# Conformal flatness of compact three-dimensional Cotton-parallel manifolds 

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#### Abstract

A three-dimensional pseudo-Riemannian manifold is called essentially conformally symmetric (ECS) if its Cotton tensor is parallel but nowherevanishing. In this note we prove that three-dimensional ECS manifolds must be noncompact or, equivalently, that every compact three-dimensional Cottonparallel pseudo-Riemannian manifold must be conformally flat.


## 1. Introduction and main result

Pseudo-Riemannian manifolds of dimensions $n \geq 4$ whose Weyl tensor is parallel are called conformally symmetric [3]. Those which are not locally symmetric or conformally flat are called essentially conformally symmetric (ECS, in short).

It has been shown by Roter in [9, Corollary 3] that ECS manifolds do exist in all dimensions $n \geq 4$, and in [4, Theorem 2] that they necessarily have indefinite metric signature. The local isometry types of ECS manifolds were described by Derdzinski and Roter in [6]. Compact ECS manifolds exist in all dimensions $n \geq 5$ and realize all indefinite metric signatures - see [8] and [7]. It is not currently known if compact four-dimensional ECS manifolds exist.

When the dimension of $M$ is $n \leq 3$, the Weyl tensor vanishes and this discussion becomes meaningless. In dimension $n=3$, however, conformal flatness is encoded in the Cotton tensor as opposed to the Weyl tensor, and so the following natural definition has been proposed in [1]: a three-dimensional pseudoRiemannian manifold is called conformally symmetric if its Cotton tensor is parallel, and those which are not conformally flat are then called ECS (note that every three-dimensional locally symmetric manifold is conformally flat). There, it is also

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shown [1, Theorem 1] that, reversing the metric if needed, any point in a threedimensional ECS manifold has a neighborhood isometric to an open subset of

$$
\begin{equation*}
(\widehat{M}, \widehat{\mathrm{~g}})=\left(\mathbb{R}^{3},\left(x^{3}+\mathfrak{a}(t) x\right) \mathrm{d} t^{2}+\mathrm{d} t \mathrm{~d} s+\mathrm{d} x^{2}\right) \tag{1.1}
\end{equation*}
$$

for some suitable smooth function $\mathfrak{a}: \mathbb{R} \rightarrow \mathbb{R}$. The coordinates $t$ and $s$ of $\widehat{M}$ are called $y$ and $t$ in [1], respectively, but have been renamed here as to make 1.1) directly resemble the corresponding local model given in [6. Section 4] for $n \geq 4$.

The pursuit of compact three-dimensional ECS manifolds quickly comes to an end in view of the following result, interesting on its own right without reference to ECS geometry:

ThEOREM A. A compact three-dimensional pseudo-Riemannian manifold with parallel Cotton tensor must be conformally flat.

While the compactness assumption here is crucial, Theorem A may be seen as a close relative (in general signature) of [2, Theorem 1]: compact Riemannian Cotton solitons are conformally flat, but nontrivial compact Lorentzian ones do exist.

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## 2. Preliminaries

Throughout this paper, we work in the smooth category and all manifolds considered are connected.
2.1. Symmetries of the Cotton tensor. The Cotton tensor of a $n$-dimensional pseudo-Riemannian manifold $(M, g)$ is the three-times covariant tensor field $C$ on $M$ defined by

$$
\begin{equation*}
\mathrm{C}(X, Y, Z)=\left(\nabla_{X} P\right)(Y, Z)-\left(\nabla_{Y} P\right)(X, Z), \quad \text { for } X, Y, Z \in \mathfrak{X}(M) \tag{2.1}
\end{equation*}
$$

Here, $P$ is the Schouten tensor of $(M, \mathrm{~g})$, given by

$$
\begin{equation*}
P=\operatorname{Ric}-\frac{\mathrm{s}}{2(n-1)} \mathrm{g} \tag{2.2}
\end{equation*}
$$

where Ric and s stand for the Ricci tensor and scalar curvature of $(M, \mathrm{~g})$, respectively. The Cotton tensor satisfies the following symmetries:
(i) $\mathrm{C}(X, Y, Z)+\mathrm{C}(Y, X, Z)=0$
(ii) $\mathrm{C}(X, Y, Z)+\mathrm{C}(Y, Z, X)+\mathrm{C}(Z, X, Y)=0$
(iii) $\operatorname{tr}_{\mathrm{g}}((X, Z) \mapsto \mathrm{C}(X, Y, Z))=0$
for all $X, Y, Z \in \mathfrak{X}(M)$. Symmetry (i) is obvious, while (ii) follows from a straightforward computation (six terms cancel in pairs), and (iii) from $\operatorname{div} P=\mathrm{d}\left(\operatorname{trg}_{\mathrm{g}} P\right)$
(which, in turn, is a consequence of the twice-contracted differential Bianchi identity $\operatorname{div}$ Ric $=\mathrm{ds} / 2$ ).
2.2. Algebraic structure in dimension 3. A routine computation shows that

$$
\begin{align*}
& \text { the Ricci and Cotton tensors of } 1.1) \text { are given } \\
& \text { by Ric }=-3 x \mathrm{~d} t \otimes \mathrm{~d} t \text { and } \mathrm{C}=3(\mathrm{~d} t \wedge \mathrm{~d} x) \otimes \mathrm{d} t \tag{2.4}
\end{align*}
$$

The expression for $C$ motivates the following result, analogous to [5, Lemma 17.1]:
Theorem 2.1. Let $(V,\langle\cdot, \cdot\rangle)$ be a three-dimensional pseudo-Euclidean space, and C be a nonzero Cotton-like tensor on $V$, i.e., a three-times covariant tensor on $V$ which formally satisfies 2.3, and consider $\mathcal{D}=\{u \in V \mid \mathrm{C}(u, \cdot, \cdot)=0\}$. Then:
(a) $\mathcal{D}$ consists only of null vectors, and hence $\operatorname{dim} \mathcal{D} \leq 1$.
(b) $\operatorname{dim} \mathcal{D}=1$ if and only if $C=(u \wedge v) \otimes u$ for some $u \in \mathcal{D} \backslash\{0\}$ and unit $v \in \mathcal{D}^{\perp}$.
(c) In (b), $u$ is unique up to a sign, while $v$ is unique modulo $\mathcal{D}$.

Here, we identify $V \cong V^{*}$ with the aid of $\langle\cdot, \cdot\rangle$.

PROOF. For (a), assuming by contradiction the existence of a unit vector $e_{1} \in \mathcal{D}$, we will show that $\mathrm{C}=0$. Considering an orthonormal basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ for $(V,\langle\cdot, \cdot\rangle)$ and using 2.3 i ) and 2.3 ii ), we see that

$$
\begin{equation*}
C_{i j k} \text { is only possibly nonzero when }\{i, j, k\}=\{2,3\} \text { with } i \neq j \tag{2.5}
\end{equation*}
$$

Now $\mathrm{C}_{322}=-\mathrm{C}_{232}$ and $\mathrm{C}_{323}=-\mathrm{C}_{233}$, while $\operatorname{tr}_{\langle, \cdot\rangle}\left(\left(w, w^{\prime}\right) \mapsto \mathrm{C}\left(e_{j}, w, w^{\prime}\right)\right)=0$ for $j=2$ and $j=3$ readily yields $C_{233}=0$ and $C_{322}=0$, respectively. Hence $C=0$, as claimed. As for (b), assume that $\operatorname{dim} \mathcal{D}=1$, fix a null vector $e_{1} \in \mathcal{D} \backslash\{0\}$, and complete it to a basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ of $V$ satisfying the relations

$$
\begin{equation*}
\left\langle e_{1}, e_{2}\right\rangle=\left\langle e_{2}, e_{3}\right\rangle=\left\langle e_{3}, e_{3}\right\rangle=0 \quad \text { and } \quad\left\langle e_{1}, e_{3}\right\rangle=\left\langle e_{2}, e_{2}\right\rangle=(-1)^{q+1} \tag{2.6}
\end{equation*}
$$

where $q \in\{1,2\}$ is the index of $\langle\cdot, \cdot\rangle$. By the same argument as in (a), we again obtain (2.5), but this time $\operatorname{tr}_{\langle\cdot, \cdot\rangle}\left(\left(w, w^{\prime}\right) \mapsto \mathrm{C}\left(e_{3}, w, w^{\prime}\right)\right)=0$ reduces to $\mathrm{C}_{232}=0$ in view of 2.6. Writing $a=\mathrm{C}_{323} \neq 0$ for the last essential component of C , it follows that $\mathrm{C}=a\left(e^{3} \wedge e^{2}\right) \otimes e^{3}$, where $\left\{e^{1}, e^{2}, e^{3}\right\}$ is the basis of $V^{*}$ dual to $\left\{e_{1}, e_{2}, e_{3}\right\}$. Applying the isomorphism $V \cong V^{*}$ and setting $u=|a|^{1 / 2} e_{1}$ and $v=\operatorname{sgn}(a) e_{2}$, we obtain the required expression $\mathrm{C}=(u \wedge v) \otimes u$. Conversely, it is straightforward to verify that the tensor $(u \wedge v) \otimes u$ with $u$ null and $v$ unit and orthogonal to $u$ is Cotton-like with $\mathcal{D}=\mathbb{R} u$ and $\mathcal{D}^{\perp}=\mathbb{R} u \oplus \mathbb{R} v$. Finally, (c) is clear from (b).

As a consequence, whenever $(M, \mathrm{~g})$ is a three-dimensional pseudo-Riemannian manifold, we may assign to each point $x \in M$ the kernel $\mathcal{D}_{x}$ of $C_{x}$ in $\left(T_{x} M, g_{x}\right)$. In
the ECS case, we have that
$\mathcal{D}$ is a smooth rank-one parallel distribution on $M$, which contains the image of the Ricci endomorphism of $(M, \mathrm{~g})$.
Indeed, we may note that (2.7) holds in the model (1.1) (as (2.4) gives us that $\mathcal{D}$ is spanned by the coordinate vector field $\partial_{s}, \widehat{\mathrm{~g}}$-dual to $\mathrm{d} t$ up to a factor of 2 ), and invoke [1, Theorem 1].

## 3. Proof of Theorem A

In this section, we fix a compact three-dimensional ECS manifold $(M, \mathrm{~g})$ and its universal covering manifold $\pi: \widetilde{M} \rightarrow M$, which equipped with the natural pull-back metric $\widetilde{g}=\pi^{*} g$ becomes an ECS manifold. We will use the same symbols Ric, $P, C, \nabla$, and $\mathcal{D}$ for the corresponding objects in both $(M, \mathrm{~g})$ and $(\tilde{M}, \widetilde{\mathrm{~g}})$. Observe that
the fundamental group $\Gamma=\pi_{1}(M)$ acts properly discontinu-
ously on ( $\widetilde{M}, \widetilde{\mathrm{~g}})$ by deck isometries, with quotient $\widetilde{M} / \Gamma \cong M$.
As $\widetilde{M}$ is simply connected, we may fix two globally defined smooth vector fields $u$ and $v$ such that $\mathrm{C}=(\boldsymbol{u} \wedge \boldsymbol{v}) \otimes \boldsymbol{u}$ on $\widetilde{M}$. Now, as $\mathcal{D}$ is parallel, item (c) of Theorem 2.1 gives us that
(i) $u$ is a null parallel vector field spanning $\mathcal{D}$;
(ii) every $\gamma \in \Gamma$ either pushes $\boldsymbol{u}$ forward onto itself or onto its opposite.

Next, as the Ricci endomorphism of ( $\widetilde{M}, \widetilde{\mathrm{~g}})$ is self-adjoint, $(2.7)$ allows us to write

$$
\begin{equation*}
\text { Ric }=-f \boldsymbol{u} \otimes \boldsymbol{u} \text {, for some smooth function } f: \widetilde{M} \rightarrow \mathbb{R} . \tag{3.3}
\end{equation*}
$$

By (3.3) and (3.2-i), ( $\tilde{M}, \tilde{g})$ is scalar-flat, and so $P=$ Ric. Combining this with $(3.2-\mathrm{i})$ again to compute $C$ via (2.1), we obtain that

$$
\begin{equation*}
\mathrm{C}=(\boldsymbol{u} \wedge \nabla f) \otimes \boldsymbol{u}, \text { where } \nabla f \text { is the } \widetilde{\mathrm{g}} \text {-gradient of } f \text {. } \tag{3.4}
\end{equation*}
$$

However, it follows from (3.2-ii) and (3.3) that $f$ is $\Gamma$-invariant, and so it has a critical point due to (3.1) and compactness of $M$. Such a critical point is in fact a zero of C by $\sqrt{3.4}$, and therefore $\mathrm{C}=0$. This is the desired contradiction: $(M, \mathrm{~g})$ must be either noncompact, or conformally flat.

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