

Three family models  
from the heterotic string

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## Outline

- 3 family orbifold GUT on  $\mathcal{M}_4 \times S_1 / (Z_2 \times Z'_2)$



- heterotic string compactified on  $[T_2]^3 / Z_6 + \text{Wilson lines}$
- 

- $E(6)$  example

Orbifold breaking  $E(6)$  to PS

PS  $\rightarrow$  SM via Higgs

3rd family in bulk

1st & 2nd families on  $SO(10)$  fixed pts.

Gauge coupling unification & proton decay

Fermion masses

- Conclusions

$$\text{SUSY E(6) on } \mathcal{M}_4 \times S_1 / (Z_2 \times Z'_2)$$

$$M_c = (\pi R)^{-1} \ll M_*$$

- E(6) breaks to PS by orbifold parities

$$P \quad P'$$

$$\text{E(6)} \rightarrow \text{SO(10)} \rightarrow \text{SU(4)}_c \times \text{SU(2)}_L \times \text{SU(2)}_R [= \text{PS}]$$

- Gauge and hypermultiplet in bulk –

$$(V, \Sigma) [\mathbf{78}] + (\mathbf{27} \oplus \overline{\mathbf{27}})$$

Consider ( + + ) modes  $\implies$  3rd family and Higgs

$$f_3^c = (\overline{4}, 1, 2), \quad f_3 = (4, 2, 1), \quad h = (1, 2, 2)$$

$$V = \mathbf{78} \rightarrow \mathbf{45} \rightarrow \text{adjoint of PS}$$

$$\Sigma = \mathbf{78} \rightarrow (\mathbf{16} \oplus \overline{\mathbf{16}}) \rightarrow f_3^c + \overline{\chi}^c$$

$$\mathbf{27} \rightarrow \mathbf{16} \rightarrow f_3$$

$$\overline{\mathbf{27}} \rightarrow \mathbf{10} \rightarrow h$$

3rd family Yukawa unification

$$\lambda_t = \lambda_b = \lambda_\tau = g \equiv \sqrt{4\pi\alpha_G}$$

$$\int_0^{\pi R} dx_5 \, g_5 (\overline{\mathbf{27}} \Sigma \mathbf{27}) \rightarrow g \, h \, f_3^c \, f_3$$

$$g = g_5 \sqrt{M_c}$$

## PS breaks to SM

- PS breaking to SM by Higgs vevs

$$\chi^c = (\bar{4}, 1, 2), \bar{\chi}^c = (4, 1, 2)$$

$$\langle \chi^c \rangle = \langle \bar{\chi}^c \rangle = M_b$$

- Need additional bulk states – 3 (  $\mathbf{27} \oplus \overline{\mathbf{27}}$  )

$$3 ( \mathbf{27} \oplus \overline{\mathbf{27}} ) \rightarrow 2 ( \mathbf{16} ) \oplus \overline{\mathbf{16}} \rightarrow 2 ( \chi^c ) + \bar{\chi}^c + 3 C$$

Now total  $\implies 2(\chi^c + \bar{\chi}^c) + 3 C [= (6, 1, 1)]$

- Superpotential  $W$  gives mass to color triplets and breaks PS to SM along D & F flat directions

$$W = M_{eff} C_1 C_2 + (\chi^c \chi^c + \bar{\chi}^c \bar{\chi}^c) C_3$$

1st & 2nd family ?

5th Dimension

SO(10) brane  
brane

E(6) bulk

SU(6) × SU(2)<sub>R</sub>

[—————]

0

πR

- In bulk or on SU(6) × SU(2)<sub>R</sub> brane —  $M_c < M_G$  \*

On SU(6) × SU(2)<sub>R</sub> brane, one family  $\in [(15, 1) \oplus (\bar{6}, 2)]$

- On SO(10) brane —  $M_c \geq M_G$  \*

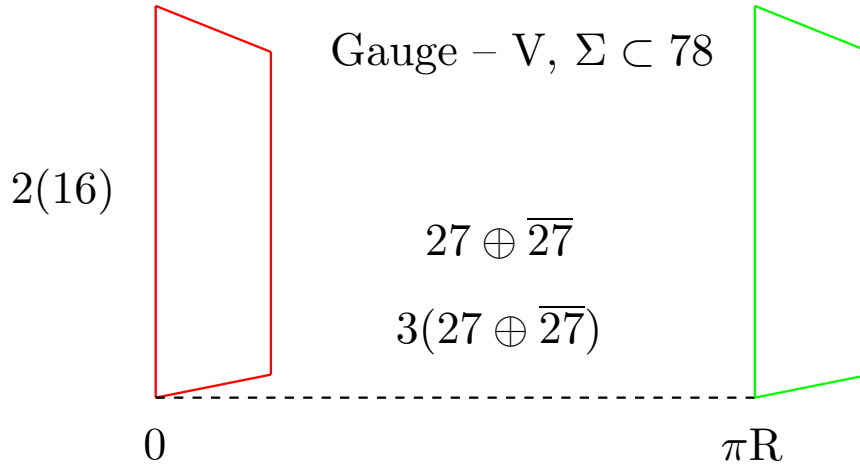
\* NO problem with proton decay !

◇ **In string theory, don't get to choose *easily***

Summary  
E(6) orbifold GUT

SO(10) brane

SU(6) × SU(2)<sub>R</sub> brane



- 3rd family and Higgs

$$\Sigma \rightarrow 16 \oplus \bar{16} \rightarrow f_3^c \oplus \bar{\chi}^c$$

$$27 \rightarrow 16 \rightarrow f_3$$

$$\bar{27} \rightarrow 10 \rightarrow h$$

- Yukawa unification

$$\int_0^{\pi R} dx_5 g_5 (\bar{27} \Sigma 27) \rightarrow g (h f_3^c f_3)$$

$$\implies \lambda_t = \lambda_b = \lambda_\tau = g = \sqrt{4\pi\alpha_G}$$

- PS breaking sector

$$3(27 \oplus \bar{27}) \rightarrow 2(16) \oplus \bar{16} \rightarrow$$

$$2(\bar{4}, 1, 2)[= 2 \chi^c] \oplus (4, 1, 2)[= \bar{\chi}^c] \oplus 3(6, 1, 1)[= 3 C]$$

## Heterotic string compactified on $[T_2]^3/Z_6$

$[T_2]^3 = G(2) \otimes SU(3) \otimes SO(4)$  root lattice

$$Z_6 = Z_2 \otimes Z_3$$

Orbifold defined by rotation on torus

$$\mathbf{Z}^i \rightarrow e^{2\pi i v_i} \mathbf{Z}^i \quad (i = 1, 2, 3)$$

$$v_3 = \frac{1}{3} (1, -1, 0), \quad v_2 = \frac{1}{2} (1, 0, -1)$$

Rotations embedded in  $E(8) \times E(8)$  root lattice via shift vectors + Wilson lines

Satisfy NON-TRIVIAL constraints of modular invariance !

Consider first  $[T_2]^3/Z_3 +$   
 $W_3$  in  $SU(3)$  torus  
 $SO(4)$  torus  $R^{-1} \gg l_s = M_*^{-1}$

$$v_3 = \frac{1}{3} (1, -1, 0)$$

$$V_3 = \frac{1}{3} (2, 2, 2, 0, 0, 0, 0, 0) \oplus (\dots)$$

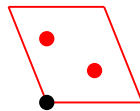
$$W_3 = \frac{1}{3} (1, -1, 0, 0, 0, 0, 0, 0) \oplus (\dots)$$

$\implies N = 2$  SUSY in 5D bulk

$G_2$

$SU_3$

$SO(4)$



$V, \Sigma \in E(6)$

$(\mathbf{27} \oplus \overline{\mathbf{27}})$

$3(\mathbf{27} \oplus \overline{\mathbf{27}})$

$G_2 \oplus SU_3 \oplus SO_4$  lattice with  $\mathbb{Z}_3$  fixed points. The fields  $V, \Sigma$ , and  $\mathbf{27}(\in U_1) + \overline{\mathbf{27}}(\in U_2)$  are bulk states from the untwisted sectors. On the other hand,  $3 \times (\mathbf{27} + \overline{\mathbf{27}})$  are “bulk” states located on the  $T_{(0,1)}/T_{(0,2)}$  twisted sector ( $G_2, SU_3$ ) fixed points.



Applying  $Z_2$  orbifold + Wilson line in  
5th Dimension

Breaks  $N = 2$  to  $N = 1$  SUSY

Defines parities  $P, P'$

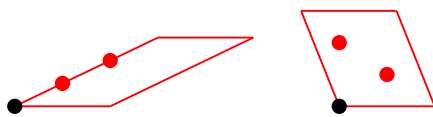
$$P \Leftrightarrow (v_6; V_6; W_3)$$

$$P' \Leftrightarrow (v_6; V_6 + W_2; W_3)$$

$$V_6 = V_2 - V_3, \quad W_2 = \frac{1}{2} (1, 0, 0, 0, 0, 1, 1, 1) \oplus (\dots)$$

Consider massless states, i.e.  $(+, +)$  modes of  $(P, P')$  from bulk and  $T_{(0,1)}/T_{(0,2)}$  twisted sector ( $G_2, SU(3)$ ) fixed points.

Find



$$V \in PS \quad (f_3^c + \bar{\chi}^c) \in \Sigma$$

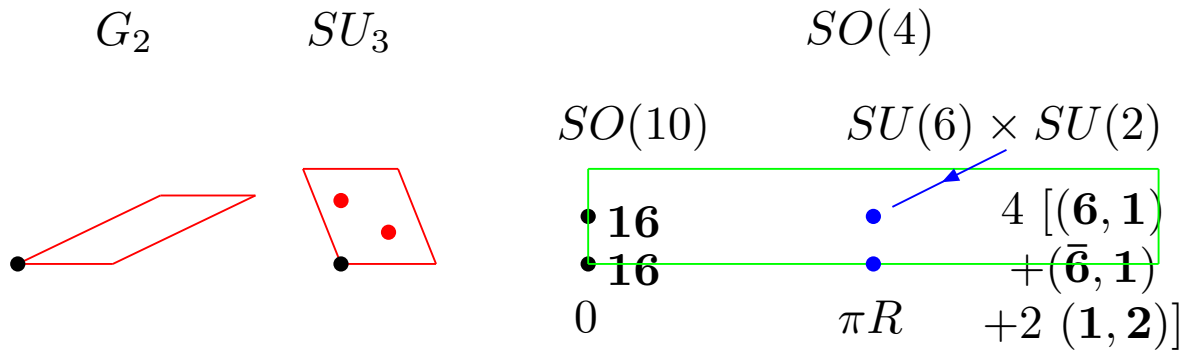
$$f_3 \in \mathbf{27} + h \in \overline{\mathbf{27}}$$

$$2(\chi^c) + \bar{\chi}^c + 3C \in 3(\mathbf{27} + \overline{\mathbf{27}})$$

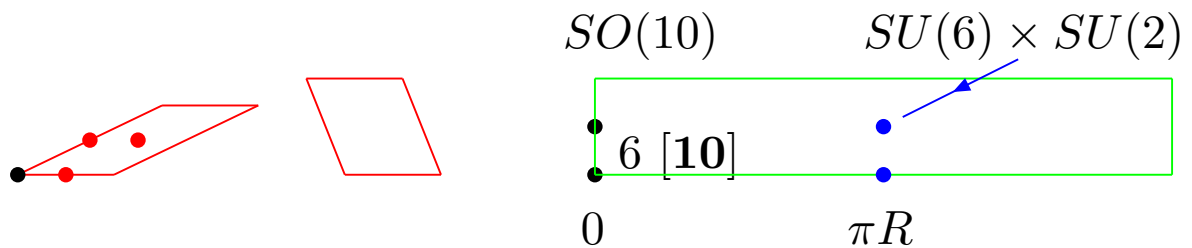
Table 1: Parities of the bulk states in model A1.

States	$P$	$P'$	States	$P$	$P'$
$V(\mathbf{15}, \mathbf{1}, \mathbf{1})$	+	+	$\Sigma(\mathbf{15}, \mathbf{1}, \mathbf{1})$	-	-
$V(\mathbf{1}, \mathbf{3}, \mathbf{1})$	+	+	$\Sigma(\mathbf{1}, \mathbf{3}, \mathbf{1})$	-	-
$V(\mathbf{1}, \mathbf{1}, \mathbf{3})$	+	+	$\Sigma(\mathbf{1}, \mathbf{1}, \mathbf{3})$	-	-
$V(\mathbf{6}, \mathbf{2}, \mathbf{2})$	+	-	$\Sigma(\mathbf{6}, \mathbf{2}, \mathbf{2})$	-	+
$V(\mathbf{4}, \mathbf{2}, \mathbf{1})$	-	+	$\Sigma(\mathbf{4}, \mathbf{2}, \mathbf{1})$	+	-
$V(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})$	-	+	$\Sigma(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})$	+	-
$V(\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	-	-	$U_3 \Sigma(\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	+	+
$V(\mathbf{4}, \mathbf{1}, \mathbf{2})$	-	-	$U_3 \Sigma(\mathbf{4}, \mathbf{1}, \mathbf{2})$	+	+
$U_1 H(\mathbf{4}, \mathbf{2}, \mathbf{1})$	+	+	$H^c(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})$	-	-
$H(\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})$	+	-	$H^c(\mathbf{4}, \mathbf{1}, \mathbf{2})$	-	+
$H(\mathbf{6}, \mathbf{1}, \mathbf{1})$	-	+	$H^c(\mathbf{6}, \mathbf{1}, \mathbf{1})$	+	-
$H(\mathbf{1}, \mathbf{2}, \mathbf{2})$	-	-	$U_2 H^c(\mathbf{1}, \mathbf{2}, \mathbf{2})$	+	+
$H(\mathbf{4}, \mathbf{2}, \mathbf{1})_+$	+	-	$H^c(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})_+$	-	+
$H(\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})_+$	+	+	$H^c(\mathbf{4}, \mathbf{1}, \mathbf{2})_+$	-	-
$H(\mathbf{6}, \mathbf{1}, \mathbf{1})_+$	-	-	$H^c(\mathbf{6}, \mathbf{1}, \mathbf{1})_+$	+	+
$H(\mathbf{1}, \mathbf{2}, \mathbf{2})_+$	-	+	$H^c(\mathbf{1}, \mathbf{2}, \mathbf{2})_+$	+	-
$H(\mathbf{4}, \mathbf{2}, \mathbf{1})_-$	-	+	$H^c(\bar{\mathbf{4}}, \mathbf{2}, \mathbf{1})_-$	+	-
$H(\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})_-$	-	-	$H^c(\mathbf{4}, \mathbf{1}, \mathbf{2})_-$	+	+
$H(\mathbf{6}, \mathbf{1}, \mathbf{1})_-$	+	+	$H^c(\mathbf{6}, \mathbf{1}, \mathbf{1})_-$	-	-
$H(\mathbf{1}, \mathbf{2}, \mathbf{2})_-$	+	-	$H^c(\mathbf{1}, \mathbf{2}, \mathbf{2})_-$	-	+

$T_{(1,2)}$  and  $T_{(1,0)}$  twisted sectors



$G_2 \oplus SU(3) \oplus SO(4)$  lattice with  $Z_6$  fixed points. The  $T_{(1,2)}$  twisted sector states sit at these fixed points.



$G_2 \oplus SU(3) \oplus SO(4)$  lattice with  $Z_2$  fixed points. The  $T_{(1,0)}$  twisted sector states sit at these fixed points.

$D_4$  family symmetry !

Note *unrequested*

$$[(6, 1, 1) \oplus (1, 2, 2)] \oplus [(4, 1, 1) \oplus (1, 2, 1)] \oplus (1, 1, 2) \oplus (1, 1, 1)$$

## Gauge coupling unification & Proton decay

- 5D RG equations

$M_s$  = string scale,  $M_{PS}$  = PS breaking scale,

$M_c$  = 5D compactification scale

$$\begin{aligned} \frac{2\pi}{\alpha_i(\mu)} &= \frac{2\pi}{\alpha_s} + b_i^{MSSM} \ln \frac{M_{PS}}{\mu} + (b_{++}^{PS} + b_{brane}^{PS})_i \ln \frac{M_s}{M_{PS}} \\ &- \frac{1}{2} (b_{++}^{PS} + b_{--}^{PS})_i \ln \frac{M_s}{M_c} + b^G \left( \frac{M_s}{M_c} - 1 \right) \end{aligned}$$

- 4D RG equations

$M_G \approx 3 \times 10^{16}$  GeV,  $\alpha_G^{-1} \approx 24$

and included threshold correction at  $M_G$ .

$$\frac{2\pi}{\alpha_i(\mu)} = \frac{2\pi}{\alpha_G} + b_i^{MSSM} \ln \frac{M_G}{\mu} + 6 \delta_{i3}$$

- $2\pi/\alpha_s = \pi/4 (M_{Planck}/M_s)^2$  \*\*

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Find —

$M_{PS} = e^{-3/2} M_G \sim 7 \times 10^{15}$  GeV,

$M_s(MAX) = e^2 M_G \sim 2 \times 10^{17}$  GeV

\*\*  $\implies \alpha_s^{-1} \sim 450$ ,  $\alpha_s$  too small — PROBLEM !

NO solution !!

Give up perturbative heterotic string  
boundary conditions

Eleven-dimensional Hořava-Witten

$$\frac{2\pi}{\alpha_s} = \frac{1}{2(4\pi)^{5/3} M \rho} \left( \frac{M_{\text{Pl}}}{M} \right)^2$$

$M$  — eleven-dimensional Newton's constant by  $\kappa^{2/3} = M^{-3}$

$\rho$  — size of the eleventh dimension

Now find solution

$$M_s \simeq M = 2M_G,$$

$$M_c \simeq M_{PS} = e^{-3/2} M_G$$

$$M\rho \simeq 2 \implies \rho \sim M_G^{-1}$$

◇ Enhanced proton decay rate — dimension-six operators

$$\tau(p \rightarrow e^+ \pi^0) = 3 \times 10^{33} \left( \frac{0.015 \text{ GeV}^3}{\beta_{\text{lattice}}} \right)^2 \text{ yrs}$$

Super-Kamiokande bound  $5.7 \times 10^{33}$  years @ 90% CL

## Yukawa couplings

$D_4 = \{\pm 1, \pm\sigma_1, \pm\sigma_3, \mp i\sigma_2\}$  family symmetry

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} : f_1 \leftrightarrow f_2$$

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} : f_2 \rightarrow -f_2$$

$$\begin{pmatrix} f_1 \\ f_2 \end{pmatrix} \text{ doublet; } f_3 \text{ singlet}$$

PS breaking VEVs

$$O_i = \langle \chi_\alpha^c \bar{\chi}_i^c \rangle, \quad i = 1, 2$$

• Fermion mass matrix [simple form]

$$(f_1 \ f_2 \ f_3) h \mathcal{M} \begin{pmatrix} f_1^c \\ f_2^c \\ f_3^c \end{pmatrix}$$

$$\mathcal{M} = \begin{pmatrix} (O_2 \tilde{S}_e + S_e) & (O_2 \tilde{S}_o + S_o) & (O_1 \ O_2 \ \phi_e + \tilde{\phi}_e) \\ (O_2 \tilde{S}_o + S_o) & (O_2 \tilde{S}_e + S_e) & (O_1 \ O_2 \ \phi_o + \tilde{\phi}_o) \\ \phi'_e & \phi'_o & 1 \end{pmatrix}$$

## Problems and Virtues

### • Virtues

- $E(6) \rightarrow SO(10) \rightarrow PS$
- Three families (**16** of  $SO(10)$ ) + Higgs (**10**) + PS breaking sector
- $D_4$  Family symmetry  $\implies$  hierarchy of masses and mixing
- Baryon and Lepton  $\neq$  violation in effective low energy field theory
 
$$f f f^c \langle \chi^c S^n \rangle \implies Q L \bar{D} + L L \bar{E}$$

$$f^c f^c f^c \langle \chi^c S^n \rangle \implies \bar{U} \bar{D} \bar{D}$$

$$f f f f \langle S^n \rangle \implies Q Q Q L ?$$
 NOT found to order  $S^8$  – inconsistent with string selection rules !

### • Problem

- Effective dimension 4 **R parity violating** operators via color triplet mixing
 
$$\langle \chi^c S^n \rangle C f^c + \langle S^m \rangle C C$$

$$\implies \hat{M}_{PS} T \bar{D} + \hat{M}_s T \bar{T} \text{ and } C = T + \bar{T}$$
 Implies – massless  $\bar{D}^0 \sim \bar{D} + \frac{\hat{M}_{PS}}{\hat{M}_s} \bar{T}$ 

$$C f^c f^c \implies \bar{D}^0 \bar{D} \bar{U}$$

$$C f f \implies \bar{D}^0 Q L$$
 B and L – **R parity violating** terms  $O(\hat{M}_{PS}/\hat{M}_s)$  — too large ??
- Vector-like exotics w/fractional charge need large mass  $O(M_G)$

## Conclusions

- Obtained UV completion of E(6) orbifold GUT in 5D
- Obtained cubic and higher order *effective Yukawa couplings*
- Need more study of baryon and lepton number violation
- We have two other three family models — one with SO(10) in bulk
- Just the beginning

Expand search to  $Z_N \otimes Z_2$  orbifolds +  
1 (or 2) Wilson lines in SO(4) direction  
 $\implies$  Effective 5 or 6 D orbifold GUTs

- Promising new directions for 3 family models !
- w/  $D_4$  family symmetry for fermion mass hierarchy and suppress flavor violation !
- and observable proton decay rate w/  $p \rightarrow \pi^0 e^+$  !!