{(2\pi)^3 2\epsilon_k} \sum_{\lambda=1}^2 \vec{\epsilon}^{(\lambda)}(\vec{k}) \left[\hat{a}_{\vec{k}, \lambda} e^{-ik \cdot x} + \hat{a}_{\vec{k}, \lambda}^\dagger e^{ik \cdot x} \right]$$

Choose polarizations basis: $\vec{k} \cdot \vec{\epsilon}^{(\lambda)}(\vec{k}) = 0$ (Coulomb gauge)

$$\vec{\epsilon}^{(\lambda)}(\vec{k}) \cdot \vec{\epsilon}^{(\lambda')}(\vec{k}) = \delta_{\lambda\lambda'}$$

$$[\hat{a}_{\vec{k}, \lambda}, \hat{a}_{\vec{k}', \lambda'}^\dagger] = (2\pi)^3 2\epsilon_k \delta_{\lambda\lambda'} \delta(\vec{k} - \vec{k}')$$

$$[\hat{a}_{\vec{k}, \lambda}, \hat{a}_{\vec{k}', \lambda'}] = [\hat{a}_{\vec{k}, \lambda}^\dagger, \hat{a}_{\vec{k}', \lambda'}^\dagger] = 0$$



$\lambda = 1, 2 \Rightarrow$ only 2 physical polarizations

$$H = \int \frac{d^3k}{(2\pi)^3 2\epsilon_k} \sum_{\lambda=1}^2 \epsilon_k \hat{a}_{\vec{k}, \lambda}^\dagger \hat{a}_{\vec{k}, \lambda}$$

(II) Lorenz gauge: $\partial_\mu A^\mu = 0 \Rightarrow A^\mu(x)$ free fields

$$\Rightarrow \pi^\mu = \frac{\delta \mathcal{L}}{\delta \dot{A}_\mu} = F^{\mu 0} \Rightarrow \pi^0 = 0 \sim \text{bad}$$

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\lambda}{2} (\partial_\mu A^\mu)^2 \sim \text{new Lagrangian}$$

EOM: $\square A^\mu - (1-\lambda) \partial^\mu \partial_\nu A^\nu = 0$

$\lambda = 1 \Rightarrow$ Feynman gauge $\Rightarrow \square A^\mu = 0$ is back.

$\pi^0 = -\partial_\mu A^\mu \sim$ non-zero (don't fix any gauge)

$$\Rightarrow [A_\mu(\vec{x}, t), \pi_\nu(\vec{y}, t)] = i g_{\mu\nu} \delta(\vec{x} - \vec{y})$$

$$[A_\mu(\vec{x}, t), A_\nu(\vec{y}, t)] = [\pi_\mu(\vec{x}, t), \pi_\nu(\vec{y}, t)] = 0$$

$\square A^\mu = 0 \Rightarrow$

$$A_\mu(x) = \int \frac{d^3k}{(2\pi)^3 2\epsilon_k} \sum_{\lambda=0}^3 \epsilon_\mu^{(\lambda)}(\vec{k}) \left[\hat{a}_{\vec{k}, \lambda} e^{-ik \cdot x} + \hat{a}_{\vec{k}, \lambda}^\dagger e^{ik \cdot x} \right]$$

$\epsilon^{(\lambda)} \cdot k \neq 0 \Rightarrow$ for $k^\mu = (k, 0, 0, k)$ have

$$\epsilon^{(0)} = (1, 0, 0, 0), \quad \epsilon^{(1)} = (0, 1, 0, 0), \quad \epsilon^{(2)} = (0, 0, 1, 0)$$

time-like

$$\epsilon^{(3)} = (0, 0, 0, 1), \quad \epsilon^{(1,2)} \cdot k = 0 \quad (\text{transverse})$$

longitudinal

$$[\hat{a}_{\vec{k}, \lambda}, \hat{a}_{\vec{k}', \lambda'}^\dagger] = -g_{\lambda\lambda'} (2\pi)^3 2\epsilon_k \delta(\vec{k} - \vec{k}')$$

(all others zero)

One then gets:

$$[\hat{a}_{\vec{k},\lambda}, \hat{a}_{\vec{k}',\lambda'}^\dagger] = -g_{\lambda\lambda'} (2\pi)^3 2\varepsilon_k \delta(\vec{k}-\vec{k}')$$

all other commutators are zero.

if $k^\mu = (k, 0, 0, k) \Rightarrow$

$$\varepsilon^0 = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad \varepsilon^1 = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}, \quad \varepsilon^2 = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \quad \varepsilon^3 = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$$

$\varepsilon^{1,2,k} = 0$ (transverse)

$\varepsilon^3 \sim$ longitudinal, $\varepsilon^0 \sim$ time-like.

Fock space: $\hat{a}_{\vec{k},\lambda}^\dagger |0\rangle, \dots$

$$|1,\lambda\rangle = \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} f(\vec{k}) \hat{a}_{\vec{k},\lambda}^\dagger |0\rangle \quad \sim \text{one-photon state}$$

Problem: $\langle 1, \lambda=0 | 1, \lambda=0 \rangle = \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} \frac{d^3k'}{(2\pi)^3 2\varepsilon_{k'}} f(\vec{k}) f^*(\vec{k}')$

$$\langle 0 | \hat{a}_{\vec{k},\lambda=0}^\dagger \hat{a}_{\vec{k},0}^\dagger |0\rangle = - \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} |f(\vec{k})|^2 \langle 0|0\rangle$$

"
$$[\hat{a}_{\vec{k},0}^\dagger, \hat{a}_{\vec{k},0}^\dagger] = - (2\pi)^3 2\varepsilon_k \delta(\vec{k}-\vec{k}')$$

\Rightarrow negative norm states!

$$H = \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} \varepsilon_k \left[\sum_{\lambda=1}^3 \hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda} - \hat{a}_{\vec{k},0}^\dagger \hat{a}_{\vec{k},0} \right] \sim \text{negative energy problem...}$$

=> way out: demand that while $\partial_\mu A^\mu \neq 0$ at the operator level, for physical states one

has $\partial_\mu A^{(+)\mu} |\psi\rangle = 0$

where (+) denotes positive-energy part of A_μ (the term with $\hat{a}_{\vec{k},\lambda}$).

$\langle \psi | \partial_\mu A^\mu | \psi \rangle = \langle \psi | \partial_\mu A^{(+)\mu} + \partial_\mu A^{(-)\mu} | \psi \rangle = 0.$

(Gupta & Bleuler method). $\langle \psi | \partial_\mu A^{(-)\mu} = [\partial_\mu A^{(+)\mu} | \psi \rangle]^\dagger = 0$

$\Rightarrow \partial_\mu A^{(+)\mu} | \psi \rangle = 0 \Rightarrow \sum_{\lambda=0}^3 k^\mu \epsilon_\mu^\lambda(\vec{k}) \hat{a}_{\vec{k},\lambda} | \psi \rangle = 0$

$\Rightarrow (k^\mu \epsilon_\mu^0 \hat{a}_{\vec{k},0} + k^\mu \epsilon_\mu^3 \hat{a}_{\vec{k},3}) | \psi \rangle = 0$

as $\epsilon^{1,2} \cdot k = 0$. Now $k \cdot \epsilon^0 = k^0 = -k \cdot \epsilon^3$

$\Rightarrow (\hat{a}_{\vec{k},0} - \hat{a}_{\vec{k},3}) | \psi \rangle = 0$

=> physical states are mixtures of longitudinal & time-like photons

$\Rightarrow \langle \psi | \hat{a}_{\vec{k},0}^\dagger \hat{a}_{\vec{k},0} | \psi \rangle = \langle \psi | \hat{a}_{\vec{k},3}^\dagger \hat{a}_{\vec{k},3} | \psi \rangle$

=> only transverse photons contribute in H:

$\langle \psi | H | \psi \rangle = \int \frac{d^3k}{(2\pi)^3 2\epsilon_k} \epsilon_k \cdot \langle \psi | \sum_{\lambda=1}^2 \hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda} | \psi \rangle \geq 0$

\Rightarrow no negative energy problem any more.

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Massive Vector Field

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{m^2}{2} A_\mu A^\mu$$

m is the mass

$$\partial_\mu F^{\mu\nu} + m^2 A^\nu = 0$$

Proca eqn's.

$\partial_\mu A^\mu = 0$ always \Rightarrow

$$A_\mu(x) = \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} \sum_{\lambda=1}^3 \varepsilon_\mu^{(\lambda)}(\vec{k}) \left[\hat{a}_{\vec{k},\lambda} e^{-i\vec{k}\cdot\vec{x}} + \hat{a}_{\vec{k},\lambda}^\dagger e^{i\vec{k}\cdot\vec{x}} \right]$$

with $[\hat{a}_{\vec{k},\lambda}, \hat{a}_{\vec{k}',\lambda'}^\dagger] = \delta_{\lambda\lambda'} (2\pi)^3 2\varepsilon_k \delta(\vec{k}-\vec{k}')$

$\vec{k} \cdot \varepsilon^{(\lambda)} = 0$ for $\lambda = 1, 2, 3 \Rightarrow$ in rest frame

have $k^\mu = (m, 0, 0, 0) \Rightarrow$

$$\varepsilon_\mu^{(1)} = (0, 1, 0, 0), \quad \varepsilon_\mu^{(2)} = (0, 0, 1, 0), \quad \varepsilon_\mu^{(3)} = (0, 0, 0, 1).$$

3 degrees of freedom (all physical)

$$H = \int \frac{d^3k}{(2\pi)^3 2\varepsilon_k} \varepsilon_k \sum_{\lambda=1}^3 \hat{a}_{\vec{k},\lambda}^\dagger \hat{a}_{\vec{k},\lambda}$$

Time evolution: $-i\partial_t A_\mu = [H, A_\mu]$ for both massive & massless vector fields.