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7701 Lecture 22

Mathematica items

- You can explicitly do 3-D integrations in other coordinate systems by including the measure.
- E.g. spherical coordinates: volume integral spherical on
 • convert integrand as needed mathematica example page.



$$z = r \cos \theta, \quad x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi$$

- Suppose $f(\vec{x}) = r^3 \sin^4 \theta \cos^3 \theta \sin^2 \phi$ to be integrated over a hemisphere of radius 2 with base in xy plane (choose spherical because limits are easy)

⇒ Integrate $\int (r^2 \times \sin[\theta]) \times r^3 \times \sin[\theta]^{14} \times \cos[\theta]^{13} \times \sin[\phi]^{12} \times r^2, \{r, 0, 2\}, \{\theta, 0, \pi/2\}, \{\phi, 0, 2\pi\}$

measure →

hemispheric
 $0 \leq \theta \leq \pi/2$

full range

- Mathematica also knows the coordinate systems, ^{your names} also "cylindrical" and more

Laplacian $[f(r, \theta, \phi), \{r, \theta, \phi\}, \text{"Spherical"}]$

- Warning: it doesn't get the delta function in $\nabla^2 1/r$!

Recap:
$$\vec{E}(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} \cdot \rho(\vec{x}')$$

where
$$\rho(\vec{x}') = \sum_{i=1}^N q_i \delta^3(\vec{x}' - \vec{x}_i)$$

} compare later to boundary conditions

- Note that this is like our convolution example:

- If we know how to find the response (here the \vec{E} field) from a point charge (we'll call this a Green's function) then the full response follows from a well-defined integral over the function of point charges (here $\rho(\vec{x}')$)

- ⇒ solution to electrostatics by "summation"
 - requires knowing $\rho(\vec{x}')$

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Let's return to calculate $\nabla^2 \frac{1}{|\vec{x}-\vec{x}'|}$:

The problem happens with $\vec{x}' = \vec{x}$ and the denominator is zero. \Rightarrow let's define it by a limit with a "regulator" ϵ (cf. an ultraviolet cutoff in field theory)

$$\nabla^2 \frac{1}{|\vec{x}-\vec{x}'|} = \lim_{\epsilon \rightarrow 0} \nabla^2 \frac{1}{\sqrt{(\vec{x}-\vec{x}')^2 + \epsilon^2}} \quad (\text{lim implicit from here on})$$

$$= (\partial_x^2 + \partial_y^2 + \partial_z^2) \frac{1}{\sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2 + \epsilon^2}} \quad \text{just carry it out!}$$

claim

$$= \frac{-3}{[(x-x')^2 + (y-y')^2 + (z-z')^2 + \epsilon^2]^{3/2}} + \frac{3[(x-x')^2 + (y-y')^2 + (z-z')^2]}{[(x-x')^2 + (y-y')^2 + (z-z')^2 + \epsilon^2]^{5/2}}$$

$$= -3 \frac{\epsilon^2}{[(x-x')^2 + (y-y')^2 + (z-z')^2 + \epsilon^2]^{5/2}} \xrightarrow{\epsilon \rightarrow 0} \begin{cases} 0 & |\vec{x}-\vec{x}'|^2 \neq 0 \\ \infty & \vec{x} = \vec{x}' \end{cases}$$

\Rightarrow looks like a delta function, we'll just check that it integrates to 1

all space

$$\int d^3x -3 \frac{\epsilon^2}{[(\vec{x}-\vec{x}')^2 + \epsilon^2]^{5/2}} \xrightarrow{\text{change variables } \vec{x} \rightarrow \vec{x}-\vec{x}'} \int d^3x -3 \frac{\epsilon^2}{(\vec{x}^2 + \epsilon^2)^{5/2}} \xrightarrow{\text{spherical}} -3\epsilon^2 4\pi \int_0^\infty \frac{r^2 dr}{(r^2 + \epsilon^2)^{5/2}}$$

Mathematica $\Rightarrow = -4\pi$

$$\Rightarrow \lim_{\epsilon \rightarrow 0} \frac{-3\epsilon^2}{(\vec{x}^2 + \epsilon^2)^{5/2}} = -4\pi \delta^3(\vec{x})$$

$$\boxed{\nabla^2 \frac{1}{|\vec{x}-\vec{x}'|} = -4\pi \delta^3(\vec{x}-\vec{x}')} \quad \text{**}$$

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Aside: Could we have done this directly in spherical coordinates? Yes

Two ways, both with $\vec{x}'=0$ (do the easier problem first!)

i) $\nabla^2 \frac{1}{\sqrt{r^2 + \epsilon^2}}$ using Jackson cover $\nabla^2 \frac{1}{r} = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} = \frac{1}{r} \frac{\partial^2}{\partial r^2} (r^2)$
if angular derivatives are zero,

$$\Rightarrow \nabla^2 (r^2 + \epsilon^2)^{-1/2} = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \left(-\frac{1}{2} (r^2 + \epsilon^2)^{-3/2} \cdot 2r \right) = -\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 (r^2 + \epsilon^2)^{-3/2})$$

$$= -\frac{1}{r^2} \left(\frac{3r^3}{(r^2 + \epsilon^2)^{3/2}} - \frac{3}{2} \frac{r^3}{(r^2 + \epsilon^2)^{5/2}} \cdot 2r \right) = \frac{-3}{(r^2 + \epsilon^2)^{3/2}} + \frac{3r^2}{(r^2 + \epsilon^2)^{5/2}}$$

$$= \frac{-3}{(r^2 + \epsilon^2)^{5/2}} \left((r^2 + \epsilon^2) - r^2 \right) = \frac{-3\epsilon^2}{(r^2 + \epsilon^2)^{5/2}} \rightarrow \begin{cases} 0 & \text{if } r > 0 \\ \frac{1}{\epsilon^3} \rightarrow \infty & \text{if } r = 0 \end{cases}$$

Check $\int d^3x \frac{-3\epsilon^2}{(r^2 + \epsilon^2)^{5/2}} = -3 \cdot 4\pi \epsilon^2 \int_0^\infty r^2 dr \frac{1}{(r^2 + \epsilon^2)^{5/2}} = -4\pi$ as before

$$\Rightarrow \boxed{\nabla^2 \frac{1}{r} = -4\pi \delta^3(\vec{x})}$$

ii) $\nabla^2 \frac{1}{r + \epsilon} = \frac{1}{r} \frac{\partial^2}{\partial r^2} (r + \epsilon) = -\left(\frac{2r}{(r + \epsilon)^2} - \frac{2r^2}{(r + \epsilon)^3} \right) = -2r \left(\frac{r + \epsilon}{(r + \epsilon)^2} - \frac{r}{(r + \epsilon)^2} \right) = \frac{-2\epsilon r}{(r + \epsilon)^3}$

Check $\int d^3x \frac{-2\epsilon r}{(r + \epsilon)^3} = -2\epsilon \cdot 4\pi \int_0^\infty r^2 dr \frac{1}{(r + \epsilon)^3} = -2\epsilon \cdot 4\pi \int_0^\infty \frac{r dr}{(r + \epsilon)^3} = -4\pi \checkmark$

Does it work to use the divergence theorem?

$\int d^3x \nabla^2 \frac{1}{|\vec{x}|} = \int d^3x \vec{\nabla} \cdot \left(\vec{\nabla} \frac{1}{|\vec{x}|} \right)$ integrate over a sphere of radius R centered at

$= \oint_{\text{surface of sphere}} \left(\vec{\nabla} \frac{1}{r} \right) \cdot \hat{n} da = \int R^2 d\Omega \left(-\frac{1}{R^2} \right) = -4\pi \checkmark$ origin

[Jackson: $\nabla^2 = \hat{r} \frac{\partial^2}{\partial r^2} + \hat{\theta} \frac{1}{r} \frac{\partial^2}{\partial \theta^2} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial^2}{\partial \phi^2}$]

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So back to Poisson's equation:

$$\begin{aligned}\nabla^2 \phi(\vec{x}) &= \frac{1}{4\pi\epsilon_0} \int d^3x' \nabla^2 \frac{1}{|\vec{x}-\vec{x}'|} \rho(\vec{x}') \\ &= -\frac{1}{\epsilon_0} \int d^3x' \delta^3(\vec{x}-\vec{x}') \rho(\vec{x}') = -\frac{1}{\epsilon_0} \rho(\vec{x}) \quad \checkmark\end{aligned}$$

So Poisson's equation works!

What about delta functions in Gauss's Law?

$$\vec{\nabla}_x \cdot \vec{E}(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' \vec{\nabla}_x \cdot \left(\frac{\vec{x}-\vec{x}'}{|\vec{x}-\vec{x}'|^3} \right) \rho(\vec{x}') = ?$$

• Specialize again to $x'=0$ and use regulator again

$$\begin{aligned}\vec{\nabla} \cdot \frac{\vec{x}}{|\vec{x}|^3} &= \frac{\partial}{\partial x_i} \frac{x_i}{(x_1^2+x_2^2+x_3^2+\epsilon^2)^{3/2}} \stackrel{\partial_i x_j = \delta_{ij}}{=} \frac{3}{(x_1^2+x_2^2+x_3^2+\epsilon^2)^{3/2}} - \frac{3}{2} \frac{\partial(x_1^2+x_2^2+x_3^2)}{\partial(x_1^2+x_2^2+x_3^2+\epsilon^2)^{5/2}} \\ &= \frac{3\epsilon^2}{(x^2+\epsilon^2)^{5/2}} \xrightarrow{\epsilon \rightarrow 0} \begin{cases} 0 & \text{if } |\vec{x}| \neq 0 \\ -\frac{3}{\epsilon^3} \rightarrow \infty & \text{if } |\vec{x}| = 0 \end{cases}\end{aligned}$$

Again,

$$\int d^3x \vec{\nabla} \cdot \left(\frac{\vec{x}}{|\vec{x}|^3} \right) = 4\pi \int_0^\infty dr r^2 \frac{3\epsilon^2}{(r^2+\epsilon^2)^{5/2}} = 4\pi \Rightarrow \vec{\nabla} \cdot \left(\frac{\vec{x}}{|\vec{x}|^3} \right) = 4\pi \delta^3(\vec{x})$$

$$\Rightarrow \vec{\nabla} \cdot \left(\frac{\vec{x}-\vec{x}'}{|\vec{x}-\vec{x}'|^3} \right) = 4\pi \delta^3(\vec{x}-\vec{x}')$$

$$\Rightarrow \vec{\nabla}_x \cdot \vec{E}(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' 4\pi \delta^3(\vec{x}-\vec{x}') \rho(\vec{x}') = \frac{\rho(\vec{x})}{\epsilon_0} \quad \checkmark$$

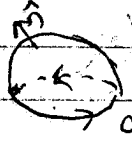
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Return to $\vec{\nabla} \times \vec{E} = 0 \dots$

We have another vector calculus theorem: Stoke's theorem, that says if we have an open surface S and C is the contour bounding it with line element $d\vec{l}$, then

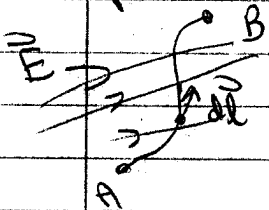
$$\int_S (\vec{\nabla} \times \vec{A}) \cdot \hat{n} da = \oint_C \vec{A} \cdot d\vec{l}$$

\hat{n} is defined by right-hand-rule given the direction of the line integral around C . eg.,  (imagine a hemisphere)

$$[\text{Also: } \int_S (\nabla \times \vec{4}) da = \oint_C \vec{4} d\vec{l}]$$

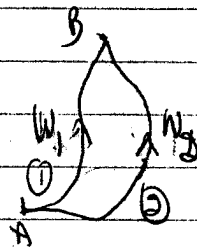
So if $\vec{\nabla} \times \vec{E}$ is zero, then the right side is zero for any closed contour $\Rightarrow \oint_C \vec{E} \cdot d\vec{l} = 0$.

But what is $\vec{E} \cdot d\vec{l}$ physically? $F = q\vec{E}$, so if we have a path from A to B , then the work we do moving a point charge q from A to B is



$$W = - \int_A^B \vec{F} \cdot d\vec{l} = -q \int_A^B \vec{E} \cdot d\vec{l}$$

But since we can close the path a different way



$$W_1 - W_2 = -q \int_A^B \vec{E} \cdot d\vec{l} + q \int_A^B \vec{E} \cdot d\vec{l} = q \oint_C \vec{E} \cdot d\vec{l} = 0$$

Stoke's theorem $= \int_S (\vec{\nabla} \times \vec{E}) \cdot \hat{n} da = 0 \Rightarrow W_1 = W_2$ for any 2 contours.

Check $W = -q \int_A^B \vec{E} \cdot d\vec{l} = q \int_A^B \vec{\nabla} \phi \cdot d\vec{l} = q \int_A^B d\phi = q(\phi_B - \phi_A) \Rightarrow$ Work is independent of path! (for electrostatics)

only depends on end points

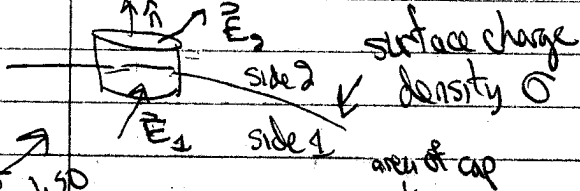
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Application of Gauss's Law: Discontinuities of \vec{E} at a surface.

Given $\nabla \cdot \vec{E} = \frac{1}{\epsilon_0} \rho$ and $\nabla \times \vec{E} = 0$

① Considers



Gauss's law applied to a "Gaussian cylinder" with sides flat go to zero \Rightarrow only surviving contributions from end caps!

make cylinder small enough so \vec{E}_2 & \vec{E}_1 can be taken as constants

$(\vec{E}_2 \cdot \hat{n} - \vec{E}_1 \cdot \hat{n}) S = \frac{1}{\epsilon_0} (\sigma S)$

$\Rightarrow (\vec{E}_2 - \vec{E}_1) \cdot \hat{n} = \frac{\sigma}{\epsilon_0}$

\Rightarrow normal component of \vec{E} has discontinuity given by surface charge.

②



$\nabla \times \vec{E} = 0 \Rightarrow \oint_C d\vec{l} \cdot \vec{E} = 0$

loop with length dl small enough so \vec{E}_2 & \vec{E}_1 constant

$\Rightarrow \vec{E}_2 \cdot d\vec{l} - \vec{E}_1 \cdot d\vec{l} = 0$

sides vanishes in limit

$\Rightarrow E_{2t} = E_{1t}$

\Rightarrow tangential component is always continuous (even if $\sigma \neq 0$)