



Smeared versus localized sources in flux compactifications

Smeared vs.
localized sources

Timm Wrase

Flux compact.

BPS case

non-BPS case

Localiz. effects

Conclusion

Timm Wrase

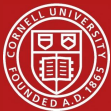


Cornell University

Based on:

TW, Zagermann 1003.0029
Blåbäck, Danielsson, Junghans, Van Riet, TW, Zagermann 1009.1877

String Vacuum Project meeting Fall 2010



Classical type II flux compactifications

- Most constructions of dS vacua use non-perturbative effects for moduli stabilization
- dS after uplift which breaks explicitly SUSY

KKLT, LARGE VOLUME

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KKLT, LARGE VOLUME

- It is in principle possible to stabilize all moduli classically

VILLADORO, ZWIRNER HEP-TH/0503169

DEWOLFE, GIRYAVETS, KACHRU, TAYLOR HEP-TH/0505160

CÁMARA, FONT, IBÁÑEZ HEP-TH/0506066

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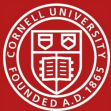
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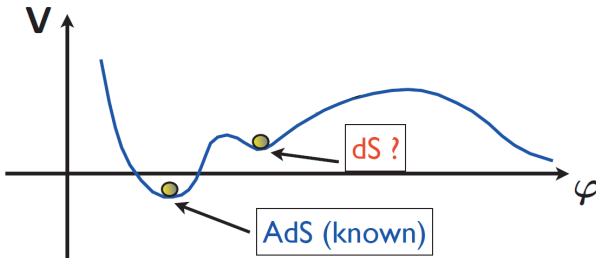
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Classical type II flux compactifications

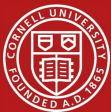
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Can we find classical dS vacua?

HERTZBERG, TEGMARK, KACHRU, SHELTON, OZCAN 0709.0002 [ASTRO-PH]



Type II supergravity

The classical ingredients for type II supergravity theories are

RR-fluxes F_p , NSNS H -flux, R_6 , Oq -planes, ...

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For smeared Oq -planes we find a 4D scalar potential

$$V(\rho, \phi, \dots) = \sum_p V_{F_p} + V_H + V_{R_6} - V_{Oq},$$

where $\rho = (vol_6)^{1/3}$ and ϕ is the dilaton.



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where $\rho = (vol_6)^{1/3}$ and ϕ is the dilaton.

When is $\partial_\rho V = \partial_\phi V = 0$ and $V > 0$ possible?

HERTZBERG, KACHRU, TAYLOR, TEGMARK 0711.2512 [HEP-TH]



We can evade a no-go theorem involving ρ and ϕ with the following minimal ingredients

Curvature	IIA	IIB
$V_{R_6} \sim -R_6 \leq 0$	O4, H , F_0	O3, H , F_1
$V_{R_6} \sim -R_6 > 0$	O4, F_0 O4, F_2 O6, F_0	O3, F_1 O3, F_3 O3, F_5 O5, F_1

HERTZBERG, KACHRU, TAYLOR, TEGMARK 0711.2512 [HEP-TH]

SILVERSTEIN 0712.1196 [HEP-TH]

HAQUE, SHIU, UNDERWOOD, VAN RIET 0810.5328 [HEP-TH]

CAVIEZEL, KOERBER, KÖRS, LÜST, TW, M. ZAGERMANN 0812.3551 [HEP-TH]

FLAUGER, ROBBINS, PABAN, TW 0812.3886 [HEP-TH]

DANIELSSON, HAQUE, SHIU, VAN RIET 0907.2041 [HEP-TH]

DE CARLOS, GUARINO, MORENO 0907.5580, 0911.2876 [HEP-TH]

CAVIEZEL, TW, ZAGERMANN 0912.3287 [HEP-TH]

TW, ZAGERMANN 1003.0029 [HEP-TH]

DANIELSSON, KOERBER, VAN RIET 1003.3590 [HEP-TH]

DANIELSSON, HAQUE, KOERBER, SHIU, VAN RIET, TW 1011.xxxx [HEP-TH]



Smearred versus localized sources

- O-planes are localized objects
- Smearing was necessary to solve equations of motion

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When is smearing $\delta(Oq) \approx 1$ a valid approximation?

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Smearred versus localized sources

- O-planes are localized objects
- Smearing was necessary to solve equations of motion

When is smearing $\delta(Oq) \approx 1$ a valid approximation?

Negative curvature $R_6 < 0$ requires (in the localized case)

large warping or large stringy corrections

DOUGLAS, KALLOSH 1001.4008 [HEP-TH]



An example with BPS sources

Giddings, Kachru and Polchinski found localized no-scale
Minkowski solutions with O3-planes

GIDDINGS, KACHRU, POLCHINSKI HEP-TH/0105097

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GIDDINGS, KACHRU, POLCHINSKI HEP-TH/0105097

smeared case

$H, F_3, O3$

$$ds^2 = ds_4^2 + ds_6^2$$

$$0 = dF_5 = H \wedge F_3 - \tilde{\mu}_3$$



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localized case

$H, F_3, O3, F_5, A$

$$ds^2 = e^{2A} ds_4^2 + e^{-2A} ds_6^2$$

$$dF_5 = H \wedge F_3 - \tilde{\mu}_3 \delta(O3)$$



An example with BPS sources

- Can solve the 10D equations of motions in both cases
- Find no-scale Minkowski vacua
- Internal space is (conformally) Ricci-flat

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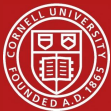


An example with BPS sources

- Can solve the 10D equations of motions in both cases
- Find no-scale Minkowski vacua
- Internal space is (conformally) Ricci-flat

But localization effects are large

$$\nabla^2 e^{-4A} = -e^{-\phi} |H|^2 + \tilde{\mu}_3 \delta(O3)$$



An example with BPS sources

- Can solve the 10D equations of motions in both cases
- Find no-scale Minkowski vacua
- Internal space is (conformally) Ricci-flat
- Complex structure moduli and ϕ are stabilized

smeared case

$$F_3 = -e^{-\phi} \star_6 H$$

localized case

$$F_3 = -e^{-\phi} \star_6 H$$



An example with BPS sources

BUT

$$F_3 = -e^{-\phi} \star_6 H = -e^{-\phi} \star_6 H$$

since warp factor cancels: $\star_6 H \approx \sqrt{\det(e^{2A} g_6)} (e^{-2A} g_6^{-1})^3 H$

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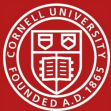
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Moduli values at minimum unchanged!

Approximation $\delta(O3) \approx 1$ is “ok”

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Moduli values at minimum unchanged!

Approximation $\delta(O3) \approx 1$ is “ok”

smeared: H and F_3 stabilize moduli

localized: $\tilde{\mu}_3 \delta(O3), F_5, A$ give corrections of equal size

\Rightarrow corrections from $\tilde{\mu}_3 \delta(O3), F_5, A$ cancel each other



A T-dual example with BPS sources

T-duality along one H -flux direction \leftrightarrow DOUGLAS, KALLOSH

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$$H \rightarrow R_6 < 0$$

$$F_3 \rightarrow F_4$$

$$O3 \rightarrow O4$$

$$F_5 \rightarrow F_4$$

$$A \rightarrow A$$



A T-dual example with BPS sources

T-duality along one H -flux direction \leftrightarrow DOUGLAS, KALLOSH

$$H \rightarrow R_6 < 0$$

$$F_3 \rightarrow F_4$$

$$O3 \rightarrow O4$$

$$F_5 \rightarrow F_4$$

$$A \rightarrow A$$

Conclusions remain unchanged:

- DOUGLAS, KALLOSH \Rightarrow warping effects are large
- But again smeared moduli values are unaffected

Note: $\int \sqrt{g_{10}} R_6 < 0 \Rightarrow V_{R_6} > 0$

(no ‘uplift’ to dS, solutions are Minkowski)

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An example with non-BPS sources

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Conclusion

Replace $O3$ -plane by $\overline{D3}$ -brane:

smeared case

$$H, F_3, \overline{D3}$$

$$ds^2 = ds_4^2 + ds_6^2$$

$$0 = dF_5 = H \wedge F_3 - \mu_3$$

localized case

$$H, F_3, \overline{D3}, F_5, A, \dots$$

$$ds^2 = e^{2A} ds_4^2 + ds_6^2$$

$$dF_5 = H \wedge F_3 - \mu_3 \delta(O3)$$



An example with non-BPS sources

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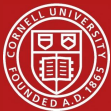
smeared case

$$H, F_3, \overline{D3}$$

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- Can solve the 10D equations of motions



An example with non-BPS sources

smeared case

$$H, F_3, \overline{D3}$$

$$ds^2 = ds_4^2 + ds_6^2$$

$$0 = dF_5 = H \wedge F_3 - \mu_3$$

- Can solve the 10D equations of motions
- Find AdS solutions $V = V_{F_3} + V_H - V_{R_6} + V_{\overline{D3}} < 0$
- Internal space is positively curved: e.g. $S^3 \times S^3$



An example with non-BPS sources

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- Internal space is positively curved: e.g. $S^3 \times S^3$
- Complex structure and ϕ stabilized ($F_3 = -e^{-\phi} \star_6 H$)
- volume moduli stabilized: e.g. $R_{ij}^{S^3} = \frac{1}{2}e^{-\phi}|H|^2 g_{ij}^{S^3}$



An example with non-BPS sources

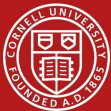
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- Can solve the 10D equations of motions
- Find AdS solutions $V = V_{F_3} + V_H - V_{R_6} + V_{\overline{D3}} < 0$
- Internal space is positively curved: e.g. $S^3 \times S^3$
- Complex structure and ϕ stabilized ($F_3 = -e^{-\phi} \star_6 H$)
- volume moduli stabilized: e.g. $R_{ij}^{S^3} = \frac{1}{2}e^{-\phi}|H|^2 g_{ij}^{S^3}$
- no SUSY but volume and dilaton masses above BF bound



An example with non-BPS sources

localized case

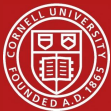
$$H, F_3, \overline{D3}, F_5, A, \dots$$

$$ds^2 = e^{2A} ds_4^2 + ds_6^2$$

$$dF_5 = H \wedge F_3 - \mu_3 \delta(O3)$$

- (Assume) $F_1 = 0$, $F_3 = -e^{-\phi} \star_6 H$ for arbitrary g_6
- Combine eoms for F_3 , H , F_5 and external Einstein:

$$e^{-2A} R_4 = -(1+1) \mu_3 \delta(\overline{D3})$$



An example with non-BPS sources

localized case

$$H, F_3, \overline{D3}, F_5, A, \dots$$

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The smeared solution cannot be localized?



BPS versus non-BPS sources

BPS condition in GKP

$$\frac{1}{4}(T_m^m - T_\mu^\mu)^{\text{loc}} \geq \mu_3 \rho_3^{\text{loc}}$$

For $O3$, $D3$ and $\overline{D3}$ we have $T_m^m = 0$.

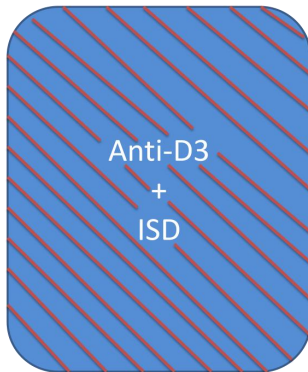
	$O3$	$\overline{D3}$	$D3$
ρ_3^{loc}	$-\frac{1}{4}$	-1	1
$-\frac{1}{4}T_\mu^\mu$	$\mu_3 \rho_3^{\text{loc}}$	$-\mu_3 \rho_3^{\text{loc}}$	$\mu_3 \rho_3^{\text{loc}}$
BPS	✓	✗	✓

Force between $\overline{D3}$ and fluxes H , F_3

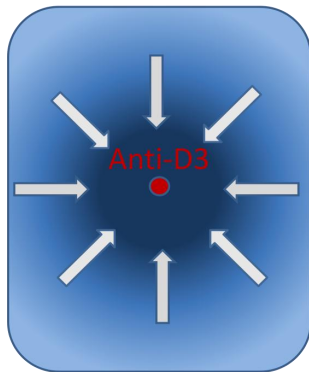
\Rightarrow no static localized solution



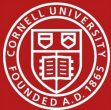
non-BPS sources



smeared case



localized case



Localization effects

Localization effects are generically large
in flux compactifications

$$\nabla^2 e^A = |\text{flux}|_p^2 - \delta(\text{source})$$

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Warping suppressed in the large volume limit $g_6 \rightarrow \lambda^2 g_6$?

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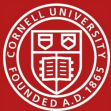
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Warping suppressed in the large volume limit $g_6 \rightarrow \lambda^2 g_6$?

Yes!

$$\frac{1}{\lambda^2} \nabla^2 e^A = \frac{1}{\lambda^{2p}} |\text{flux}|_p^2$$



Localization effects

Localization effects are generically large
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$$\nabla^2 e^A = |\text{flux}|_p^2 - \delta(\text{source})$$

Warping suppressed in the large volume limit $g_6 \rightarrow \lambda^2 g_6$?

Yes! But so are the fluxes!

$$\frac{1}{\lambda^2} \nabla^2 e^A = \frac{1}{\lambda^{2p}} |\text{flux}|_p^2$$



Localization effects

Localization effects are generically large
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$$\nabla^2 e^A = |\text{flux}|_p^2 - \delta(\text{source})$$

Are there regions of small warping $e^A \approx 1$ and $(\nabla A)^2 \ll 1$?

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$$\nabla^2 e^A = |\text{flux}|_p^2 - \delta(\text{source})$$

Are there regions of small warping $e^A \approx 1$ and $(\nabla A)^2 \ll 1$?

Not really:

$$\nabla^2 e^A = e^A \nabla^2 A + e^A (\nabla A)^2 \approx \nabla^2 A = |\text{flux}|_p^2$$



Conclusion:

- Explicit examples:
 - smeared BPS sources are ok
 - non-BPS solutions problematic
- Localization effects comparable to fluxes

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Outlook:

- Study solutions close to BPS point
- Construct localized, non-BPS examples
- Generalize findings to intersecting branes



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