# THE RICCI IDENTITY

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#### 1 Setup

We work on the smooth category, and fix a vector bundle  $E \to M$  equipped with a connection  $\nabla \colon \mathfrak{X}(M) \times \Gamma(E) \to \Gamma(E)$ , where  $\Gamma(E)$  stands for the  $\mathbb{C}^{\infty}(M)$ -module of smooth sections of E. The curvature tensor of  $\nabla$  is  $R \colon \mathfrak{X}(M) \times \mathfrak{X}(M) \times \Gamma(E) \to \Gamma(E)$ defined by

$$R^{\nabla}(X,Y)\psi = \nabla_X \nabla_Y \psi - \nabla_Y \nabla_X \psi - \nabla_{[X,Y]} \psi.$$
(1.1)

Relative to a local trivialization  $e_a$  for *E* and local coordinates  $x^j$  for *M*, we write

(a) 
$$\nabla_{\partial_j} e_a = \Gamma_{ja}^b e_b$$
 (b)  $R(\partial_j, \partial_k) e_a = R_{jka}^{\ \ b} e_b$ , (1.2)

while repeatedly substituting (1.2-a) into (1.1) to compute (1.2-b) yields

$$R_{jka}^{\ \ b} = \partial_j \Gamma_{ka}^b - \partial_k \Gamma_{ja}^b + \Gamma_{ka}^c \Gamma_{jc}^b - \Gamma_{ja}^c \Gamma_{kc}^b.$$
(1.3)

Whenever  $\psi \in \Gamma(E)$ , we may consider  $\nabla \psi$  as a section of Hom(*TM*, *E*). Applying the Leibniz rule for  $\nabla$  together with (1.2-a), we have that

$$\nabla_{\partial_j}\psi = (\partial_j\psi^b + \Gamma^b_{ja}\psi^a)e_b, \tag{1.4}$$

and this motivates defining

$$\nabla_{j}\psi^{b} = \partial_{j}\psi^{b} + \Gamma^{b}_{ja}\psi^{a}, \text{ so that } \nabla\psi = \nabla_{j}\psi^{b}dx^{j} \otimes e_{b}.$$
(1.5)

Another notation for  $\nabla_j \psi^b$  is  $\psi^b_{jj'}$  but I prefer the former to the latter as it allows us to think of " $\nabla_j = \partial_j + \Gamma_j$ " as an operator on its own right.

The vector-bundle index *a*, for concrete choices of *E*, may actually stand for a collection of indices coming from *M*. For example, assume that  $E = TM^{\otimes 2}$  has a connection induced from some connection in *TM*. If  $e_a$  is a local frame for *TM*, then  $e_a \otimes e_b$  forms a local trivialization for *E*, in which case we have that

$$\begin{aligned} \nabla_{\partial_j}(e_a \otimes e_b) &= \nabla_{\partial_j} e_a \otimes e_b + e_a \otimes \nabla_{\partial_j} e_b \\ &= \Gamma_{ja}^c e_c \otimes e_b + e_a \otimes \Gamma_{jb}^d e_d \\ &= (\Gamma_{ja}^c \delta_b^d + \delta_a^c \Gamma_{jb}^d) e_c \otimes e_d \end{aligned} \tag{1.6}$$

says that  $\Gamma_{j,\underline{a}\underline{b}}^{c\underline{d}} = \Gamma_{ja}^{c}\delta_{b}^{d} + \delta_{a}^{c}\Gamma_{jb}^{d}$ , where the underlines on  $\Gamma_{j,\underline{a}\underline{b}}^{c\underline{d}}$  are meant to emphasize that ab is a single  $TM^{\otimes 2}$ -index, as is cd. The point of observing this is that it suffices to develop the theory for vector bundles, and then it may be applied to any tensor bundle over M.

## 2 Second derivatives

When *TM* is also equipped with a connection  $\nabla^{\circ}$ , the bundle Hom(*TM*, *E*) inherits a connection  $\nabla'$  from both  $\nabla$  and  $\nabla^{\circ}$ . More precisely, if  $e_a$  is a local trivialization for *E* and  $x^j$  are local coordinates for *M*, then we may write

$$F = F_k^a \, \mathrm{d} x^k \otimes e_a, \text{ where } F_k^a \text{ is the } a\text{-th component of } F(\partial_k), \tag{2.1}$$

for any section *F* of Hom(*TM*, *E*). The pair  ${}^{a}_{k}$  is a single *upper* Hom(*TM*, *E*)-index (while  ${}^{k}_{a}$  as in  $dx^{k} \otimes e_{a}$  would be a *lower* one). By definition of the induced connection on Hom(*TM*, *E*), we have that

$$(\nabla'_X F)(Y) = \nabla_X (F(Y)) - F(\nabla^\circ_X Y)$$
(2.2)

Mimicking what was done in (1.4), we set  $X = \partial_i$  and  $Y = \partial_k$  on (2.2) and compute

$$(\nabla_{\partial_{j}}^{\prime}F)(\partial_{k}) = \nabla_{\partial_{j}}(F(\partial_{k})) - F(\nabla_{\partial_{j}}^{\circ}\partial_{k})$$

$$= \nabla_{\partial_{j}}(F_{k}^{a}e_{a}) - F(\mathring{\Gamma}_{jk}^{\ell}\partial_{\ell})$$

$$= (\partial_{j}F_{k}^{a})e_{a} + F_{k}^{a}\partial_{j}e_{a} - \mathring{\Gamma}_{jk}^{\ell}F(\partial_{\ell})$$

$$= (\partial_{j}F_{k}^{a})e_{a} + F_{k}^{b}\Gamma_{jb}^{a}e_{a} - \mathring{\Gamma}_{jk}^{\ell}F_{\ell}^{a}e_{a}$$

$$= (\partial_{j}F_{k}^{a} + F_{k}^{b}\Gamma_{jb}^{a} - \mathring{\Gamma}_{jk}^{\ell}F_{\ell}^{a})e_{a},$$

$$(2.3)$$

giving us that

$$\nabla'_j F^a_k = \partial_j F^a_k + F^b_k \Gamma^a_{jb} - \mathring{\Gamma}^\ell_{jk} F^a_\ell.$$
(2.4)

Usually, the left side of the above is written just as  $\nabla_j F_k^a$ , with the dependence of the right side on the auxiliary connection  $\nabla^\circ$  being understood. Now, we may take  $F = \nabla \psi$  for some section  $\psi \in \Gamma(E)$ , so that

$$\nabla_{j}\nabla_{k}\psi^{a} = \partial_{j}(\nabla_{k}\psi^{a}) + (\nabla_{k}\psi^{b})\Gamma^{a}_{jb} - \mathring{\Gamma}^{\ell}_{jk}(\nabla_{\ell}\psi^{a})$$
(2.5)

Substitute (1.5) into the first two terms on the right side of (2.5) to compute

$$\nabla_{j}\nabla_{k}\psi^{a} = \partial_{j}(\partial_{k}\psi^{a} + \Gamma^{a}_{kb}\psi^{b}) + (\partial_{k}\psi^{b} + \Gamma^{b}_{kc}\psi^{c})\Gamma^{a}_{jb} - \mathring{\Gamma}^{\ell}_{jk}(\nabla_{\ell}\psi^{a})$$

$$= \partial_{j}\partial_{k}\psi^{a} + (\partial_{j}\Gamma^{a}_{kb})\psi^{b} + \Gamma^{a}_{kb}\partial_{j}\psi^{b} + \Gamma^{a}_{jb}\partial_{k}\psi^{b} + \Gamma^{a}_{jb}\Gamma^{b}_{kc}\psi^{c} - \mathring{\Gamma}^{\ell}_{jk}(\nabla_{\ell}\psi^{a})$$

$$= (\partial_{j}\partial_{k}\psi^{a} + \Gamma^{a}_{kb}\partial_{j}\psi^{b} + \Gamma^{a}_{jb}\partial_{k}\psi^{b}) + (\partial_{j}\Gamma^{a}_{kb})\psi^{b} + \Gamma^{a}_{jb}\Gamma^{b}_{kc}\psi^{c} - \mathring{\Gamma}^{\ell}_{jk}(\nabla_{\ell}\psi^{a}).$$
(2.6)

Noting that the first group of three terms in the last line of (2.6) is symmetric in the pair (j, k), we may directly obtain that

$$\nabla_{j}\nabla_{k}\psi^{a} - \nabla_{k}\nabla_{j}\psi^{a} = (\partial_{j}\Gamma^{a}_{kb} - \partial_{k}\Gamma^{a}_{jb})\psi^{b} + (\Gamma^{a}_{jb}\Gamma^{b}_{kc} - \Gamma^{a}_{kb}\Gamma^{b}_{jc})\psi^{c} - \tau^{\ell}_{jk}\nabla_{\ell}\psi^{a}, \qquad (2.7)$$

where  $\tau_{ik}^{\ell}$  are the components of the torsion tensor of  $\nabla^{\circ}$ , defined by

$$\tau(X,Y) = \nabla_X^\circ Y - \nabla_Y^\circ X - [X,Y].$$
(2.8)

If  $\nabla^{\circ}$  is torsionfree (i.e.,  $\tau = 0$ ), substituting (1.3) into (2.7) and renaming  $b \leftrightarrow c$ , it follows that

$$\nabla_j \nabla_k \psi^a - \nabla_k \nabla_j \psi^a = R_{jkb}^{\ a} \psi^b, \qquad (2.9)$$

as required.

### **3** Application to tensor bundles

In this section, we let  $\nabla$  be a *torsionfree* connection on M, and again denote by  $\nabla$  the induced connections on all tensor bundles over M. We also choose the "auxiliary" connection to be  $\nabla^{\circ} = \nabla$ . Set  $E = TM^{\otimes r} \otimes T^*M^{\otimes s}$ . A *r*-times contravariant and *s*-times covariant tensor field on M is a section of E, which may be written relative to a local coordinate system  $x^j$  for M as

$$T = T_{j_1 \dots j_s}^{i_1 \dots i_r} \partial_{i_1} \otimes \dots \otimes \partial_{i_r} \otimes \mathrm{d} x^{j_1} \otimes \dots \otimes \mathrm{d} x^{j_s}. \tag{3.1}$$

As we are not dealing with metrics and the process of raising/lowering indices, there is no need to stick to the more precise notation  $T^{i_1...i_r}_{j_1...j_s}$ . To apply (2.9), we need an expression for the curvature tensor of the induced connection on *E*. Such expression, however, will follow from the Leibniz rule combined with relations

(a) 
$$R(\partial_j, \partial_k)\partial_{i_a} = R_{jki_a}^{\ \ p}\partial_p$$
 (b)  $R(\partial_j, \partial_k)dx^{j_b} = -R_{jkq}^{\ \ b}dx^q$  (3.2)

from the base cases (r, s) = (1, 0) and (r, s) = (0, 1). Indeed, we have that

$$R(\partial_{j},\partial_{k})\left(\partial_{i_{1}}\otimes\cdots\otimes\partial_{i_{r}}\otimes dx^{j_{1}}\otimes\cdots\otimes dx^{j_{s}}\right) = \\ = (R_{jki_{1}}^{p_{1}}\partial_{p_{1}})\otimes\cdots\otimes\partial_{i_{r}}\otimes dx^{j_{1}}\otimes\cdots\otimes dx^{j_{s}} \\ +\cdots+\partial_{i_{1}}\otimes\cdots\otimes(R_{jki_{r}}^{p_{r}}\partial_{p_{r}})\otimes dx^{j_{1}}\otimes\cdots\otimes dx^{j_{s}} \\ +\partial_{i_{1}}\otimes\cdots\otimes\partial_{i_{r}}\otimes(-R_{jkq_{1}}^{j_{1}}dx^{q_{1}})\otimes\cdots\otimes dx^{j_{s}} \\ +\cdots+\partial_{i_{1}}\otimes\cdots\otimes\partial_{i_{r}}\otimes\cdots\otimes dx^{j_{1}}\otimes\cdots\otimes(-R_{jkq_{s}}^{j_{s}}dx^{q_{s}}).$$

$$(3.3)$$

With the aid of Kronecker deltas, we may informally rewrite (3.3) as

$$\frac{R(\partial_{j},\partial_{k})\left(\partial_{i_{1}}\otimes\cdots\otimes\partial_{i_{r}}\otimes dx^{j_{1}}\otimes\cdots\otimes dx^{j_{s}}\right)}{\partial_{p_{1}}\otimes\cdots\otimes\partial_{p_{r}}\otimes dx^{q_{1}}\otimes\cdots\otimes dx^{q_{s}}} = 
= \sum_{\ell=1}^{r} R_{jki_{\ell}}{}^{p_{\ell}}\delta_{i_{1}}{}^{p_{1}}\dots\widehat{\delta_{i_{\ell}}}{}^{p_{\ell}}\dots\delta_{i_{r}}{}^{p_{r}}\delta_{q_{1}}{}^{j_{1}}\dots\delta_{q_{s}}{}^{j_{s}} - \sum_{\ell=1}^{s} R_{jkq_{\ell}}{}^{j_{\ell}}\delta_{i_{1}}{}^{p_{1}}\dots\delta_{i_{1}}{}^{p_{r}}\delta_{q_{1}}{}^{j_{1}}\dots\widehat{\delta_{q_{\ell}}}{}^{q_{\ell}}\dots\delta_{q_{s}}{}^{j_{s}}.$$
(3.4)

The right side of (3.4) is the curvature term present in the right side of (2.9) for our choice of *E* (after suitably adjusting dummy indices). Already contracting all the Kronecker deltas possible, we obtain that

$$\nabla_{j}\nabla_{k}T_{j_{1}\dots j_{s}}^{i_{1}\dots i_{r}} - \nabla_{k}\nabla_{j}T_{j_{1}\dots j_{s}}^{i_{1}\dots i_{r}} = R_{jkp}^{i_{1}}T_{j_{1}\dots j_{s}}^{pi_{2}\dots i_{r}} + \dots + R_{jkp}^{i_{r}}T_{j_{1}\dots j_{s}}^{i_{1}\dots i_{r-1}p} - R_{jkj_{1}}^{i_{1}}T_{qj_{2}\dots j_{s}}^{i_{1}\dots i_{r}} - \dots - R_{jkj_{s}}^{i_{s}}T_{j_{1}\dots j_{s-1}q}^{i_{1}\dots i_{r}}.$$
(3.5)

The placement of signs follows (3.2). We list some particular cases of (3.5) below for the reader's convenience:

(i) 
$$\nabla_k \nabla_\ell T_{ij} - \nabla_\ell \nabla_k T_{ij} = -R_{k\ell i}^{\ \ q} T_{qj} - R_{k\ell j}^{\ \ q} T_{iq}.$$

(ii) 
$$\nabla_k \nabla_\ell T_i^r - \nabla_\ell \nabla_k T_i^r = R_{k\ell p}^{\ r} T_i^p - R_{k\ell i}^{\ q} T_q^r.$$

(iii) 
$$\nabla_{\ell} \nabla_{m} T_{ij}^{k} - \nabla_{m} \nabla_{\ell} T_{ij}^{k} = R_{\ell m p}^{\ \ k} T_{ij}^{p} - R_{\ell m i}^{\ \ q} T_{qj}^{k} - R_{\ell m j}^{\ \ q} T_{iq}^{k}.$$

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### 4 Geometric consequences

Here, we let (M, g) be a *n*-dimensional pseudo-Riemannian manifold, and  $\nabla$  be its Levi-Civita connection. One consequence of the Ricci identity (3.5) is that

if (M, g) has nonzero constant sectional curvature, then M does not admit nontrivial parallel vector fields, or parallel 1-forms. (4.1)

If X is a parallel vector field on M, then  $R_{ijp}^{\ \ k}X^p = 0$  reads  $K(g_{jp}\delta_i^k - g_{ip}\delta_j^k)X^p = 0$ , where  $K \neq 0$  is the sectional curvature of (M, g). Evaluating it and cancelling K gives us that  $\delta_i^k X_j - \delta_j^k X_i = 0$ , and making k = i yields  $(n - 1)X_i = 0$ , so that  $X_i = 0$ . The argument for 1-forms  $\alpha$  is dual:  $R_{ijk}^{\ \ q}\alpha_q = 0$  reads  $K(g_{jk}\delta_i^q - g_{ik}\delta_j^q)\alpha_q = 0$ , and therefore  $g_{jk}\alpha_i - g_{ik}\alpha_j = 0$ . Contracting against  $g^{jk}$  leads to  $(\dim M - 1)\alpha_i = 0$ , and hence  $\alpha_i = 0$ .