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7701 Lecture 18

Lecture Plan

- Continue through Fourier transform discussion
 - key examples from last time if incomplete
 - FT as a linear operator \rightarrow apply to PDE's
 - 3D FT and homework problem
 - Wave equation
 - Sine and Cosine transforms
- Wave equation with initial conditions, eqn (7.1)
- Wave equation with inhomogeneous source, eqn (7.5) (7.6)

Before class:

- Start up 7701 page + Mathematica
- notebook: diffusion.nb, p100-101

On board:

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \quad \text{cf.} \quad G(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(t) e^{i\omega t} dt$$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k) e^{ikx} dk \quad \text{cf.} \quad g(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} G(\omega) e^{-i\omega t} dt$$

- Recap: Carrying out transforms are just exercises in doing integrals.
- often contour integrals because $\int_{-\infty}^{\infty}$ integrals
 - Mathematica can do many using Fourier Transform or Inverse Fourier Transform (convention $e^{i\omega t} / e^{-i\omega t}$)
 - \Rightarrow be careful of convention \rightarrow you may get complex conjugate
 - Several examples you need to be familiar with (problem set):
 - exponential (11), Gaussian (completing the square) (14),
 - delta function (13). Also, use of δ Functions.

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Revisit application to partial differential equations (diffusion example)

Key: consider linear differential equations \Rightarrow taking a Fourier transform is a linear operation. [generalize: other integral transforms]

Denote it $F[\] \Rightarrow F[f] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx$

Apply to diffusion equation from last time
(mass m of dye initially at $x=l$ in long pipe of water. Find $t>0$ distribution)

$F\left[\frac{\partial p(x,t)}{\partial t} - D \frac{\partial^2 p(x,t)}{\partial x^2} = 0\right]$ \leftarrow cf. Bk 6, diffusion equation with additional term

linear: $F[f+g] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (f(x)+g(x)) e^{-ikx} dx$
 $= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx + \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(x) e^{-ikx} dx$
 $= F[f] + F[g]$

Important: f and g here are just patterns to match; they can be anything.
This is how Mathematics works! (example notebook if time)

$\Rightarrow F\left[\frac{\partial p}{\partial t}\right] - F\left[D \frac{\partial^2 p}{\partial x^2}\right] = 0$ ($F[0]=0$)

Now $F[af(x)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (af(x)) e^{-ikx} dx = a \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx\right) = a F[f]$

$\Rightarrow F\left[\frac{\partial p}{\partial t}\right] - D F\left[\frac{\partial^2 p}{\partial x^2}\right] = 0$ \leftarrow still partial since dependence on k

But $F\left[\frac{\partial p}{\partial t}\right] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial p(x,t)}{\partial t} e^{-ikx} dx = \frac{d}{dt} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} p(x,t) e^{-ikx} dx = \frac{d}{dt} F[f]$

But $F\left[\frac{\partial p}{\partial x}\right] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial p}{\partial x} e^{-ikx} dx \stackrel{int. by parts}{=} - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} p(x,t) \frac{d}{dx} e^{-ikx} dx + \left. \frac{p(x,t) e^{-ikx}}{\sqrt{2\pi}} \right|_{-\infty}^{\infty}$
 $= +ik \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} p(x,t) e^{-ikx} dx = +ik F[f]$
two minus signs

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So we have $\frac{d}{dt} [F_L] - D(k)^2 [F_L] = 0$
 or finally $\frac{d}{dt} \tilde{f}(k,t) + k^2 \tilde{f}(k,t) = 0$

Consider this for each $k \Rightarrow 1^{st}$ order like $\frac{df}{dt} = -\alpha^2 f \Rightarrow f(t) = f_0 e^{-\alpha^2 t}$

$\Rightarrow \tilde{f}(k,t) = \tilde{f}(k,0) e^{-Dk^2 t}$ here k^2 in exponent is just a constant

Find $\tilde{f}(k,0)$ from initial conditions

$\tilde{f}(k,0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} p(x,0) e^{-ikx} dx = \frac{1}{\sqrt{2\pi}} \frac{m}{A} e^{-ikl}$
cross sectional area $\rightarrow \frac{m}{A} \delta(x-l)$

$\Rightarrow \tilde{f}(k,t) = \frac{m}{A} \frac{e^{-ikl}}{\sqrt{2\pi}} e^{-k^2 t}$

Inverse transform: $p(x,t) = \frac{m}{A} \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikx - ikl - k^2 t} dk = \frac{m}{A} \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i(\sqrt{t}k + \frac{i(k-x)^2}{2\sqrt{t}}) - \frac{k^2 t}{4}} dk$
complete sphere

check $-(\sqrt{t}k + \frac{i(k-x)^2}{2\sqrt{t}})^2 = -t k^2 - 2\sqrt{t}k \frac{i(k-x)^2}{2\sqrt{t}} - \frac{i^2(k-x)^2}{4} = -t k^2 - i(k-x)^2 k - \frac{(k-x)^2}{4}$ ✓

$\Rightarrow p(x,t) = \frac{m}{A} \frac{1}{2\pi} e^{-\frac{i}{4}(k-x)^2} \frac{1}{\sqrt{t}} \cdot \sqrt{\pi} = \frac{m}{A} \frac{1}{2\sqrt{t}} e^{-(x-l)^2/40t}$
shift integral; no poles

How do we check?

(1) satisfies initial condition? $\delta(y) = \frac{n}{\sqrt{\pi}} e^{-ny^2} \xrightarrow{n \rightarrow \infty} \delta(y)$

$y = x-l, n^2 = 1/40t \Rightarrow n = 1/2\sqrt{t} \Rightarrow \frac{m}{A} \frac{n}{\sqrt{\pi}} e^{-n(x-l)^2} \xrightarrow{n \rightarrow \infty} \frac{m}{A} \delta(x-l)$ ✓

(2) satisfies differential equation? Not obviously! Only $\frac{(x-l)^2}{40t}$ matters

$\frac{d}{dt} \left[\frac{1}{2} e^{-\frac{(x-l)^2}{40t}} \right] = \left[\frac{1}{\sqrt{40}} \frac{(x-l)^2}{40} - \frac{1}{2} \right] e^{-\frac{(x-l)^2}{40t}} = 0 \left[\frac{-2}{40t^{3/2}} + \frac{4(x-l)^2}{160t^2 \sqrt{40}} \right] e^{-\frac{(x-l)^2}{40t}}$ ✓ works

Show spreading of initial distribution.

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diffusion 1.nb

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Clear["Global`*"]
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Animate the solution to the diffusion example of a mass m of dye starting at $x=l$ and diffusing in a pipe of cross sectional area A with diffusion constant Df .

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rho[x_, t_, Df_] := (m / A) * 1 / (2 Sqrt[Pi Df t]) Exp[-(x - l)^2 / (4 Df t)]
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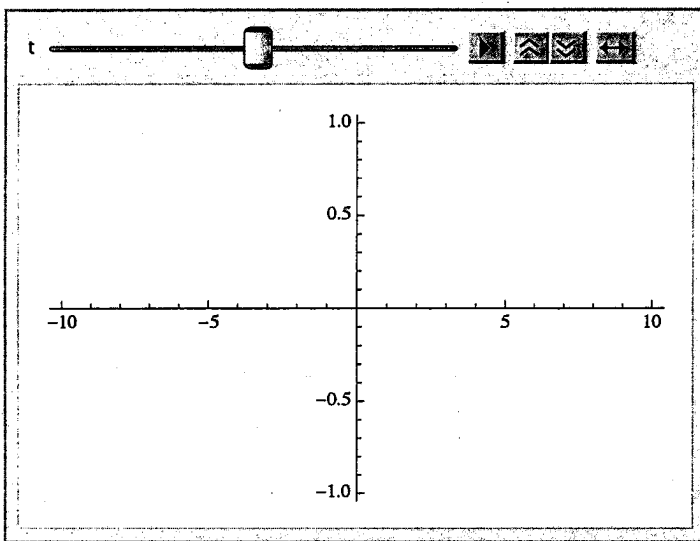
Pick some values:

```
m = 1; A = 1; l = 1 / 2;
```

```
xmin = -10; xmax = +10;  
tmin = 0.001; tmax = 10;
```

```
Df1 = 0.5;
```

```
Animate[Plot[rho[x, t, Df1], {x, xmin, xmax}, PlotRange -> Full],  
{t, tmin, tmax}, AnimationRunning -> False, AnimationDirection -> ForwardBackward]
```



```
Df2 = 0.1;
```

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Fourier Transform of Wave Equation

Let's revisit the wave equation $\frac{\partial^2 y}{\partial t^2} = v^2 \frac{\partial^2 y}{\partial x^2}$. Should we Fourier transform with respect to t or x ? Answer: both!

$$\tilde{y}(k, \omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx e^{-ikx} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dt e^{i\omega t} y(x, t) \quad \text{note } e^{-ikx} e^{i\omega t} \text{ convention}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx dt e^{-ikx + i\omega t} y(x, t)$$

inverse: $y(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk \int_{-\infty}^{\infty} d\omega \tilde{y}(k, \omega) e^{ikx - i\omega t} dk d\omega$

$$\Rightarrow (-i\omega)^2 \tilde{y} = v^2 (ik)^2 \tilde{y} \Rightarrow \tilde{y}(\omega^2 - v^2 k^2) = 0 \Rightarrow \boxed{\tilde{y}(\omega + vk)(\omega - vk) = 0}$$

Now what? $\tilde{y} = 0$ is a solution, but not the one we want. Looks like $\omega = \pm vk$, but what does that imply?

When is a Fourier transform zero except for special values of ω ? Answer: when it has delta functions!

$$\Rightarrow \boxed{\tilde{y}(k, \omega) = A(k) \delta(\omega - vk) + B(k) \delta(\omega + vk)}$$

check: $(\omega + vk)(\omega - vk) \tilde{y}(k, \omega) = A(k) (\omega + vk)(\omega - vk) \delta(\omega - vk) + B(k) (\omega + vk)(\omega - vk) \delta(\omega + vk)$

$$\Rightarrow y(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk \int_{-\infty}^{\infty} d\omega e^{ikx - i\omega t} [A(k) \delta(\omega - vk) + B(k) \delta(\omega + vk)]$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dk [A(k) e^{ik(x-vt)} + B(k) e^{ik(x+vt)}]$$

\nwarrow moves to right \swarrow moves to left

What if we had just transformed wrt x and then solved the equation wrt t ?

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Extension to 3-dimensions

eg. homework problem $\rho(\vec{r}) = \rho_0 \frac{e^{-r/a}}{4\pi r}$

If we think of this as $\rho(x, y, z)$, the extension is trivial $\Rightarrow F_x F_y F_z [\rho]$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx e^{-ik_x x} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dy e^{-ik_y y} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dz e^{-ik_z z} \rho(x, y, z)$$

$$= \left(\frac{1}{2\pi}\right)^{3/2} \int d^3r e^{-i\vec{k}\cdot\vec{r}} \rho(\vec{r})$$

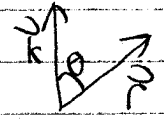
Linear properties carry through. Gradients and the like become:

$$F[\vec{\nabla} f] = i\vec{k} F[f]$$

$$F[\nabla^2 f] = -k^2 F[f]$$

$$F[\vec{\nabla} \times \vec{f}] = i\vec{k} \times F[\vec{f}]$$

But what if $\rho(\vec{r}) = \rho(r)$?
Then spherical coordinates makes sense!

$\vec{k} \cdot \vec{r} = kr \cos\theta$ so choose \hat{z} -axis aligned with \vec{k} : 

$$\Rightarrow \int d^3r = \int_0^{\infty} r^2 dr \int_{\cos\theta}^{\pi} \sin\theta d\theta \int_0^{2\pi} d\phi = \int_0^{\infty} r^2 dr \int_{-1}^1 dx \int_0^{2\pi} d\phi \quad x = \cos\theta$$

Homework problem is typical example of what you'll encounter.

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Sine and Cosine Fourier Transforms

Just as with Fourier series, we can also have Sine and Cosine Fourier transforms, which add an assumption about $f(x)$ compared to $f(-x)$. Two general cases:

1. $f(x)$ is known to be even or odd about $x=0$.
2. $f(x)$ is defined only for $x > 0$.

Cosine

$$F_c(k) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(x) \cos kx \, dx = F_c(-k)$$

For cosine \rightarrow

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} F_c(k) \cos kx \, dk = f(-x) \quad [\text{even extension to } x < 0]$$

Sine

$$F_s(k) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(x) \sin kx \, dx$$

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} F_s(k) \sin kx \, dk$$

What is different? We have integrals from $\int_0^{\infty} \Rightarrow$ surface terms (in sum).

remember
[]'s \rightarrow
mean the linear
operator

$$F_c \left[\frac{df}{dx} \right] = -\sqrt{\frac{2}{\pi}} f(0) + k F_s[f] \quad \text{check!}$$

$$F_s \left[\frac{df}{dx} \right] = -k F_c[f]$$

$$\begin{aligned} F_c \left[\frac{d^2f}{dx^2} \right] &= -\sqrt{\frac{2}{\pi}} f'(0) + k F_s \left[\frac{df}{dx} \right] = -\sqrt{\frac{2}{\pi}} f'(0) - k^2 F_c[f] \\ &= -\sqrt{\frac{2}{\pi}} f'(0) - k^2 F_c(k) \end{aligned}$$

$$F_s \left[\frac{d^2f}{dx^2} \right] = -k F_c \left[\frac{df}{dx} \right] = k \sqrt{\frac{2}{\pi}} f(0) - k^2 F_s(k)$$

\Rightarrow Most useful when only even or odd derivatives (so that $F_s \Rightarrow F_s$ only, etc.) and choose according to initial conditions (eg. whether $f(0)$ or $f'(0)$ is known)

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So suppose we have $\frac{d^2y}{dx^2} - \alpha^2 y = 0$ [so we know the solutions]

For two cases: i) $y(0) = y_0$ and $y \rightarrow 0$ as $x \rightarrow \infty$
ii) $y'(0) = y'_0$ and $y \rightarrow 0$ as $x \rightarrow \infty$ } boundary (initial) conditions

Note: Why don't we need two initial conditions?

Because using FT implies that $F(x) \rightarrow 0$ as $x \rightarrow \infty$.

Otherwise use Laplace transform (more later)

i) $y(0)$ is known, so use F_s

$$\Rightarrow F_s \left[\frac{d^2y}{dx^2} - \alpha^2 y = 0 \right] \Rightarrow -k^2 F_s(y) + k \sqrt{\frac{2}{\pi}} y_0 - \alpha^2 F_s(y) = 0$$

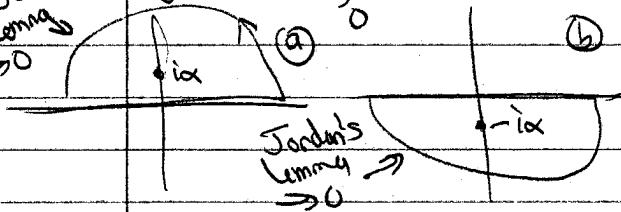
$$\Rightarrow F_s(y) = \sqrt{\frac{2}{\pi}} \frac{ky_0}{k^2 + \alpha^2}$$

$k \sin kx$ is even (factor of $\frac{1}{2}$)

Now we just find the inverse

Jordan's Lemma $\rightarrow 0$

$$y(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} dk \sqrt{\frac{2}{\pi}} \frac{ky_0}{k^2 + \alpha^2} \sin kx = \frac{y_0}{\pi} \int_{-\infty}^{\infty} dk \frac{k}{k^2 + \alpha^2} \frac{e^{ikx} - e^{-ikx}}{2i}$$



$$= \frac{y_0}{2\pi i} \left[(+2\pi i) \frac{i\alpha}{2i\alpha} e^{i(i\alpha)x} - (-2\pi i) \frac{-i\alpha}{2i\alpha} e^{-i(-i\alpha)x} \right]$$

$$= y_0 e^{-\alpha x}$$

which is what we would have written down.

ii) Now $y'(0)$ is known, so use F_c

check? $y(0) = y_0$ ✓

$$\Rightarrow F_c \left[\frac{d^2y}{dx^2} - \alpha^2 y = 0 \right] \Rightarrow -k^2 F_c(y) - \sqrt{\frac{2}{\pi}} y'_0 - \alpha^2 F_c(y) = 0$$

$y \rightarrow 0$ as $x \rightarrow \infty$ satisfies diff eq.

$$\Rightarrow F_c(y) = -\sqrt{\frac{2}{\pi}} \frac{y'_0}{k^2 + \alpha^2} \Rightarrow \text{find inverse} \Rightarrow \text{same contours}$$

$$y(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} dk -\sqrt{\frac{2}{\pi}} \frac{y'_0}{k^2 + \alpha^2} \cos kx = -\frac{y'_0}{\pi} \int_{-\infty}^{\infty} dk \frac{1}{k^2 + \alpha^2} \frac{e^{ikx} + e^{-ikx}}{2}$$

$$= -\frac{y'_0}{2\pi i} \left[(+2\pi i) \frac{1}{2i\alpha} e^{i(i\alpha)x} + (-2\pi i) \frac{1}{2i\alpha} e^{-i(-i\alpha)x} \right] = -\frac{y'_0}{\alpha} e^{-\alpha x}$$

Check $\frac{dy}{dx} = y'_0 e^{-\alpha x} /_{x=0} = y'_0$ ✓