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7701 Lecture 27

- Emphasis points for 2nd midterm
 - go over problem sets #5-#8 together with lecture notes
 - Formula sheet plus Jackson covers are available
 - ⇒ extra copies at the exam, but you can bring your own.
- Focus on the core competencies for each topic
 - a. delta and Heaviside functions
 - b. Fourier transforms
 - c. Vector calculus theorems and div, grad, curl, ∇^2
 - d. Electrostatics Gauss's Law, scalar potential
 - e. Poisson/Laplace equations, capacitance, energy

} look through emphasis points in lecture notes

a. delta and Heaviside functions

- Spot the Error! problems
- change (or other) densities in other coordinates
- derivatives of θ -functions
- representations as Fourier transforms
- delta functions of functions:

$$I = \int_{-\infty}^{\infty} e^{-x^2} \delta(x^2 - 6) dx = \int_{-\infty}^{\infty} e^{-x^2} \delta(x-2) \delta(x+3) dx = \frac{1}{5} (e^{-4} + e^{-9})$$

$x_0=2$ $x_1=3$

b. Fourier transforms

- carrying out a transform and inverse (often contour integration)
- Heaviside functions from closing in one half plane or the other and getting different results (often zero)
- solving $\ddot{x} + 2\alpha\dot{x} + \omega_0^2 x = f(t)$ by applying $\int_{-\infty}^{\infty} dt e^{i\omega t}$ to both sides, finding the resulting $\hat{x}(\omega)$, then inverting to find $x(t)$.
- focus on setting up integrals mostly, but some features in carrying them out, like θ functions
- 3-dimensional transform in spherical coordinates
 - eg. $f(\vec{x}) = f_0 \frac{e^{-r/a}}{4\pi r} \Rightarrow f(\vec{k}) = \frac{1}{(2\pi)^3} \int d^3r f(\vec{r}) e^{-i\vec{k}\cdot\vec{r}} = \frac{1}{(2\pi)^3} \int_0^{\infty} dr r^2 \int d\Omega e^{-i\vec{k}\cdot\vec{r}} f_0 \frac{e^{-r/a}}{4\pi r}$
 - choosing \hat{z} axis along \vec{k}

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c. Vector calculus theorems and div, grad, curl, ∇^2

- carrying out operators using Jackson covers (A_1, A_2, A_3)
- applying Stokes and Green's and divergence theorems

d. Electrostatics I

- Coulomb potential in vector form
- Gauss's law in symmetric cases; \vec{E} -field BC's
- scalar potential $\vec{E} = -\vec{\nabla}\Phi$, point charge and $\rho(\vec{x})$ formulas

e. Electrostatics II

- $\vec{\nabla} \cdot \vec{E} = -\nabla^2 \Phi = \frac{1}{\epsilon_0} \rho(\vec{x})$, $\vec{\nabla} \times \vec{E} = 0$, $\nabla^2 \frac{1}{|\vec{x}-\vec{x}'|} = -4\pi \delta(\vec{x}-\vec{x}')$
- electrostatic energy in terms of $\rho, \Phi, |\vec{E}(\vec{x})|^2$
- capacitance

• Recap of Neuman BC's and master formula: (183)

- Dirichlet Green function for region outside of sphere
 - use method of images (184)-(185)
 - also force and surface charge density

• Return to general discussion of solving for Green function (176)-(179)

- Aside: nonlinear equation (180)
- conducting sphere in uniform field \vec{E}_0

• Expansions \Rightarrow separation of variables (191)+

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Let's apply our knowledge of Fourier series (and transforms) to find solutions in Cartesian coordinates, to Poisson/Laplace equations.

- use separation of variables
- apply to problems with a rectangular box

• Laplace's equation:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = 0$$

⇒ look for $\phi(x, y, z) = X(x)Y(y)Z(z)$

$$\Rightarrow X''YZ + XY''Z + XYZ'' = 0$$

Isolate dependence on x or y or z by dividing by XYZ

$$\Rightarrow \frac{X''}{X}(x) + \frac{Y''}{Y}(y) + \frac{Z''}{Z}(z) = 0$$

\uparrow only on x \uparrow only on y \uparrow only on z

As we vary x, y, z , each term must stay the same, so each is a constant. Positive or negative constants? Can't all be the same sign!

- This will lead to multiple representations of ϕ or G_0 .
- Here choose x, y separation constants negative and z constant positive:

$$\frac{X''}{X} = -\alpha^2 \quad \Rightarrow \quad X(x) = C_1 e^{i\alpha x} + C_2 e^{-i\alpha x} \quad (\text{or sin's and cos's})$$

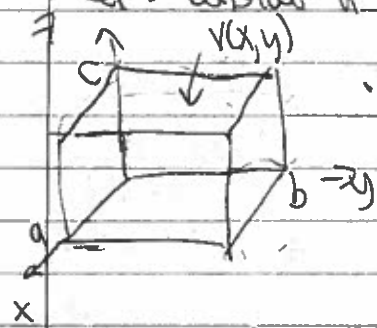
$$\frac{Y''}{Y} = -\beta^2 \quad \Rightarrow \quad Y(y) = \tilde{C}_1 e^{i\beta y} + \tilde{C}_2 e^{-i\beta y} \quad (\text{or sin's and cos's})$$

$$\frac{Z''}{Z} = \alpha^2 + \beta^2 \equiv \gamma^2 \quad \Rightarrow \quad Z(z) = \hat{C}_1 e^{\gamma z} + \hat{C}_2 e^{-\gamma z} \quad (\text{or sinh's and cosh's})$$

- At this stage a general solution to Laplace's equation in a rectangular system.

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Let's consider a definite example to see how it plays out.



- Rectangular box $a \times b \times c$
- All surfaces grounded except top, which has $V(x,y)$
- Note: $V(0,y) = V(a,y) = V(x,0) = V(x,b) = 0$

Plan! Use the boundary conditions at $x=0, y=0, z=0$ to determine the correct combination of exponentials, then $X(a) = Y(b) = 0$ to constrain α and β . Finally, fix the expansion coefficients by requiring $V(x,y)$ at $z=c$.

$$X(0) = 0 \Rightarrow X(x) \propto \sin \alpha x; \quad X(a) = 0 \Rightarrow \sin(\alpha a) = 0 \Rightarrow \alpha_n = \frac{n\pi}{a}$$

$$Y(0) = 0 \Rightarrow Y(y) \propto \sin \beta y; \quad Y(b) = 0 \Rightarrow \sin(\beta b) = 0 \Rightarrow \beta_m = \frac{m\pi}{b}$$

$$Z(0) = 0 \Rightarrow Z(z) \propto \sinh(\gamma z) \text{ and } \gamma_{nm} = \sqrt{\alpha_n^2 + \beta_m^2} = \sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{b}\right)^2} \equiv \gamma_{nm}$$

So ϕ is a combination of $\phi_{nm}(x,y,z) \propto \sin(\alpha_n x) \sin(\beta_m y) \sinh(\gamma_{nm} z)$, so write a general linear combination:

$$\phi(x,y,z) = \sum_{n,m \in \mathbb{N}} A_{nm} \sin(\alpha_n x) \sin(\beta_m y) \sinh(\gamma_{nm} z)$$

A_{nm} is determined by requiring $\phi(x,y,z=c) = V(x,y)$ for every x,y . This is a double Fourier series \Rightarrow find A_{nm} by projecting coefficients:

$$V(x,y) = \sum_{n,m=1}^{\infty} [A_{nm} \sinh(\gamma_{nm} c)] \sin(\alpha_n x) \sin(\beta_m y)$$

Recall solving wave equation for a string with fixed ends using Fourier series

$$f(x) = \sum_{n=1}^{\infty} a_n \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \text{ with } f(0) = f(L) = 0 \Rightarrow a_n = \left[\frac{2}{L}\right] \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx \text{ using } \int_0^L \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi x}{L}\right) dx = \delta_{nm}$$

$$\Rightarrow A_{nm} \sinh(\gamma_{nm} c) = \left(\frac{2}{a}\right) \left(\frac{2}{b}\right) \int_0^a \int_0^b V(x,y) \sin(\alpha_n x) \sin(\beta_m y) dx dy$$

$$\Rightarrow A_{nm} = \frac{4}{ab \sinh \gamma_{nm} c} \int_0^a \int_0^b V(x,y) \sin(\alpha_n x) \sin(\beta_m y) dx dy \quad \text{Problem solved!}$$

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The sines here are one example of an orthonormal set of functions. We will discuss these more generally later as solutions to a Sturm-Liouville equation:

$$\hat{L}_x u_n(x) + \lambda w(x) u_n(x) = 0 \quad n=1, 2, \dots, \infty$$

$\hat{L}_x \equiv \frac{d}{dx} f(x) \frac{d}{dx} - g(x)$ weight function
 eigenvalue
 normalization constant chosen to be 1 \Rightarrow orthonormal

Properties:

(1) orthogonal w.r.t $w(x)$: $\int_a^b w(x) u_m^*(x) u_n(x) dx = \delta_{mn}$

(2) completeness on $[a, b]$: $f(x) = \sum_{n=1}^{\infty} a_n u_n(x) \Rightarrow a_n = \int_a^b f(x) u_n^*(x) w(x) dx$
 and $\delta(x-x') = \sum_n u_n^*(x') u_n(x) w(x)$

(3) real eigenvalues

(4) self-adjoint $\int_a^b u_n^* (\hat{L}_x u_m) dx = \int_a^b (\hat{L}_x u_n)^* u_m dx$ [from boundary conditions on the u_n 's]

For $\hat{L}_x = \frac{d^2}{dx^2}$ [so $f(x)=1, g(x)=0$] and $w(x)=1$ with Dirichlet BC's at $x=0$ and L , we have

the case just considered:

$$u_n(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L} \Rightarrow f(x) = \sum_{n=1}^{\infty} a_n \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L} \text{ with } a_n = \sqrt{\frac{2}{L}} \int_0^L \sin \frac{n\pi x}{L} f(x) dx$$

As we already know, the $\{\sin \frac{n\pi x}{L}\}$ are a complete set on $x \in [0, L]$ with these boundary conditions. (With Neumann or mixed BC's, we have to consider cosines as well.)

From (2), we have $\delta(x-x') = \sum_{n=1}^{\infty} \sqrt{\frac{2}{L}} \sin \left(\frac{n\pi x}{L} \right) \sin \left(\frac{n\pi x'}{L} \right)$

We'll use this next time to find the Dirichlet Green function $G_D(\vec{x}, \vec{x}')$ for the rectangular box.