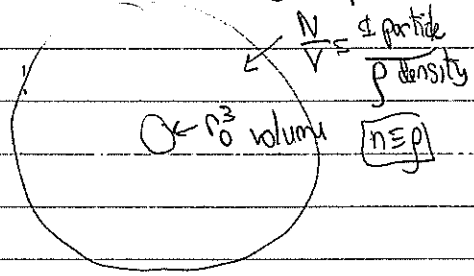


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Comments on Piotr Magerski\_path-integral-gmc-rtar.pdf

slide 1: always attraction, but different regimes based on value of  $k_F a_0 \Rightarrow$  our expansion parameter, ( $a_0 \equiv a_s$ )  
note that  $1/k_F a_0$  is plotted, so the origin is very large scattering length: just bound or just unbound state at zero energy.

slide 2: Joaquim Druot is Casey's supervisor for REU/thesis project.  
former OSU postdoc in NTG, now faculty at UNC

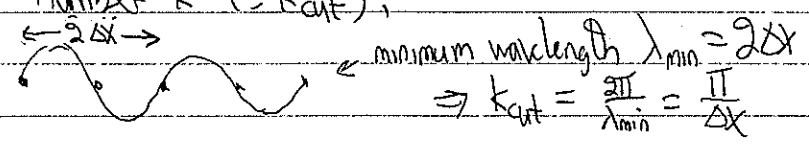
slide 4:   $\Rightarrow n \rho^3 \ll 1$  means the ratio of the volume over which a particle interacts is  $\ll$  the volume for that particle,  $\Rightarrow$  dilute gas!

universality follows from lack of scale:  $r_0 \rightarrow 0, a_0 \rightarrow \pm \infty$   
so dimensional analysis implies  $E \propto E_{\text{fermi gas}} \Rightarrow$  dimensionless ratio  
and for  $x = T/E_F$  scaling parameter (dimensionless)  
 $\uparrow$  fermi energy  $\Rightarrow E(x) = f(x) E_{FG}$

slide 8: quantities of interest: energy, pressure, momentum distribution  $n_s(k)$   
talk of  $n_s(k) \sim C/k^4 \Rightarrow C$  is called the "contact"

slide 9: Same Hamiltonian we've considered  $g \rightarrow -C_0$  (so  $g > 0$  is attractive)

- $\nu = 2$ : only spin up and down.
- $\Delta \equiv \nabla^2$
- Given a spacing  $\Delta x$  and periodic boundary conditions, what is the minimum wavelength  $\Rightarrow$  maximum momentum wave number  $k$  ( $\equiv k_{\text{cut}}$ )?



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slide 10 |  $k_{cut} \equiv \Lambda_{UV} = \frac{\pi}{\Delta x}$

What about  $k_{min} \equiv \Lambda_{IR}$ ?  $e^{ik_n x} = e^{ik_n(x+L)}$  PBC  
 $\Rightarrow k_n L = 2\pi n \Rightarrow k_{min} = \frac{2\pi}{L} = \Lambda_{IR}$

How big and fine a lattice do we need?

Want the minimum kinetic energy and spacing to be much smaller than the Fermi energy

$$\Rightarrow \frac{\Lambda_{IR}^2}{2m} \ll E_F$$

Want the energy cutoff to be much larger than the pairing gap (?)

$$\frac{\Lambda_{UV}^2}{2m} \gg \Delta$$

• real space  $\xleftrightarrow{FFT}$  momentum space means use "Fast Fourier Transform" algorithm to switch back and forth. Kinetic energy is trivial in momentum space, potential is trivial in coordinate (real) space. [e.g. trivial  $\Rightarrow$  diagonal matrix]

### Adell: Running coupling

• See equation (29) from section 7 b on "Non-perturbative matching"

$$C_0(\Lambda_c) = \frac{4\pi}{m} \frac{1}{a_0 - \frac{2}{\pi}\Lambda_c} \Rightarrow \frac{1}{C_0} = \frac{m}{4\pi} \left( \frac{1}{a_0} - \frac{2}{\pi}\Lambda_c \right)$$

$\uparrow$   
sharp cutoff

$$= \frac{m}{4\pi a_0} - \frac{m}{2\pi^2} \Lambda_c$$

$$= -\frac{1}{g} \left( = \frac{m}{4\pi^2 a} - \frac{mk_{cut}}{2\pi^2 \hbar^2} \right) \checkmark$$

The unitary limit is when  $C_0 \rightarrow \infty$  so

$$\boxed{\frac{1}{g} = \frac{mk_{cut}}{2\pi^2} = \frac{m \pi / \Delta x}{2\pi^2} = \frac{m}{2\pi \Delta x}}$$

Trotter expansion: breaks up partition function into the product of many small imaginary time / temperature steps.

(3)

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slide 12: Splitting up the kinetic and potential parts costs only a  $\mathcal{P}^3$  error (you might have thought it was worse  $\Rightarrow$  see Piazza).

• Eliminate fermion integration in favor of auxiliary field  $\sigma$ .

\* for basic idea, see (5) on 11/3/14